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Tuning the work function of Si(100) surface by halogen absorption: a DFT study

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<p>Please submit a plain text version of your cover letter here.</p> <p>Please note, if you are submitting a revision of your manuscript, there is an opportunity for you to provide your responses to the reviewers later; please do not add them to the cover letter.</p>	<p>Daniel Hiller, Dirk König, Walter Weber and Zachary Holman Guest Editors E-MRS 2017 Spring Symposium P Proceedings in pss (c) Silicon & Silicon nanostructures: from recent fundamental research to novel applications</p> <p>27 June 2017</p> <p>Dear colleagues,</p> <p>Please find attached the original research article entitled Tuning the work function of Si(100) surface by halogen absorption: a DFT study by M. Bertocchi, M. Amato, I. Marri and S. Ossicini, which we would like to be considered for publication as E-MRS 2017 Spring Symposium P Proceedings in Physica Status Solidi (c).</p> <p>Recently, the study of the surface passivation of group IV semiconductor surfaces is back to being a hot topic of debate because of its wide range of technological applications in particular in the field of nanoelectronics, sensing and energy. Accurate knowledge of the physical and chemical process occurring at the interface between the passivating atom and the semiconductor is crucial for the designing novel and efficient devices.</p> <p>Here, we contribute to this research field by employing atomistic Density Functional Theory (DFT) simulations in order to investigate the halogen absorption effect on the Si (100) surface. Our results, in agreement with state-of-the-art experiments, shed light on two main aspects:</p> <p>1)First, the presence of halogen atoms (I, Br, Cl, and F) leads to an increase of the work function if compared to the fully hydrogenated surface. This is due to their high electronegativity that causes a large charge transfer with an associated dipole moment.</p> <p>2)Second, the generality of this behavior is confirmed by the discussion and comparison of our results with other calculations and experiments on both Si(100) and Si(111). This sheds light on the major role of the type of passivation with respect to the surface crystal orientation.</p> <p>Due to the great interest of our results and their significance for future theoretical and experimental investigation on chemical functionalization of semiconductor surfaces, we do believe that our manuscript meets the criteria of novelty and significance to justify publication in Physica Status Solidi C.</p> <p>We are looking for a rapid decision on our manuscript.</p>

	Best regards, Stefano Ossicini
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Abstract:	<p>First-principles calculations of work function tuning induced by different chemical terminations on Si(100) surface are presented and discussed. We find that the presence of halogen atoms (I, Br, Cl, and F) leads to an increase of the work function if compared to the fully hydrogenated surface. This is a quite general effect and is directly linked to the chemisorbed atoms electronegativity as well as to the charge redistribution at the interface. All these results are examined with respect to previous theoretical works and experimental data obtained for the (100) as well as other Si surface orientations. Based on this analysis, we argue that the changes in the electronic properties caused by variations of the interfacial chemistry strongly depend on the chemisorbed species and much less on the surface crystal orientation.</p>

Tuning the work function of Si(100) surface by halogen absorption: a DFT study

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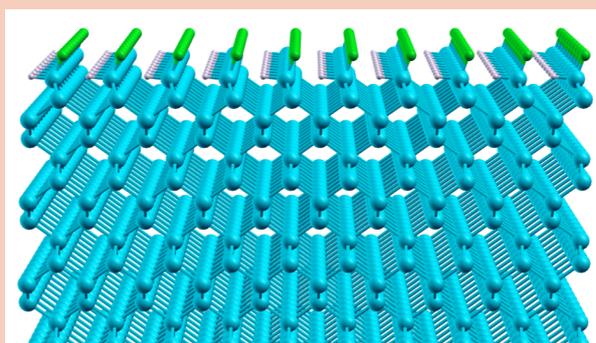
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First-principles calculations of work function tuning induced by different chemical terminations on Si(100) surface are presented and discussed. We find that the presence of halogen atoms (I, Br, Cl, and F) leads to an increase of the work function if compared to the fully hydrogenated surface. This is a quite general effect and is directly linked to the chemisorbed atoms electronegativity as well as to the charge redistribution at the interface. All these results are examined with respect to previous theoretical works and experimental data obtained for the (100) as well as other Si surface orientations. Based on this analysis, we argue that the changes in the electronic properties caused by variations of the interfacial chemistry strongly depend on the chemisorbed species and much less on the surface crystal orientation.



A Si(100) surface supercell passivated by H and F atoms. Light blue spheres represent Si atoms, grey spheres H atoms and green spheres F atoms.

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1 Introduction Silicon is one of the most used semiconductors. It is cheap, non-toxic and largely engineered in electronic and micro-electronic devices for applications in different fields, from photonics to photovoltaics [1,2]. In the last thirty years a large number of experimental and theoretical works have been dedicated to the study of the electronic and optical properties of Si-based systems with different dimensionality.

In this context first principles techniques, based on the density functional theory (DFT), have emerged as efficient tools that can be used to complement experimental activ-

ities. The application of DFT, sometimes combined with advanced methodologies including many-body effects has allowed to increase our knowledge about the structure, electronic, optical and transport properties of different semiconductor materials [3–5], to investigate how such properties are modified by strain, passivation, and doping [6,7] and to predict and quantify new effects, as the second-harmonic generation in silicon [8] and the carrier multiplication in isolated and interacting nanocrystals [9–12].

In this work we investigate the (100) surface of silicon.

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This system exhibits a variety of reconstructions that are mainly related to the different ordering of surface dimers. Since the Si(100) system is largely used in micro- and nano-electronic industry, strong efforts have been dedicated to the study of its chemical and physical properties [13–15]. From a broader perspective the study of the chemical passivation of group IV semiconductor surfaces is back to being a hot topic of debate. In particular, the analysis of the electronic modifications at the semiconductor surface induced by the adsorption/chemisorption of atoms and molecules has been the object of many different works. It has been pointed out that the functionalization of these surfaces is a key process in micro- and nano-electronics, energy conversion, charge storage, process information, sensing and electrochemical catalysis [14, 16–21]. Indeed, chemisorption of atoms and molecules on semiconductor surfaces has a strong influence on some properties that are fundamental for technological applications, i.e. chemical reactivity, surface conductance, band edge profile and work function (WF). More specifically, the possibility to tune ionization potential, electron affinity and WF of both the (111) and (100) surfaces of silicon [21–38] has been investigated and exploited.

The aim of this work is to study, within DFT, the effect induced by the chemisorption of halogen atoms on the electronic properties of the Si(100) surface. In particular, we aim at investigating, on an atomic scale, the mechanisms that are responsible of the WF modulation. The accurate determination of the WF provides an absolute electron energy level reference relative to the vacuum energy, which is important for applications like hetero-interfaces based devices, electrocatalysts, oxide and graphene based electronics, solar cells, Schottky junctions and thermionic devices. The term of comparison throughout our work is the hydrogenated surface, which often represents the starting point for the preparation of samples with different chemical termination [28, 35]. In this work we consider only the unreconstructed (1x1) Si(100) surface. This is because of two main reasons: firstly, as often observed in experiments, the (2x1) reconstruction, upon monolayer chemical absorption, undergoes a *deconstruction* to the (1x1) structure, with one adatom present for each semiconductor surface atom [13, 39, 40]; this condition makes possible a direct comparison between our results and experimental measurements. Secondly, without the need of considering complicated reconstructions, this choice renders our calculations less computationally demanding and less cumbersome.

2 Theoretical Approach Si(100) surface was modeled through the supercell method using a slab of n semiconductor layers, grown along the (100) direction. In our calculations we adopt the bulk Si lattice parameter ($a_{Si\ bulk}=5.40$ Å). The determination of n is generally done to satisfy specific requirements that are connected

with both the system and the physical properties under consideration; depending on them, n can be tuned on a large range of values. The energy gap E_G , for instance, strongly depends on the slab thickness, as consequence of the quantum confinement effect [41–43]; a detailed determination of E_G requires therefore an accurate convergence on n . The same considerations hold for the WF that, in general, requires a large number of layers to be converged. To avoid any ambiguity all the results presented in this work refer to a slab made of $n = 36$ layers, that corresponds to a slab thickness of about 50 Å. By assuming $n = 36$, we can correctly reproduce the electronic structure of bulk silicon thus ensuring an achieved convergence for all the calculated properties.

In our simulations, Si slabs have been embedded in very large supercells, in order to avoid spurious interactions among periodic replicas that may influence the computed electronic properties. In particular, we have inserted a vacuum region of more than 35 Å [44, 45]. Moreover, surfaces have been modeled by adopting centrosymmetric slabs that keeps from the formation of spurious and artificial electric fields components. Since we considered an unreconstructed Si(100) surface, both the top and the bottom surface layers of our slabs contain two-dangling bonds. Initially, all these dangling bonds have been saturated with hydrogen atoms. This hydrogenated structure represents our reference system, accordingly with the recipes often used in experiments (see Section 1). In order to understand how the properties of the Si surface depend on the atomic termination, we have investigated surfaces terminated with different adsorbed species, labeled by X . These species replace one of the initial capping H atoms placed at the top and bottom surface sites. The substitution is accurately done to preserve both the symmetry and charge neutrality of the supercell. The structures are then relaxed, keeping the atomic position of the central layers fixed, in order to reproduce a bulk-like region in the slab. In our investigation different X terminations have been chosen by considering a series of halogen atoms (I, Br, Cl, and F) with different electronegativity. As previously done in several theoretical and experimental studies concerning the Si(111) surface [35, 36], we discuss our results in term of electronegativity and charge rearrangement at the surface.

The calculations were performed employing the plane-wave pseudopotential PWscf code of the QuantumESPRESSO distribution [46]. This code is largely used by scientific community (and also by our group [47–51]) and allows to calculate structural and electronic properties of metal and semiconductor materials using different exchange-correlation functionals. In our work we adopt the local density approximation (LDA) and we employed norm-conserving pseudopotential for all the chemical elements considered. As is known, in semiconductors, LDA severely underestimates the electronic band gap with respect to the experimental values. For the bulk silicon, for instance, we obtain $E_{G,bulk}^{Si} = 0.56$ eV, to be compared to

the experimental $E_{G,bulk,exp}^{Si} = 1.12$ eV [44]. This underestimation can be, in principle, overcome by introducing quasiparticle corrections using the so-called GW approximation [52,53]. This procedure is, however, quite demanding from a computational point of view. The accuracy of the method employed in this work relies on the fact that in several cases the use of the GW correction results in an almost rigid shift of the band, corresponding to a constant opening of the gap value. Moreover, as proven in previous works [35], the WF can be accurately determined also in the framework of the DFT at the LDA level. Being the scope of this manuscript the study of the WF in Si(100) surfaces with different passivations, we hence limit our approach to the DFT-LDA.

Noticeably, the calculated values for the band gaps of Si surfaces with different passivation are very similar (within 40 meV) to the one obtained for the Si bulk. This result points out that, for the considered systems, passivation do not influence the energy gap and also that our slabs correctly describe the bulk properties of the surface.

A careful analysis of the convergence of both the electronic and structural properties with respect to the plane-wave basis set cutoff has been conducted. In particular the kinetic cut-off for the plane-wave basis has been set equal to 35 Ry and all the atomic positions in the supercell were relaxed until the forces acting on each atom were less than 0.003 Ry/a.u. All the results discussed in the next sections concern the electronic structure calculated for the optimized geometries.

3 Electronic structure and work function

The WF of a material is usually defined as the minimum energy required to remove an electron from the bulk to the vacuum outside across the surface. Consequently, it can be calculated as the difference between the vacuum energy level E_{vac} and the Fermi energy E_F of the system under consideration:

$$WF = E_{vac} - E_F \quad (1)$$

In order to calculate the surface WF by employing this equation, accurate values of both the vacuum potential and the Fermi energy are needed. In a semiconductor, since there are no allowed electronic energy levels in the gap, the Fermi energy is a somewhat theoretical construct. Moreover, in a vacuum-slab-supercell calculation, the E_F can be calculated in different ways. For example, one can consider as Fermi energy the value directly derived from the slab calculation [29,54]. On the other hand, one could employ the Fermi energy of the corresponding bulk system whose value will not suffer from finite-size effect [37,54,55]. We have tested both the methods and proved that, for the considered systems, they give quite similar results if, in the first case, a slab formed by a sufficient number of layers and an adequate vacuum thicknesses are taken into account.

In this manuscript WF is calculated considering the E_F of bulk Si and using a procedure consisting of few simple steps [56]: (i) firstly, starting from the unrelaxed passi-

ivated surface, we remove hydrogens, halogen atoms and the vacuum region, thus building a bulk Si supercell that contains the same number of Si atoms of the passivated surface. For this system we calculate the Fermi level E_F taken at half of the energy band gap (see Fig. 1). (ii) Secondly, we consider H and X passivated relaxed surface and we calculate electronic properties and the vacuum energy. The vacuum level is determined by calculating the planar average of the electrostatic potential in the slab supercell (dashed lines in Fig. 1), far away from the surface, from which the macroscopic average along the z direction is then deduced (solid lines in Fig. 1). (iii) As third step, we superimpose the oscillating planar average of the electrostatic potential calculated for the bulk Si supercell with the one calculated in the middle region of the passivated slab. We then calculate the WF as the difference between the vacuum energy level and the E_F of the bulk Si supercell [56]. Results of the application of this method are shown in Fig. 1.

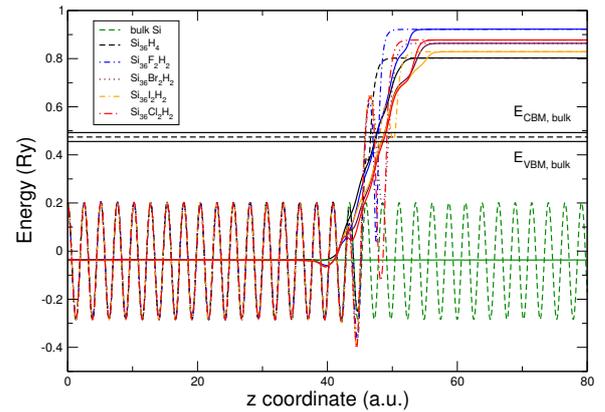


Figure 1 (Color online) Planar (dashed lines) and macroscopic (solid lines) averages of the electrostatic potential for the different chemisorbed slabs. All the curves are superimposed in the bulk-like region of the slab to the oscillating planar average of the electrostatic potential calculated for the bulk Si supercell. The position in energy of the valence band maximum and conduction band minimum for the bulk Si supercell are highlighted. The Fermi level is taken at half of the bulk Si supercell energy gap. Energies are in Rydbergs.

Table 1 summarizes our calculated values of the WF for the H, I, Br, Cl and F terminated Si(100) surface. This table also contains ΔWF , i.e. the differences between the calculated WF values for all the considered chemisorbed surfaces with respect to the hydrogenated ones. Results of Table 1, also reported in Fig. 2 to improve the visibility of trends, point out that the calculated value for the hydrogen terminated surface is not so far from the value of the WF calculated by Sgiarovello et al. [26] for the (2x1) recon-

Table 1 DFT-LDA calculated values (in eV) of the WF for chemically modified Si(100) surface. The results (second column) are shown as function of X , the different atomic termination (first column). The third column reports the WF changes, ΔWF (in eV), induced by the chemisorbed species with respect to the H-terminated surfaces.

X	WF	ΔWF
H	4.47	0
I	4.83	+0.36
Br	5.30	+0.83
Cl	5.49	+1.02
F	6.10	+1.63

structed Si(100) surface (5.12 eV), i.e. without hydrogen passivation. The agreement is noticeable if one takes into account the experimentally observed lowering (about 0.4 eV) of the WF values of clean surfaces due to hydrogen adsorption [57,58].

Fig. 2 clearly indicates a similar trend for the different chemisorbed Si surfaces. All the considered halogen passivated surfaces, indeed, show an increment of WF with respect to the H-terminated one.

The comparison between the calculated values of WF and the experimental results is not trivial due to several reasons. Firstly, the methodologies adopted to measure the WF can sometimes lead to different results. It is worth mentioning, for instance, that the WF can be measured both by using photoelectron spectroscopy and Kelvin probe technique [59]. The first method allows the measurement of the absolute WF , whereas the second gives only a contact potential difference between the probe and the sample surface. However, using a calibration procedure with photoelectron spectroscopy, it is possible to transform the Kelvin probe results into absolute values, thus permitting a more uniform interpretation of the experimental data. Secondly, experimental values of the WF are often influenced by the characteristic of the sample, like the presence of defects, doping and impurities concentration, and so on. These parameters affect the experimental results but cannot be easily included in a numerical ab-initio model. Finally, the intrinsic approximations present in the methodology adopted for the calculation of WF (in our case the DFT-LDA approximation) can lead to an over-estimation or to a under-estimation of the experimental data.

Experimental results for the WF of the Si(100)2x1 clean surface range from 4.6 eV to 4.91 eV [60–62]. These values are larger (with differences from 0.13 eV to 0.44 eV) than our 4.47 eV calculated for the H-terminated Si(100)1x1 surface. We have to consider, however, the experimentally measured decrease of the WF on going from the clean Si(100)2x1 to the H-Si(100)2x1 (from 0.34 eV to 0.40 eV) [57,58]. Additionally, He et al. [28], by employing photoelectron spectroscopy, reported WF val-

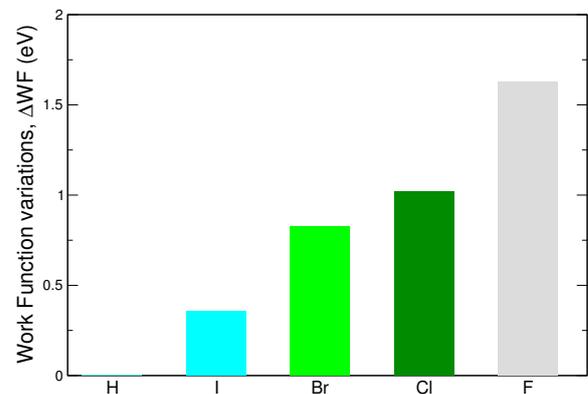
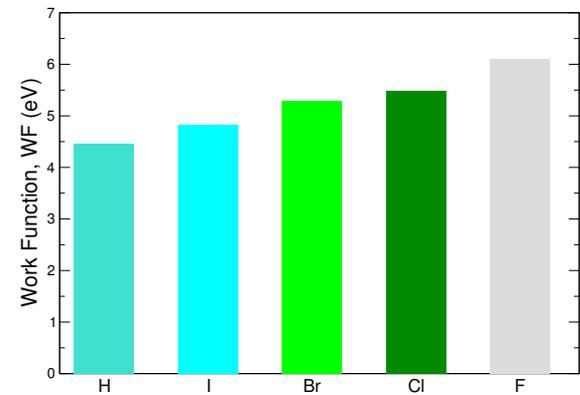


Figure 2 (Color online) Top panel: Trends in the WF values at DFT-LDA level (see text) for the Si(100) surfaces terminated by different atomic species. Bottom panel: Variations in the WF values at DFT-LDA level (see text) for the Si(100) surfaces passivated with different atomic species with respect to the hydrogenated one.

ues of H-passivated Si(100) surfaces (for different doping concentration) in the range between 4.21 eV to 4.55 eV, values that go along with our result of 4.47 eV.

It also possible to compare, in some cases, our results with other theoretical calculations. Anagaw et al. [29] have investigated, using DFT, the change in the WF induced by chemisorption on Si(100)2x1 surfaces. They found 4.6 eV for the H-terminated surface that is in good agreement with our results for the Si(100)1x1 surface (4.47 eV). Ng et al. [31] have investigated the modulation of the WF of silicon nanowires (NWs) grown in the (110) direction, whose dangling bonds at the interface are firstly covered by hy-

drogens that are then substituted with OH or F. For NWs with diameter of about 1 nm they found WF values of 4.61 eV for the fully H-covered system and 5.31 eV - 6.73 eV for the F-terminated case, depending on the percentage of coverage. This behaviour, that consists of an increase of the WF for the F-covered NWs, is congruent with what we have found for F terminated surfaces.

In Table 2 we correlate our calculated WF with theoretical and experimental results obtained for the Si(111)

Table 2 WF values calculated for the differently terminated Si(100) surfaces. The value in parenthesis are related to the WF changes, ΔWF , induced by the different chemisorbed species with respect to the H-terminated surfaces. Our results for the Si(100) surfaces are compared with theoretical [36] and experimental results [35,63–67] obtained for the Si(111) surface. All the values are in eV.

X	Our	Th.(50%) [36]	Th.(100%) [36]
H	4.47 (0.0)	4.35 (0.0)	4.35 (0.0)
I	4.83 (+0.36)	4.72 (+0.37)	4.62 (+0.27)
Br	5.30 (+0.83)	4.98 (+0.63)	5.12 (+0.77)
Cl	5.49 (+1.02)	5.09 (+0.74)	5.43 (+1.12)
F	6.10 (+1.63)	5.36 (+1.01)	6.39 (+2.04)

X	Exp. [35]	Exp. [63]	Exp. [64–67]
H	4.16-4.24 (0.0)	4.17-4.35 (0.0)	4.42(0.0)
I			
Br	4.32 (+0.16)		
Cl	4.60 (+0.36)		(+1.12-1.15)
F			

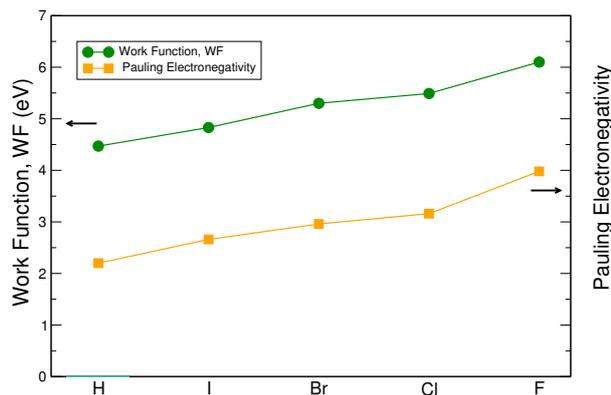


Figure 3 (Color online) Calculated WF values (in eV) for Si(100) surfaces with different atomic termination (green dots). Orange squares represent the Pauling electronegativity of the chemisorbed atoms. The lines are guide for the eyes.

surfaces [35,36,63–67]. Noticeably, in Ref. [36], the WF has been calculated for two different coverages of the Si(111) surface (50% and 100%): the first one is the same of our Si(100) surface, where only one of the two capping H atoms has been substituted by other species. We see that the agreement is indeed quite good. Moreover Li et al. [35] have calculated at G_0W_0 level the variation in the ionization potential with respect to a H-terminated Si(111) surface due to the substitution of H with Br and Cl. They found variation of 0.8 and 1.1 eV to be compared to our 0.83 and 1.02 eV values, respectively. Since the variations in the ionization potential cannot be very different from those of the WF , these results confirm the accuracy of LDA in evaluating the WF for the considered systems. The agreement between our calculated WF changes for the Si(100) surfaces with those determined for Si(111) surfaces is quite remarkable: we can observe a very similar increase of the WF in all the halogen passivated surfaces. These results are a strong indication that, for this quantity, the major role is played by the chemisorbed species present on the surfaces and not by the different surface orientation.

As previously stated, the dependence of WF values on the passivating species for the Si surface (see Fig. 3) shows that halogen atom chemisorption (I, Br, Cl and F) always induces an increase in the WF . With the purpose of finding a rationale for the observed trend, in Fig. 3 we also reported the Pauling electronegativity values for each halogen atom. Looking at the figure, it is clear that the augmentation of the WF value for the passivated Si surface follows the same trend for the increase of the electronegativity going from I to F.

Table 3 $\Delta_q(Si)$ indicates the calculated Löwdin charge change for the Si atom in contact with the halogen element, while q_x is the calculated partial charge of halogen atom. All the values are in units of electron charge. Theoretical results for the Si(111) surface [36] are also reported.

X	$\Delta_q(Si)$	q_x	$\Delta_q(Si)$ [36]	q_x [36]
I	-0.04	-7.02	-0.03	-7.07
Br	-0.17	-7.12	-0.13	-7.16
Cl	-0.25	-7.18	-0.22	-7.22
F	-0.53	-7.36	-0.53	-7.46

In Table 3 we report the Löwdin charge changes $\Delta_q(Si)$ calculated for the Si atom in contact with the halogen element with respect to the fully hydrogenated configuration together with q_x the computed partial charge of the halogen atoms. The data, in agreement with results of Ref. [36] for the (111) Si surface, indicates an increment of $\Delta_q(Si)$ when we move from I to F meaning that the charge transfer is more pronounced when the halogen electronegativity increases, reflected also in the augmentation of q_x . These results point out the direct relation that exists between ΔWF and the redistribution of charge at the surface. Furthermore, this confirms again the dominant

role of the chemisorbed species with respect to the surface orientation in the determination of the WF .

4 Conclusions We presented and discussed first principles simulations of halogen passivated Si(100) surface. By calculating work function and charge transfer we were able to draw up a general trend for the halogen absorption effect on the Si electronic properties. In particular, we have found that halogens have a tendency to increase the WF of Si surfaces as a consequence of their high electronegativity that causes a large charge transfer with an associated dipole moment. The behavior observed moving from I to F fairly follows the increase in the corresponding Pauling electronegativity. The generality of this effect is confirmed by the discussion and comparison of our results with other calculations and experiments on both Si(100) and Si(111). This sheds light on the major role of the type of passivation with respect to the surface crystal orientation.

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References

- [1] R. M. Swanson, Prog. Photovoltaics Res. Appl. **14**, 443 (2006).
- [2] L. Pavesi and D. Lockwood, Silicon Fundamentals for Photonics Applications (Silicon photonics, 2004).
- [3] M. Bernardi, D. Vigil-Fowler, J. Lischner, J. B. Neaton, and S. G. Louie, Phys. Rev. Lett. **112**, 257402 (2014).
- [4] E. Degoli, R. Guerra, F. Iori, R. Magri, I. Marri, O. Pulci, O. Bisi, and S. Ossicini, C. R. Phys. **10**, 575 (2009).
- [5] M. Govoni, I. Marri, and S. Ossicini, Phys. Rev. B **84**, 075215 (2011).
- [6] F. Iori, E. Degoli, R. Magri, I. Marri, G. Cantele, D. Ninno, F. Trani, O. Pulci, and S. Ossicini, Phys. Rev. B **76**, 085302 (2007).
- [7] R. Guerra, I. Marri, R. Magri, L. Martin-Samos, O. Pulci, E. Degoli, and S. Ossicini, Phys. Rev. B **79**, 155320 (2009).
- [8] M. Cazzanelli, F. Bianco, E. Borga, G. Pucker, M. Ghulinyan, E. Degoli, E. Luppi, V. Véniard, S. Ossicini, D. Modotto, S. Wabnitz, R. Pierobon, and L. Pavesi, Nat. Mater. **11**, 148154 (2012).
- [9] M. Govoni, I. Marri, and S. Ossicini, Nat. Photon. **6**, 672 (2012).
- [10] I. Marri, M. Govoni, and S. Ossicini, J. Am. Chem. Soc. **136**, 13257 (2014).
- [11] I. Marri, M. Govoni, and S. Ossicini, Beilstein J. Nanotechnol. **6**, 343–352 (2015).
- [12] I. Marri, M. Govoni, and S. Ossicini, Solar Energy Material & Solar Cells **145**, 162 (2016).
- [13] J. Dabrowski and H. J. Müssig, Silicon Surfaces and Formation of Interfaces (World Scientific, 2000).
- [14] J. M. Buriak, Chem. Rev. **102**, 1271 (2002).
- [15] M. Amato and R. Rurali, Prog. Surf. Sci. **91**, 1 (2016).
- [16] S. P. Cummings, J. Savchenko, and T. Ren, Coord. Chem. Rev. **255**, 1587 (2011).
- [17] A. V. Teplyakov and S. F. Bent, J. Vac. Sci. Technol. A **31**, 050810 (2013).
- [18] K. T. Wong and N. S. Lewis, Acc. Chem. Rev. **47**, 3037 (2014).
- [19] M. Amato, M. Palumbo, R. Rurali, and S. Ossicini, Chem. Rev. **114**, 1371 (2014).
- [20] W. Peng, S. M. Rupich, N. Shafiq, Y. N. Gartstein, A. V. Malko, and C. Y. J., Chem. Rev. **115**, 12764 (2015).
- [21] B. Fabre, Chem. Rev. **116**, 4808 (2016).
- [22] Y. Morikawa, Phys. Rev. B **51**, 14802 (1995).
- [23] E. Sicilia, M. Toscano, T. Mineva, and N. Russo, Int. J. Quant. Chem. **61**, 571 (1997).
- [24] T. M. Bhave and S. V. Bhoraskar, J. Vac. Sci. Technol. B **16**(4), 2073 (1998).
- [25] S. Y. Davydov and A. V. Pavlyk, Tech. Phys. **49**, 1050 (2004).
- [26] C. Sgiarovello, N. Binggeli, and Baldereschi, Phys. Rev. B **64**, 195305 (2001).
- [27] S. Kajita, T. Nakayama, and J. Yamauchi, J. Phys.: Conf. Ser. **29**, 120–123 (2006).
- [28] T. He, H. Ding, N. Peor, M. Lu, D. A. Corley, B. Chen, Y. Ofir, Y. Gao, S. Yitzhai, and J. M. Tour, J. Am. Chem. Soc. **130**, 1699 (2008).
- [29] A. Y. Anagaw, R. A. Wolkow, and G. A. DiLabio, J. Phys. Chem. C **112**, 3780 (2008).
- [30] S. N. Novikov and S. P. Timoshenkov, Russ. J. Phys. Ch. A **84**, 1266 (2010).
- [31] M. F. Ng, L. Y. Sim, D. X. H. Jin, K. H. Lim, and S. W. Yang, Theor. Chem. Acc. **127**, 689 (2010).
- [32] A. J. Cooper, K. Keyvanfar, A. Deberardinis, L. Pu, and J. C. Bean, Appl. Surf. Sci. **257**, 6138 (2011).
- [33] C. H. Kuo, C. P. Liu, S. H. Lee, H. Y. Chang, W. C. Lin, Y. W. You, L. H. Y., and S. J. J., Phys. Chem. Chem. Phys. **13**, 15122 (2011).
- [34] H. S. Huang, W. Y. Chan, W. B. Su, G. Hoffman, and C. S. Chang, J. Appl. Phys. **114**, 214308 (2013).
- [35] Y. Li, L. E. OLeary, N. S. Lewis, and G. Galli, J. Phys. Chem. C **117**, 5188 (2013).
- [36] H. H. Arefi and G. Fagas, J. Phys. Chem. C **118**(26), 14346 (2014).
- [37] H. H. Arefi, M. Nolan, and G. Fagas, Phys. Chem. Chem. Phys. **18**(19), 12952 (2016).
- [38] S. M. Nawaz, S. Dutta, and A. Mallik, IEEE T. Electron Dev. **62**, 3951 (2015).
- [39] A. Papageorgopoulos and M. Kamaratos, Surf. Sci. **466**, 173 (2000).
- [40] X. Zhang, E. Garfunkel, Y. J. Chabal, S. B. Christman, and E. E. Chaban, Appl. Phys. Lett. **79**, 4051 (2001).
- [41] B. Delley and E. F. Steigmeier, Appl. Phys. Lett. **67**, 2370 (1995).
- [42] A. N. Kholod, S. Ossicini, V. N. Borisenko, and F. Arnaud D'Avitaya, Surf. Sci. **527**, 30 (2003).
- [43] Y. Li and G. Galli, Phys. Rev. B **82**, 045321 (2010).

- [44] M. Amato, M. Bertocchi, and S. Ossicini, *J. Appl. Phys.* **119**, 085705 (2016).
- [45] This allows us to compare the results of this work with previous calculations related to SiGe slabs [44], where a vacuum of about 20 Å was used. Notably, a difference in the vacuum region of nearly 15 Å induces variations in the WF values of the order of 0.04 eV.
- [46] P. Giannozzi, S. Baroni, N. Bonini, M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, G.L. Chiarotti, M. Cococcioni, I. Dabo, A.D. Corso, S. de Gironcoli, S. Fabris, G. Fratesi, R. Gebauer, U. Gerstmann, C. Gougousis, A. Kokalj, M. Lazzeri, L. Martin-Samos, N. Marzari, F. Mauri, R. Mazzarello, S. Paolini, A. Pasquarello, L. Paulatto, C. Sbraccia, S. Scandolo, G. Scaluzero, A.P. Seitsonen, A. Smogunov, P. Umari, and R.M. Wentzcovitch, *J. Phys.: Condens. Matter* **21**, 395502 (2009).
- [47] I. Marri and S. Ossicini, *Solid State Commun.* **147**, 205 (2008).
- [48] A. Iacomino, G. Cantele, D. Ninno, I. Marri, and S. Ossicini, *Phys. Rev. B* **78**, 075405 (2008).
- [49] S. Ossicini, O. Bisi, E. Degoli, I. Marri, F. Iori, E. Luppi, R. Magri, R. Poli, G. Cantele, D. Ninno, F. Trani, M. Marsili, O. Pulci, V. Olevano, M. Gatti, K. Gaal-Nagy, A. Incze, and G. Onida, *J. Nanosci. Nanotechn.* **8**, 479 (2008).
- [50] R. Guerra, I. Marri, R. Magri, L. Martin-Samos, O. Pulci, E. Degoli, and S. Ossicini, *Superlattice and Microstructure* **46**, 246 (2009).
- [51] F. Iori, E. Degoli, E. Luppi, R. Magri, I. Marri, G. Cantele, D. Ninno, F. Trani, and S. Ossicini, *J. Lumin.* **121**, 335 (2006).
- [52] L. Hedin, *Phys. Rev.* **139**, A796 (1965).
- [53] G. Onida, L. Reining, and A. Rubio, *Rev. Mod. Phys.* **74**(2), 601 (2002).
- [54] N.E. Singh-Miller and N. Marzari, *Phys. Rev. B* **80**, 235407 (2009).
- [55] P.C. Rusu, G. Giovanetti, C. Weijtens, R. Coehoorn, and G. Brocks, *J. Phys. Chem. C* **113**, 9974 (2009).
- [56] I. Borriello, G. Cantele, D. Ninno, G. Iadonisi, M. Cossi, and V. Barone, *Phys. Rev. B* **76**, 035430 (2007).
- [57] K. Fukuiwara, *Phys. Rev. B* **26**, 2036 (1982).
- [58] A.E. Souzis, M. Sidl, W.E. Carr, and H. Huang, *J. Vac. Sci. Tech. A* **7**, 720 (1989).
- [59] S. Pouch, M. Amato, M. Bertocchi, S. Ossicini, N. Chevalier, T. Mélin, J.M. Hartmann, O. Renault, V. Delaye, D. Mariolle, and L. Borovik, *J. Phys. Chem. C* **119**, 26776 (2015).
- [60] S. Kono, Y. Enta, T. Abukawa, and T. Kinoshita, *Appl. Surf. Sci.* **41**, 75 (1989).
- [61] J. Günster, T. Mayer, and V. Kempter, *Surf. Sci.* **359**, 155 (1996).
- [62] Q.B. Lu, R. Souda, D.J. O'Connor, and B.V. King, *Phys. Rev. B* **54**, R17347 (1996).
- [63] C.A. Hacker, *Solid State Electron.* **54**, 1657 (2010).
- [64] R. Hunger, C. Pettenkofer, and R. Scheer, *J. Appl. Phys.* **91**, 6560 (2002).
- [65] R. Hunger, R. Fritsche, B. Jaeckel, W. Jaegermann, L. Webb, and N. Lewis, *Phys. Rev. B* **72**, 045317 (2005).
- [66] R. Hunger, W. Jaegermann, A. Merson, Y. Shapira, C. Pettenkofer, and J. Rappich, *J. Phys. Chem. B* **110**, 15432 (2006).
- [67] G. Lopinski, B. Eves, O. Hul'ko, C. Mark, S. Patitsas, R. Boukherroub, and T. Ward, *Phys. Rev. B* **71**, 125308 (2005).