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1 *Perspective article*

2 **Molecular Spins in the context of Quantum** 3 **Technologies.**

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13 **Abstract:** Molecular spins have shown interesting quantum features which make
14 them potential candidates for the implementation of quantum information processing.
15 New challenges related to possible applications in broader class of quantum technologies
16 are currently under discussion. Here, we revisit some key features trying to learn
17 something from experiences in near fields.

18 **Keywords:** Molecular NanoMagnets; Quantum Technologies.

19 **PACS:** J0101

20

21 **1. Introduction**

22 Quantum Computing (QC) is one of the Quantum Technologies (QT) [1] that aims at
23 exploiting genuine quantum features of systems and devices. QC was envisaged in early Eighties
24 by Richard Feynman and other pioneers and since then other possible applications of quantum
25 systems became feasible also thanks to technological progress. Among QTs, quantum-
26 communications, sensing, cryptography and metrology are now attracting much interest. Besides
27 the discreteness of energy levels, the possibility to create and maintain superposition of states and
28 quantum correlation (entanglement) are considered two of the main features of quantum systems
29 that do not exist in the classical world. Quantum systems such as photons, cold atoms, spin
30 impurities in solids, semiconducting and superconducting devices have been intensively studied in
31 the last decades and several applications based on these systems are in rapid development and
32 some of these already appeared in the market.

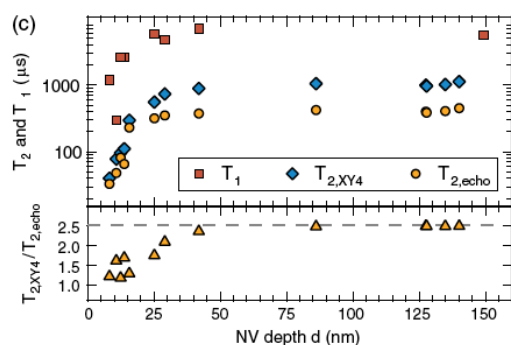
33 Different spins centers in solid state have been -and they currently are- intensively studied for
34 potential applications: spin impurities in Si have been studied and proposed for qubit encoding [2];
35 Nitrogen Vacancies (NV) centers, as well as other colour centers in diamond or SiC, are now
36 attracting much interest for their potentialities as atomic scale sensors, besides the possibility to use
37 them as hardware for QC [3]. In spite of different preparation technologies, experiments (eg.
38 manipulation and read out of spin) and modeling (eg. spin dynamics, source of decoherence, spin
39 entanglement etc.) all these spins systems obviously share many analogies, thus we can probably
40 learn something by comparing the properties of molecular spins to those of other spin systems. This
41 is intended to be the main stream of this perspective article. We shall focus on three main trends in
42 the field: coherent manipulation of spin ensembles; molecular quantum spintronics and perspective
43 to embed (molecular) spins in quantum circuits.

2. Coherent spin dynamics of spin ensembles.

Quantum effects in molecular spin systems have been explored in the last two decades. The first milestone was the discovery of quantum tunneling of the magnetization in Nineties that demonstrated that molecular spins are real systems on which quantum effects can be controlled and studied in laboratory conditions [4]. Next challenge was to control the dynamics of molecular spins. Molecules with low spin ground state (ideally $S=1/2$) well isolated by excited states are prototypical cases for this type of study [5] and long-lived coherent oscillations between two spin states is target experiment that can be performed by pulsed Electron Spin Resonance (ESR) techniques. The understanding and control of decoherence mechanisms are of fundamental importance for the dynamics of spins in solid state environment. Thus, in the last decade there has been an intense activity in searching molecular spin systems with improved performances, more specifically with the longer relaxation (T_1) and dephasing (T_2) times, as discussed in previous reviews [6, 7] and themed issues of specialized Journals and book series [8, 9, 10]. Inter-molecular interactions are, in general, detrimental, thus working with spin ensembles diluted in non magnetic matrix is mandatory. In this way, coherent spin oscillations were soon observed in low-spin Cr_7Ni rings [11] and V_{15} [12] at liquid He temperature.

Most of the solid state hardware for quantum technology works at cryogenic temperatures yet this limits all of them to niche applications and increasing the working temperature –and reducing the use of magnetic field- is certainly appealing in view of widely spread applications. Working at 2K instead of 20mK is a gain of two orders of magnitude in temperature and this allows to avoid the use of dilution refrigerators. Working at 20K will allow a gain of another order of magnitude and, for instance, the use of closed cycle refrigerators that can be developed even on one chip. At 80K we may use liquid nitrogen and room temperature operation will open to much broader class of applications. A first breakthrough along this line was obtained with the report of persistence of long T_2 up to 100K in commercial CuPc derivative [13]. This result was followed by rational design of mononuclear derivatives for which T_2 was further improved [14, 15, 16]. The relevant figure of merit here is the ratio between the manipulation time and the spin coherence time. Since for typical pulsed ESR set up the time to manipulate an electron spin is about 10ns, both T_1 and T_2 need to be much longer in order to observe coherent spin oscillations and, possibly to perform some error corrections. As a matter of facts, Rabi oscillations at room temperature have been recently reported in VOPc [17]. These results represent an important milestone for molecular magnetism.

In ref.13 one can find a direct comparison of relevant coherence time measured at different temperatures in different spin systems. Yet, care should be taken in making direct comparison since data should be taken in similar conditions and experiments. One relevant point is the dependence of decoherence times of spin defects on the depth of impurity in the solid. For instance, it is well documented that T_1 and T_2 decrease for shallow impurities in both Si and diamond and they dramatically drop when spin is closer than 5nm to the surface. For instance, fig.1 (from ref.18) shows that both T_1 and T_2 drop below 100 μs for NV positioned less than 10nm from the diamond surface. Interestingly, the coordination of the spin center is not altered but the shallow defects were demonstrated to be sensitive to both the magnetic and the electric [19] noise on the diamond surface. This should draw our attention to control the environment of our spin: if impurities are diluted in nuclear free environment, magnetic noise can be drastically reduced. Same attention should be taken to reduce (electrical) noise produced by local vibrations. From this point view, the seminal work reported in ref.20 demonstrated that molecular engineering can be a powerful tool to control the closest environment of the spin. On the other hand, if we need to expose spins to an external environment, such as a biological system or an electronic circuit, the ligand shell may, in some ways, protect or, at least, define the closest surrounding around the sensitive spin.



94
 95 figure 1 Coherence time T_2 and relaxation time T_1 as a function of the NV depth in diamond. The plot shows
 96 strong suppression of coherence for shallow NV centres. The lower panel shows $T_{2,XY4}=T_{2,echo}$ ratio that is
 97 reduced as well with decreased depth, also indicating that dynamical decoupling with $N=4$ pulses is less
 98 efficient for shallower NVs. (Reproduced from ref.18 with permission from American Physical Society)
 99

100 In view of possible applications for which spins need to be embedded in external environment,
 101 the next challenge seems to be the consolidation of performances with an overall (chemical,
 102 structural) robustness of the molecules in different working conditions. To this end, tests on
 103 isolated molecules on surfaces or under different stimuli or thermal cycling need to be performed.
 104 Important achievements have done by studying isolated molecules on surfaces by different
 105 techniques such as STM or X-ray spectroscopies. Not all, but few molecules resulted to be robust
 106 enough to substantially maintain their (static) magnetic features when dispersed. However subtle
 107 effects may occur when the environment change and these need to be carefully checked case by
 108 case [21]. We expect that different molecules can be designed for specific ambient/applications,
 109 thus, for instance, molecules with external organic ligand and specific linkers will be more suitable
 110 for biological applications; molecules that can be sublimed and with linkers designed to graphitic or
 111 Si surfaces will better work embedded in electronic circuits; while molecules with robust periphery
 112 (eg. oxygen ligand) can be possible applied in ambient (air) conditions.
 113

114 Quantum computation may require the implementation of quantum gates [22]. A Universal
 115 scheme for computation that, in principle, solves a broad class of computational problems, can be
 116 reduced to basic gates using one or two qubits. For spins, a one-qubit gate corresponds to rotation
 117 along two orthogonal directions and therefore relevant tests are observation of Rabi oscillations in
 118 pulsed ESR experiments, as previously discussed. Several experiments on mononuclear molecular
 119 spin centers have been performed as also reviewed in one article of this issue [23].

120 Next challenge is to perform conditioned quantum operations involving more than one spin
 121 center. In general terms, we need to demonstrate that the dynamics of one (*target*) spin is
 122 conditioned by the state of the second (*control*) spin [22]. This may require the two spins to be
 123 distinguishable either spatially or spectroscopically and a typical sequence for two-qubit gate
 124 encoding encompasses: the initialization of the system, the rotation of the *target* spin under different
 125 conditions of the *control* spin qubit and finally the read out of the system state. According to the
 126 specific pulse sequence and rotation, a number of two-qubit gates can be performed, similarly
 127 classical two-bit gates, such as controlled-NOT. This conditional spin dynamics is possible by
 128 exploiting quantum correlation (entanglement) of two or more spin centers and, since this is a
 129 genuine quantum property, quantum gates essentially different from classical gates.

130 Spin correlation (entanglement) is intrinsic property of system states that needs to be properly
 131 quantified by using suitable experimental quantities or mathematical functions [24]. Spin
 132 entanglement can be obtained within a single molecule or at supramolecular level (i.e. between
 133 molecules) by controlling spin topology and magnetic interactions. Along this line, dimers of
 134 molecular spins with weak permanent magnetic coupling could be designed and synthesized [25]
 135 and sizable spin entanglement was demonstrated [26]. One issue of this approach is related to the
 136 fact that, during the gate operation, the coupling between the two qubits needs to be switched on

137 and off in order to allow independent rotation of the two spin qubits. In other words, switchable
138 links would be required. This problem can be solved by engineering the molecular states in such a
139 way to make use of auxiliary states [27] as also described in the realistic case of antiferromagnetic
140 rings [28]. An alternative way to entangle two spins is to simply position two spin centers at fixed
141 distance (few *nm*) and exploit dipolar (i.e. *through space*) interaction as done for spin defects in
142 diamond [29] and Si. The drawback of using dipolar interaction is, however, its persistent character
143 and its strong dependence on both distance and orientation of the two spin centers. These make
144 such an approach not simply scalable for defects. Maybe this problem is less critical for molecular
145 assemblies. A further alternative way to entangle two spin centers far apart is, however, the use of
146 flying qubits (photons) as demonstrated in a recent experiment in which entanglement between two
147 NV centers distant 25nm from one another was achieved by engineering the pulse sequence and
148 using dynamic decoupling technique [30].

149 Proposals for the implementation of two-qubit gates with specific molecular spins systems
150 have appeared and preliminary experiments have been reported using binuclear lanthanides [31],
151 radicals [32], antiferromagnetic rings [33]. Moreover, an intense search of suitable bi-nuclear
152 molecular system is currently under way and many potential candidates are ready to be tested. In
153 terms of experiments, this would require the use of multi-frequency spectrometers that allow to
154 distinguish the two qubits or to activate switchable interaction. Commercial instrumentation, such
155 as that for Pulsed Electron-Electron Double Resonance (PELDOR), is well developed but, in most of
156 the cases, it needs be adapted to working conditions (frequency, temperature, magnetic field) of
157 molecular qubit. At present, dedicated instrumentation is accessible only to few laboratories
158 worldwide whilst the implementation of two-qubit gates with molecular spin ensembles would
159 require more flexibility/tunability on both molecular qubits and instrumentation. Moreover, in spite
160 of the progress achieved in using pulse ESR techniques to molecular electron spins [34], much work
161 still remains to be done to encompass inhomogeneities of both molecular features and applied field
162 and, more in general, to efficiently decouple spin to the environment. For instance, very interesting
163 results have been obtained in the optimization of pulse sequences for dynamical decoupling [35]. In
164 view of using multi-frequency pulse spectrometer, it is worth to remind that magnetic molecules
165 generally possess both electron and nuclear spins and probably the best approach will be to
166 combine the long coherence time of nuclear spins with the faster manipulation and read out of
167 electronic spins. For this ENDOR technique offers several interesting solutions as discussed in
168 ref.36.

169 Final general remark should be done concerning the use of spin ensembles. The main
170 criticism, well known in the QC community for NMR, is that working with ensembles makes the
171 systems not –easily- scalable. In other words, the resources, (eg spins) required for solving
172 problems and for the relative error correction, scale fast with increasing number of data, thus
173 making this type of architectures not suitable for a universal use. Within a more pragmatic
174 perspective, we may focus our interest on quantum simulators [37], i.e. small quantum computers
175 that solve only specific problems which would require huge amount of classical computing
176 resources, or even intractable by classical computers. In practice this requires the use of molecular
177 derivatives comprising small but well defined spin clusters whose dynamics may solve a complex
178 problem or emulate the behavior of system of interest in other fields (for instance a chemical
179 reaction or a problem in solid state physics, see also review article in this issue [38]). This topic is
180 still largely unexplored and deserves joint effort from chemists and physicists.

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183 3. Molecular Quantum Spintronics

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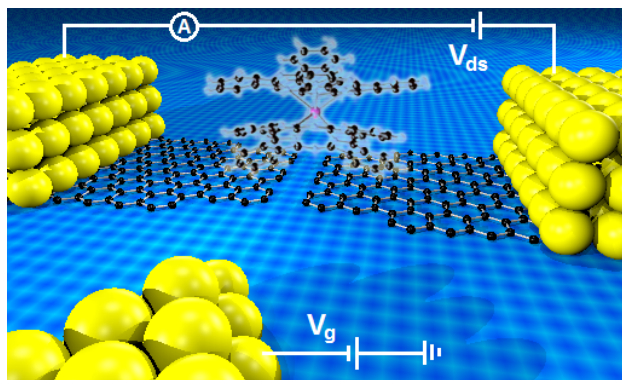
185 The addressing, manipulation and read out of **single molecular objects** constitute the next big
186 conceptual and technological challenge. If we can avoid some of the drawbacks found in using spin
187 ensembles, here we have to find efficient ways to detect tiny magnetic signals and to individually
188 manipulate spins by protecting fragile quantum states from the environment at the same time.

189 Different approaches have been tested to detect tiny magnetic signals: one milestone was the
190 development of **nano-SQUID** made with carbon nanotube [39]. Despite the magnetic flux
191 sensitivity of these devices can achieve the quantum limit, the main limitation of this type of
192 magnetometers is constituted by the magnetic coupling - through space- of the molecule with the
193 sensor. Direct single molecule detection by charge current seems more appealing, at least for
194 spintronic devices, and it may benefit from the progresses achieved in close fields such as molecular
195 electronics, scanning microscopy and single-electron semiconducting devices.

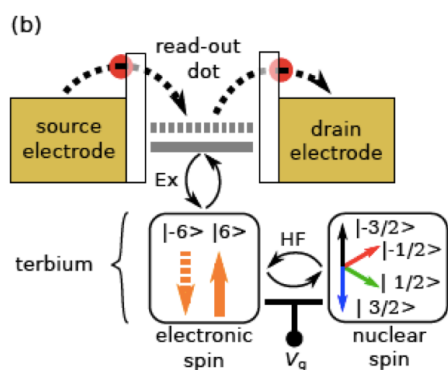
196 The use of **scanning probes**, more specifically tunneling tips, holds much promise for their
197 potentialities to localize, move and read out single magnetic atoms and molecules [40]. With respect
198 to bare magnetic atoms, the presence of an organic ligand seems to delocalize the magnetic features
199 of the molecule [41] and data interpretation is still under debate. Recently interesting experiments
200 on the manipulation and read out of molecular spin (TbPc₂) by radio frequency [42] and single
201 magnetic atom on a surface by pulse-MW sequences have been reported [43].

202 At the same time, the development of **tunnel junctions** that may host a single molecule (fig.2)
203 have been independently reported by different groups. Magnetic features of individual molecules
204 have been observed and reported for Fe₄ [44, 45], TbPc₂ [46] organic radicals [47]. An open question
205 here is how the charge current from the leads perturbs the magnetic state of the molecule. Changes
206 of valence (and spin) state of the magnetic core are indeed expected. Since each type of molecule
207 behaves in different way, this issue needs to be evaluated case by case. An alternative read out
208 scheme comprises a quantum dot whose conductivity is affected by the spin state of the magnetic
209 center that is in close proximity and coupled with the device. This scheme is analogous to the spin
210 dependent tunneling that is used to read out spin qubits also in semiconductors, but here the
211 **quantum dot** can be made by CNT [48], graphene nano-constrictions [49] or even the organic
212 radical present in the ligand of the molecule as demonstrated in the case of TbPc₂ molecular spin
213 transistor [50]. In this scheme, the charge channel and the spin center are two separate bodies,
214 coupled by exchange interaction (see fig.3).

215 The experiments reported by the Grenoble team went a step ahead demonstrating the
216 possibility to read out and manipulate the nuclear spin state. The nuclear spin of Tb ($I=3/2$) is
217 indeed coupled to the electron magnetic moment (ground state $J=6$, $m_J=\pm 6$) by the hyperfine
218 interaction (fig.3) giving rise to hybrid electron-nuclear states whose level anticrossing (LAC) are
219 well visible at low magnetic field. The nuclear states can be identified by sweeping magnetic field
220 and comparing the position of the LACs with those measured on bulk samples. Note that, since
221 measurements are performed on a single qubit, the process must be repeated several times in order
222 to get a significant statistics.
223



224
225 *Figure 2: Molecular spin transistor made by graphene electrodes and TbPc₂ molecule. In this case*
226 *graphene based electrodes are used to contact TbPc₂ molecule. (Reproduced from Ref. 51 with permission*
227 *from the Consiglio Nazionale delle Ricerche (CNR) and The Royal Society of Chemistry.)*
228



229
 230 Figure 3 Scheme of functioning for molecular spin transistor with separate quantum dot (Pc radical) in
 231 which charge can tunnel from electrodes, the electron magnetic moment $J=6$ and nuclear spin ($I=3/2$) of Tb^{3+}
 232 ion (Reproduced from Ref. 54 with permission from The American Association for the Advancement of
 233 Science)

234
 235 By exploiting the long coherence time of nuclear spin (coherence time exceeding $60\mu s$) and
 236 the hyperfine electric Stark effect, the same team has then demonstrated that it is possible to
 237 manipulate the nuclear spin of a single molecule. More specifically nuclear spin trajectories [52] and
 238 Ramsey fringes have been observed [53]. These results compare well with similar experiments
 239 reported for spin impurity (P donor) in Si [54] and witness a tangible contribution of molecular
 240 magnetism to QT. At the same time, we expect that the next steps along this research line will move
 241 in parallel with similar devices based on spin impurities in semiconductors [55] or NV centers in
 242 diamond [56], that is, the implementation of two-qubit gates or multi-qubit algorithms.

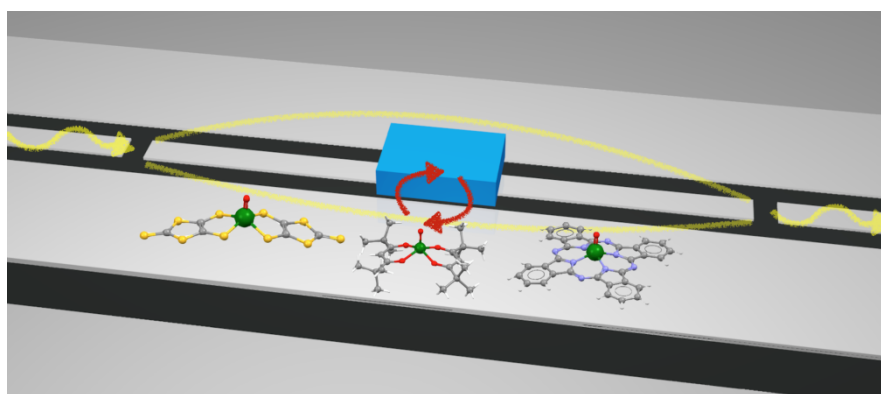
243 Different schemes can be envisaged for the implementation of two-qubit gate with molecular
 244 spintronic devices: a CNT may host two (or more) molecules and several gate electrodes that can
 245 act to switch on and off interaction between molecules. Alternatively, a suitable bi-nuclear molecule
 246 within a junction can be used as molecular hardware. In both cases, spin manipulation can be done
 247 by addressing each spin center by microwaves using suitable pulse sequences.

248 Working devices need to be reliable and suitable **quantum error correction protocols** are to be
 249 developed. The problem of error correction is well known in computer science and some schemes
 250 are discussed in textbooks [22]. Briefly, it may occur that during a quantum operation a qubit
 251 accidentally flips, invalidating the whole process. To mitigate and correct these errors we may encode
 252 the qubit in more than one (typically 3) processor and then use the majority rule: if the probability
 253 of accidental flip is relatively low, 2 qubits over 3 remain in the correct state while one contains
 254 error. In this case the majority (2 over 3) determine the correct state while the third one can be
 255 corrected. Further methods to correct different type of quantum errors have been reported in the
 256 literature for other spin systems, in particular those tested for NV centers [57].

257 The previous discussion leads us to another technological issue: soon or later we'd need to
 258 have arrays of similar devices working at the same time in order to guarantee the scalability of our
 259 computing machine. Although this aspect is less discussed in literature, the rate of success R in the
 260 device fabrication process is relatively low (often $<10\%$). These numbers are typical for tunnel
 261 junctions in molecular electronics although they related to the specific fabrication process and
 262 certainly need -and can- be improved in future by exploiting novel bottom-up fabrication methods.
 263 It should be noted, however, that also the fabrication of CMOS-compatible quantum devices below
 264 10nm by top-down (lithographic) techniques are affected by low yields. In view of applications, it is
 265 worth to fix some benchmarks that may assess the reliability of the fabrication process. As an
 266 example, let us suppose we intend to test a quantum error correction protocol. As mentioned above,
 267 we'd need at least three identical devices and in this case the probability to get all these working at
 268 the same time will drop at R^3 . Thus, we would need to fabricate 10^3 devices to get at least 1 working
 269 machine to test quantum error correction code!

270 Concerning multi-qubit devices, recently the implementation of **Grover's algorithms** within
 271 the ground state multiplet of a single molecule [58] has been carried out at the L. Néel institute

272 laboratory [59]. The Grover's algorithm efficiently solves a specific problem by performing a search
273 of an item within a set of data. The sub-levels of the ground multiplet inequivalently spaced in
274 energy, such as the mixed nuclear-electronic states of TbPc₂, constitute an excellent playground to
275 test this scheme since each of them can be addressed separately by microwave pulses. Note that in
276 this scheme, spin entanglement is not required, so the scheme can be implemented within a single
277 molecule. As a proof of concept, few sub-levels can be considered but it is worth to mention that the
278 implementation is hardly scalable by using a single molecule. It is nevertheless worth to highlight
279 here that, so far, the Grover's algorithm has been implemented with trapped ions, photons and
280 superconducting qubits but, for the best of our knowledge, not with other solid state platform such
281 as impurity in Si: this, once more, testifies the pioneering role of molecular spin in the field of
282 Quantum Technologies.
283



284 Figure 4. Artistic view of molecular spin coupled to planar superconducting resonators.
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288 4. Molecular spins in hybrid quantum architectures. 289

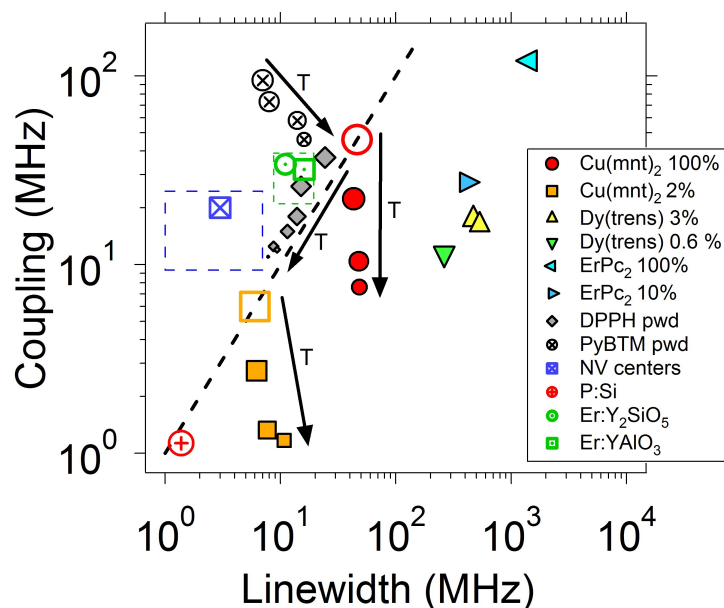
290 The ability to manipulate and read out an arbitrary spin state in a molecule is certainly
291 pre-requisite to be good candidates for quantum information processing. Yet, other features are
292 required in view of wider exploitation of molecular spins for quantum technologies. One of this is
293 the possibility to exchange quantum information between solid state registers (spins) and flying
294 qubits (photons). Molecules offer a broad spectrum of frequencies for an efficient coupling with
295 photons: nuclear spins are active at radio frequencies (MHz) while the pattern of magnetic energy
296 levels of electron spins fully covers the microwave (MW) range (GHz). Several molecules are also
297 active in the visible range and, interestingly, their response can be sensitive to their spin state.
298 Again, the main challenge is to *coherently* couple the spin with photons, that is the match in energy
299 should occur along with the transfer of phase information. This implies that the spin manipulation
300 should be fast enough to overcome the decoherence mechanisms of both the spin and photon
301 systems. Key experiments in this context make use of microwave resonant cavities with high
302 quality factor in which long-lived photons trapped in the cavity couple with spins (fig.4). First tests
303 are typically performed in continuous wave operation mode, although the final goal is to get hybrid
304 devices in which quantum information are exchanged through MW pulse sequences. Coherent
305 spin-photon states are obtained in the so-called *strong coupling regime* for which the spin-photon
306 coupling is stronger than each decoherence rate of spins and photons [7]. Again, only molecular
307 spins with the longest coherence time are suitable to pass this test. Spin can couple with the
308 magnetic component of radiation B_1 through their dipole $g\mu_S$, yet this interaction is in general very
309 weak and we need to develop strategies to enhance it. Molecular engineering may allow to enhance
310 both the Landé g -factor and the total magnetic moment of the ground multiplet beyond ordinary
311 values obtained for single atoms. Another strategy to strengthen the spin-photon coupling is to use
312 spin ensembles. It has been demonstrated in fact that the spin-photon coupling can be enhanced by
313 a factor \sqrt{N} – with N total number of spins- by using collective modes [7]. Although the general
314 problem of coupling photons with two-level system (eg. atoms, molecules with *electric* dipoles etc.)

315 is well documented in literature, only recently experimental investigations focused on spin systems.
316 Strong coupling regime was achieved by using conventional superconducting resonators and NV
317 centers [60] or Er spin impurities in inorganic matrix [61] at mK temperature. We are currently
318 using high T_c superconducting planar resonators [62] that show excellent performances at finite
319 temperature (up to 50K, at least) and in strong magnetic field and, with these, we can achieve
320 strong coupling regime with DPPH [60] and PyBTM [63] organic radicals at liquid helium
321 temperature and even above. It is worth to note that in this latter cases the strong coupling regime
322 was obtained by using concentrate samples that present sharp EPR lines due to exchange
323 narrowing. As concerns mononuclear molecular metal-spin centers, preliminary results on diluted
324 crystals show that we get close -but not yet fully in- to the strong coupling regime [64]. For
325 comparison, a snapshot summary of these preliminary results is plotted in fig.5. It should be
326 emphasized, however, that data reported in fig.5 depend only in part on the intrinsic features of the
327 spin centers. Other parameters, such as temperature, number of spins, cavity characteristics,
328 influence both the spin-photon coupling and the spin resonance linewidth. Thus, for instance, in
329 fig.4 different values are reported for same derivative at different temperatures or spin
330 concentrations. Results on $\text{Cu}(\text{nmt})_2$ crystals are quite encouraging: if we extrapolate trend at low
331 temperature (open symbols) we can safely expect that the strong coupling regime can be achieved
332 below 1K.

333 An alternative approach is to locally enhance the intensity of MW radiation. It has been
334 proposed indeed that nano-structured superconducting strip lines may allow to achieve strong
335 coupling regime even with a single (molecular) spin and this can be used for scalable architectures
336 [65]. Experiments are currently testing different solutions since the detection of tiny signals requires
337 extraordinary sensitivity. By using Josephson Junctions amplification, the detection of small
338 ensemble of about 10^3 spins has been reported and this is, at present, the best performance of EPR
339 nano-detection [66]. The use of electric field component to manipulate spins is a further attractive
340 alternative and several mechanisms have been proposed [67] and they are currently under
341 investigation [68].

342 Achieving strong coupling regime with microwave photons allow to integrate molecular spins
343 in hybrid quantum devices. Superconducting circuits are normally used as a bus to transfer
344 quantum information between different quantum memories and registers. In a recent experiment,
345 we have demonstrated that we can couple two or more distinguished spin ensembles through
346 resonant microwaves photons by spectroscopic measurements [63]. Next steps along this line could
347 be the integration of molecular spins in more complex superconducting circuits with the encoding
348 of sequences of MW pulses. This will eventually allow to perform quantum algorithms [69] or
349 simply to transfer qubit from fast quantum registers to memories. Along this line, pulse MW
350 sequences have been used to efficiently transfer qubits in hybrid superconducting circuits with NV
351 centers [70]. Again, we may learn a lot from fields close to molecular magnetism!

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358 *Figure 5 Coupling strength against spin linewidth parameters as measured in different molecular spins systems.*
359 *Results can be compared with those obtained with NV centers and Er spin defects in YSiO or YAlO. The parameters of*
360 *DPPH and PyBTM organic radicals are taken from ref. 62 and 63 respectively. Data taken at different temperatures are*
361 *indicated by the black arrows and by symbols of different sizes that range from 2 K (larger symbols) to 50 K (smaller*
362 *symbols). Empty symbols display the parameters extrapolated to 0.3 K. The dashed line represent the threshold above which*
363 *strong coupling is achieved. Dashed rectangles show the typical working ranges used for NV centers (blue) and Er*
364 *spin centers (green) coupled to superconducting Nb planar resonators at mK temperature region. (Reproduced from Ref.*
365 *64 with permission from the Consiglio Nazionale delle Ricerche (CNR) and The Royal Society of Chemistry.)*
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372

373 **Author Contributions:** All authors contributed to write this article.

374

375 **Conflicts of Interest:** The authors declare no conflict of interest.

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