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BLACKHOLEWEATHER – Jet-regulated chaotic cold accretion across the meso scale: Variability and kinematics

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ABSTRACT

Context. Chaotic cold accretion (CCA) predicts that supermassive black holes are fed by multiphase clouds condensing from turbulent hot atmospheres. In jet-regulated systems, however, cold gas must also remain dynamically connected to the central accretion region. **Aims.** We investigate how a self-regulated kinetic jet modifies the kinematics, radial transport, and variability of CCA across the meso-scale of a typical galaxy-group atmosphere.

Methods. We develop and analyse two hyper-zoom hydrodynamical simulations of a 100 kpc group atmosphere with radiative cooling, driven subsonic turbulence, sub-pc accretion, and a kinetic mass-loaded bipolar jet. The runs differ only in turbulent driving strength. We measure accretion histories, Eddington ratios, power spectra, phase-separated mass fluxes, projected k-plots, and cooling-to-eddy-time (*C*-ratio) profiles.

Results. Both runs become CCA-fed once precipitation begins, with accretion rising from Bondi-like to strongly super-Bondi values while remaining mostly low-Eddington and mechanically dominated. The strongly stirred run develops an early *stormy* phase with extended condensation, bursty feeding, and strong inflow/outflow variability, but later enters a *cloudy* phase in which cold and warm gas persist at meso- and inner macro-scales while sink coupling weakens. The calmer run instead maintains a compact *rainy* state with a longer-lived central reservoir. Accretion-rate spectra show flicker-like low-frequency slopes and red-noise tails; in the *cloudy* phase, the normalization drops and the low-frequency slope flattens. Phase-separated fluxes show fountain-like recycling in the strongly stirred run, but inner-kpc recycling in the calmer run. The jet excavates a hot channel where sustained condensation is suppressed, while $C \sim 1$ is reached most systematically outside the cone and near the jet-ambient interface.

Conclusions. Jet-regulated CCA is controlled by meso-scale transport, not only by cold-gas production. Within the BLACK-HOLEWEATHER framework, combined k-plot and *C*-ratio diagnostics are crucial to distinguish cold gas that is merely present from cold gas dynamically linked to SMBH feeding.

Key words. black hole physics – accretion, accretion disks – hydrodynamics – methods: numerical – galaxies: active – galaxies: groups: general

1. Introduction

How supermassive black holes (SMBHs) are fed by multiphase gas and how their jets redistribute energy, momentum, and metals through the surrounding atmosphere remain central open problems in galaxy evolution. Within galaxy evolution models, Active Galactic Nuclei (AGNs) are commonly invoked as

the drivers of galactic-scale outflows, powered by the release of gravitational energy during gas accretion onto the SMBH at the centre of the potential well (e.g. [Silk & Rees 1998](#); [King & Pounds 2015](#)). The associated AGN feedback processes, together with stellar feedback from supernovae and stellar winds (e.g. [Barbani et al. 2023, 2025](#)), are believed to significantly alter the thermodynamic and dynamical state of the interstellar and circumgalactic media (ISM and CGM), thereby regulating

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galaxy formation over cosmic time. In physically motivated scenarios, fast winds and relativistic jets launched by AGNs are coupled to the ambient medium (see, e.g. Menci et al. 2008; Zubovas & King 2012; Faucher-Giguère & Quataert 2012). Nevertheless, the ultimate outcome of this coupling remains unclear, with proposed effects ranging from positive feedback through shock-induced star formation to negative feedback via the suppression of cooling and fragmentation driven by turbulence and radiative heating (McNamara & Nulsen 2012; Morganti 2017; Mukherjee et al. 2018; Gaspari et al. 2020, for reviews).

Modeling the feedback cycle driven by accretion and outflows in galaxy groups and clusters is particularly challenging because it requires simultaneously capturing physical processes operating across an extreme range of scales. On the largest end, groups and clusters extend over megaparsec distances, whereas the mechanisms responsible for gas accretion onto the SMBH and the subsequent launch of jets and winds are set on sub-parsec scales, with horizon radii $\sim 10^{-5}$ pc for a $10^8 M_{\odot}$ black hole. In the immediate vicinity of the SMBH, within the $\sim 10\text{--}10^2$ Schwarzschild radii, only a minor fraction of the inflowing gas is ultimately accreted, while the majority is re-ejected through feedback (e.g. Gaspari et al. 2017). The gravitational energy released in this region is efficiently transformed into mechanical power, driving either ultrafast nuclear outflows with velocities of the order $0.1c$ or highly collimated relativistic jets, especially in systems hosting rapidly spinning black holes (e.g. Tombesi et al. 2010; Fabian 2012; Gianolli et al. 2024). These processes are commonly identified as the initial trigger for AGN feedback. In parallel, radiation produced in the innermost regions of the accretion disk contributes to an additional feedback channel, although becoming significant only at high Eddington ratios ($\dot{M}/\dot{M}_{\text{Edd}} \gtrsim 0.1$) (e.g. Ciotti et al. 2010).

Observationally, AGN perturb their gaseous atmospheres through shocks, bubbles, cavities, uplift, and multiphase outflows, with mechanical powers up to $\sim 10^{46}$ erg s^{-1} inferred from X-ray cavity and shock measurements (Fabian 2012; Gitti et al. 2012; Eckert et al. 2021). Over the last decade, ALMA, MUSE, XMM-Newton, and Chandra have shown that these feedback events are not single-phase phenomena: molecular, warm ionized, and X-ray-emitting gas often coexist over the same central regions, with disturbed velocity fields and broad line profiles (Alatalo et al. 2013; Cicone et al. 2014; Fiore et al. 2017; Mainieri et al. 2021; Maiolino et al. 2017; Morganti 2017; Tombesi et al. 2010; Veilleux et al. 2020). Observed outflow rates span roughly $\dot{M}_{\text{out}} \sim 0.1\text{--}100 M_{\odot} \text{ yr}^{-1}$, while characteristic velocities range from $\sim 10^2$ to $\sim 10^4$ km s^{-1} , generally decreasing from highly ionized/X-ray gas to the molecular phase (King & Pounds 2015; Tombesi et al. 2013). However, the same observed multiphase structures can trace direct outflow, turbulent stirring, uplift, fallback, or precipitation-driven inflow. Physically grounded kinematic diagnostics are therefore required to distinguish bulk flows from turbulence and to determine whether condensed gas is actually feeding the SMBH or merely circulating through the jet-regulated atmosphere.

A physically motivated framework connecting kinematics, variability, and SMBH feeding is *chaotic cold accretion* (CCA). In this scenario, radiative cooling in a turbulent hot atmosphere drives nonlinear condensation, producing cold clouds and filaments that “rain” toward the SMBH when cooling and turbulent mixing operate on comparable timescales (Gaspari et al. 2013, 2015, 2017). Collisions among multiphase condensates and chaotic torques can cancel angular momentum, boosting accretion above the smooth hot-mode expectation and producing rapidly variable, flickering SMBH feeding. The resulting jet re-

sponse is therefore intermittent rather than steady. Here we focus on the *meso-scale*, the intermediate $\sim 0.1\text{--}1$ kpc region linking the unresolved SMBH feeding zone to the larger hot halo. This is the scale where condensation, turbulence, jet-driven uplift, and cloud-cloud interactions determine whether cold gas remains connected to the central sink or is instead stirred and recycled. CCA therefore predicts not only multiphase morphology, but also specific kinematic and temporal signatures across this meso-scale: enhanced velocity dispersion, broadened and multi-component line profiles, coexistence of inflowing, uplifted, and outflowing gas, recurrent feedback cycles, and strong variability in accretion and jet power. A key goal of this paper is to determine which of these signatures remain robust in a jet-regulated atmosphere, and how they can diagnose efficient SMBH feeding versus meso-scale recycling.

On the numerical side, general relativity magneto-hydrodynamic (GRMHD) simulations have made major progress in quantifying near-horizon accretion efficiencies, jet launching, and the dependence of mechanical power on magnetic flux and SMBH spin (Tchekhovskoy et al. 2012; Sądowski & Gaspari 2017). Conversely, semi-analytic and idealized galaxy-scale models show that ultrafast nuclear outflows and jets can shock, entrain, and accelerate the ambient medium, producing cooler and broader gas components as energy and momentum are transferred from the nuclear to galactic scales (Faucher-Giguère & Quataert 2012; Wagner et al. 2013; Zubovas & King 2014; Serafinelli et al. 2019). The remaining challenge is to bridge these regimes. CCA is intrinsically stochastic: condensation, cloud collisions, angular-momentum cancellation, jet response, uplift, and fallback evolve over many local dynamical times. Capturing this behaviour therefore requires long-duration, high-resolution simulations that follow cooling, turbulence, jet feedback, and multiphase kinematics with sufficient cadence to measure both velocity fields and accretion variability across the meso-scale.

This motivates state-of-the-art computational approaches. Here we leverage the open-access AthenaPK¹ code (Grete et al. 2023), which combines portability, scalability, and astrophysical specialization with GPU-resident memory handling, asynchronous one-sided MPI communication, and block-level tasking. This work is part of the BLACKHOLEWEATHER project (PI: Gaspari), which aims to build a unified, multiscale description of SMBH feeding and feedback in multiphase atmospheres. The present paper, C26b, is the second of two companion studies based on the same jet-regulated simulation suite. The companion paper, Cammelli et al. (2026), hereafter C26a, focuses on morphology and thermodynamics, including phase structure, condensation morphology, cavity evolution, and shock/cocoon phenomenology. Here, we investigate the kinematic and time-domain response of the atmosphere: velocity-field statistics, phase-separated inflow and outflow rates, turbulence-related diagnostics, and the time-resolved coupling between SMBH feeding and jet power. Our goal is to identify which kinematic and variability signatures trace efficient chaotic cold accretion, and which instead indicate that condensed gas is being stirred, uplifted, or recycled across the meso-scale rather than delivered to the central engine. A central novelty of this work is therefore to separate multiphase gas formation from multiphase gas delivery: cold gas at meso-scales is not by itself a sufficient proxy for SMBH feeding, unless it remains dynamically coupled to the micro-scale sink. Detailed synthetic spec-

¹ <https://github.com/parthenon-hpc-lab/athenapk>

troscopic and integral-field diagnostics are deferred to future BLACKHOLEWEATHER work.

In this paper, we investigate how self-regulated AGN jet feedback reshapes the multiphase gas in a galaxy-group atmosphere. We perform three-dimensional hydrodynamical simulations with AthenaPK, resolving the central sink and jet-injection region with sub-pc cell size while following the intra-group medium over tens of kpc. The unresolved disk-to-horizon physics is encoded through an effective accretion and feedback prescription, rather than explicitly resolving the black-hole horizon or the jet-launching region. The simulations include gravity, radiative cooling, externally driven turbulence, sink accretion, and mass-loaded kinetic jet injection. Static mesh refinement (SMR) provides the required dynamic range, allowing us to study how jet-driven feedback couples micro-scale feeding to meso-scale multiphase transport and group-scale gas circulation.

In practice, we perform consistent three-dimensional hydrodynamic (HD) simulations targeting galaxy-group environments, spanning from the inner launching region (resolved to sub-parsec scales via static mesh refinement, SMR) up to group/cluster scales (tens of kpc). We follow multiphase gas evolution under gravity, cooling, turbulence driving, jet launch and energy deposition within AthenaPK, and extract kinematic and time-domain diagnostics relevant to CCA and feedback self-regulation.

Within the BLACKHOLEWEATHER series, the present work is also tied to two complementary sets of companion simulations that provide useful reference points for later comparisons. The no-jet CCA simulations of Barbani et al. (2026a,b), hereafter B26a,b, isolate how ambient turbulence regulates multiphase condensation, radial transport, and pure feeding/rain in the absence of self-regulated jet feedback, thereby providing the natural baseline against which we assess what the jet changes. The spin-regulated simulations of Piana et al. (2026a,b), hereafter P26a,b, instead explore how unresolved angular-momentum transport, SMBH spin evolution, and relativistic prescriptions affect the innermost feeding-feedback loop.

The paper is organized as follows. §2 summarizes the subgrid accretion and jet-feedback prescription that links sink-scale gas inflow to SMBH feeding and mechanical outflow, and describes the numerical framework, initial conditions, and reference simulations. §3 presents the main results on accretion variability, radial inflow/outflow rates, projected gas kinematics, and CCA diagnostics. §4 discusses the implications for jet-regulated CCA, observational interpretation, and AGN self-regulation across different weather states. Finally, §5 summarizes our conclusions.

2. Numerical methods

Our simulations are performed with AthenaPK, the GPU-accelerated extension of Athena++, built on the Parthenon and Kokkos frameworks (Stone et al. 2020; Grete et al. 2023; Edwards et al. 2014). Since the full numerical setup is presented in C26a and B26a, here we summarize only the ingredients most relevant for the present analysis.

We simulate a 100 kpc box with static mesh refinement (SMR) over 10 levels, starting from a 128^3 root grid and reaching a finest spatial resolution of $\Delta x_{\min} \simeq 0.78$ pc. The inner sink region is centred at the origin and has radius $r_{\text{sink}} = 4\Delta x_{\min}$ (cf. Gaspari et al. 2013 and B26a). We adopt the same gravitational potential, radiative cooling, turbulence driving, and initial conditions as in C26a and B26a. Magnetic fields are neglected

in this work in order to isolate the role of kinetic feedback and ambient turbulence.

The gas dynamics are evolved in conservative Eulerian form as

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = \mathcal{S}_\rho^{\text{jet}}, \quad (1)$$

$$\frac{\partial (\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \otimes \mathbf{v} + P \mathbb{I}) = -\rho \nabla \Phi + \rho \mathbf{f}_{\text{turb}} + \mathcal{S}_{\rho \mathbf{v}}^{\text{jet}}, \quad (2)$$

$$\frac{\partial E}{\partial t} + \nabla \cdot [(E + P)\mathbf{v}] = -\rho \mathbf{v} \cdot \nabla \Phi - \mathcal{L} + \mathcal{S}_{\text{turb}} + \mathcal{S}_E^{\text{jet}}, \quad (3)$$

where $\mathcal{L} = n^2 \Lambda(T)$ is the radiative cooling term and $\mathcal{S}_{\text{turb}} = \rho \mathbf{v} \cdot \mathbf{f}_{\text{turb}}$ is the turbulence-driving source term. We use the same second-order Godunov scheme, PLM reconstruction, HLLC solver, and fallback strategy described in C26a.

2.1. Subgrid feedback prescription

Because the sink particle does not resolve the flow from the accretion disc down to the black-hole horizon or the innermost stable circular orbit (ISCO), we map the sink accretion rate, \dot{M}_{sink} , to an SMBH accretion rate, \dot{M}_\bullet , through

$$\dot{M}_\bullet = (1 - \epsilon_{\text{disk}}) \dot{M}_{\text{sink}}, \quad (4)$$

where ϵ_{disk} denotes the fraction of gravitational energy dissipated between the sink radius and the ISCO. Under the assumption of a radiatively efficient disc, this is equivalent to the disc radiative efficiency. In principle these efficiencies depend on BH spin; here we keep them fixed to isolate the impact of feedback (see P26a for more details on spin-jet coupling).

The unresolved SMBH feeding rate is coupled, neglecting magnetic fields, to the injected feedback power through

$$P_{\text{tot}} = P_{\text{kin}} + P_{\text{th}} = \epsilon_{\text{tot}} \dot{M}_\bullet c^2 \equiv (\epsilon_{\text{kin}} + \epsilon_{\text{th}}) \dot{M}_\bullet c^2, \quad (5)$$

with the corresponding mass-loaded jet rate

$$\dot{M}_{\text{jet}} \equiv \eta_{\text{jet}} \dot{M}_\bullet, \quad (6)$$

where η_{jet} is the adopted mass-loading factor (B26a). In these fiducial runs we set $\eta_{\text{jet}} = 1 - \epsilon_{\text{tot}}$, so that most of the gas reaching the unresolved nuclear region is returned to the resolved atmosphere through the jet. This choice is motivated by GRMHD models, in which only a small fraction of the inflowing gas is ultimately accreted while the rest is expelled in mechanically dominated outflows (e.g. Sądowski & Gaspari 2017; Gaspari & Sądowski 2017).

In this work, we consider a kinetically-dominated jet, and retain the thermal term only as an effective internal-energy contribution at injection. In this limit

$$P_{\text{tot}} \simeq P_{\text{kin}} + P_{\text{th, jet}} = \epsilon_{\text{tot}} \dot{M}_\bullet c^2, \quad (7)$$

where

$$P_{\text{kin}} = \frac{1}{2} \dot{M}_{\text{jet}} v_{\text{jet}}^2, \quad (8)$$

$$P_{\text{th, jet}} = e_{\text{jet}} \dot{M}_{\text{jet}}, \quad (9)$$

and $e_{\text{jet}} = k_B T_{\text{jet}} / [(\gamma - 1)(\mu m_p)]$ is the internal energy per unit mass of the injected gas. Combining the above equations yields

$$v_{\text{jet}} = \sqrt{\frac{2}{1 - \epsilon_{\text{tot}}} [\epsilon_{\text{tot}} c^2 - (1 - \epsilon_{\text{tot}}) e_{\text{jet}}]}. \quad (10)$$

This prescription should be interpreted as a controlled sub-grid coupling between the pc-scale sink and the resolved atmosphere: it does not attempt to model the disc-to-horizon launching process directly, but imposes a fixed relation between unresolved feeding and injected mechanical power.

2.2. Jet injection

The feedback module injects mass, momentum, and energy in two oppositely directed cylindrical regions centred on the sink and aligned with the $\pm z$ axis. The jet geometry and numerical implementation follow C26a. Within the injection region, the source terms are

$$\mathcal{S}_{\rho}^{\text{jet}} = \dot{\rho}_{\text{jet}}(W_+ + W_-), \quad (11)$$

$$\mathcal{S}_{\rho v}^{\text{jet}} = \dot{\rho}_{\text{jet}} v_z \hat{z} W_+ - \dot{\rho}_{\text{jet}} v_z \hat{z} W_-, \quad (12)$$

$$\mathcal{S}_E^{\text{jet}} = \dot{\rho}_{\text{jet}} \left(e_{\text{jet}} + \frac{v_z^2}{2} \right) (W_+ + W_-), \quad (13)$$

where $v_z = v_{\text{jet}}$ is given by Eq. (10), $\dot{\rho}_{\text{jet}}$ is set by Eq. (6) and the injection volume, and W_{\pm} are normalized kernels describing the two jet lobes. As in C26a, the jet is injected with a fixed axis and no opening angle, allowing us to isolate the role of ambient turbulence in shaping the resolved kinematics and variability.

2.3. Reference runs

We analyze the same two reference runs presented in C26a, which differ only in the level of initial turbulence driving. They bracket two atmospheric ‘‘weather’’ regimes: a **low-turb** case with weaker non-thermal support and earlier, more centrally retained condensation, and a **high-turb** case with stronger stirring, delayed precipitation, and more extended multiphase development. In both runs, turbulence is first allowed to develop before radiative cooling and jet feedback are activated; the subsequent evolution is analyzed in terms of the rain time t_{rain} defined in C26a.

2.4. Initial conditions & simulation suite

As introduced above, all simulations carried out in this work have been performed with the AthenaPK code using a 100 kpc side box at 10 static refinements levels, which aligns with the resolution of C26a. Throughout the whole set of simulations, we assume the same initial density and temperature profiles as for a static hot gaseous halo in hydrostatic equilibrium. The static gravitational potential includes the dark matter halo, the central dominant group galaxy, and the SMBH, whose mass is updated according to Eq. 4. The initial profiles are discussed in detail in B26a, including the data sample and the fitting procedure used to obtain the initial profiles. Here we note that we opt to simulate a galaxy group (reference parameter values as in B26a) as this configuration shows a consistently faster evolution with respect to a galaxy cluster. In fact, the cooling time in a galaxy group, which drives the criterion for the condensation of gas into CCA episodes (Gaspari et al. 2013, 2017), results in about

Table 1. Summary of the numerical setup and reference runs. Parameters common to both runs are listed once; run-dependent quantities are given in the last two columns. The acceleration field has an RMS amplitude a_{rms} and is injected with a solenoidal fraction $\zeta = 1$, corresponding to purely incompressible driving peaking at mode number $n_{\text{peak},v}$. The $N_{\text{modes},v}$ random Fourier modes refresh every correlation time t_{corr} .

Parameter	low-turb	high-turb
<i>Common numerical setup</i>		
L_{box} [kpc]		100
Root grid		128^3
SMR levels		10
Δx_{min} [pc]		0.78
r_{sink} [pc]		≈ 3.1
T_{sink} [K]		1 K
ρ_{sink} [g cm^{-3}]		1×10^{-30}
d_{jet} [pc]		≈ 1.5
M_{\bullet} [M_{\odot}]		2.8×10^8
$\epsilon_{\text{jet}} (= \epsilon_{\text{tot}})$		0.028
T_{jet} [K]		5×10^8 K
v_{jet}		$\approx 0.25 c$
$n_{\text{peak},v}$		4
l_{inj} [kpc]		$L_{\text{box}}/n_{\text{peak}}$
$N_{\text{modes},v}$		64
t_{corr} [Myr]		30
<i>Run-dependent quantities</i>		
a_{rms} [$10^{-8} \text{ cm s}^{-2}$]	0.62	1.55
3D Mach number, \mathcal{M}	~ 0.15	~ 0.4
$\sigma_v(t_{\text{restart}})$ [km s^{-1}]	60–90	210–230
t_{rain} [Myr]	~ 9	~ 16
Total duration after restart [Myr]	~ 65	~ 115
Total simulated duration [Myr]	~ 100	~ 150

$t_{\text{cool}} \approx 10 - 20$ Myr in the central kpc in our simulations, typically one or two orders of magnitude shorter than a galaxy cluster. This allows us to test our implemented physics in a computationally efficient manner and, on the other hand, the large scale structure of the group/cluster is not expected to affect the micro-scale phenomena we have implemented as described in the previous section.

Following the analysis carried out in C26a, in this work we employ two different scenarios in terms of initial conditions at both strong (**high-turb**) and weak (**low-turb**) turbulence to better study SMBH feeding and feedback processes in realistic environments. For reference, the main parameters of the turbulence driving used in C26a are shown in Table 1. Following B26a, we drive turbulence solenoidally at $n_{\text{peak}} = 4$, corresponding to $l_{\text{inj}} = L_{\text{box}}/n_{\text{peak}}$. This idealized forcing scale is motivated by the large-scale vortical stirring observed and inferred in cool-core hot halos, where AGN bubbles, jet-driven uplift, minor mergers, and sloshing generate subsonic motions over core scales. The purely solenoidal prescription allows us to isolate the role of turbulent mixing in CCA condensation and jet-regulated feeding (Fabian 2012; McNamara & Nulsen 2012; Vazza et al. 2012; Gaspari & Churazov 2013; Hofmann et al. 2016; Hitomi Collaboration et al. 2016; XRISM Collaboration et al. 2025).

It is important to note that in this work we investigate a specific case of the above described injection mechanism. In particular, we limit our efforts to studying the effects and the consequences of a kinetically-dominated jet, i.e. setting the kinetic fraction to 1 in Eq. 7. This choice is motivated by the need to isolate how injected kinetic energy is redistributed into thermal, turbulent, and multiphase gas motions across resolved scales. In

addition, we will make use of our findings in future companion papers in the `BLACKHOLEWEATHER` series, in accordance with the project aim to link and study AGN accretion and feedback at different scales and phases as a whole, following a *bottom-up*, first-principles approach (Gaspari et al. 2020). However, we also note that the current implementation does not account for magnetic fields and extra physics, such as relativistic effects and AGN winds, which could potentially alter the results. We defer the study of these limitations and caveats to a future work.

3. Results

In this section, we present the main results of the analysis for the two reference simulations, `high-turb` and `low-turb`, introduced in §2. The purpose of this work is to investigate how AGN feedback, injected on the smallest resolved scales, propagates through and reshapes the surrounding medium under two different turbulence regimes. In particular, turbulence driving and the ensuing cascade regulate both the degree of mixing and the thermodynamic fluctuations of the gas (e.g. Gaspari & Churazov 2013). As a result, the `high-turb` run develops a stronger level of non-thermal support than the `low-turb` case, delaying the onset of local thermal instability and therefore postponing the formation of condensed structures. The production runs are initialized from pre-evolved turbulent atmospheres with the target velocity dispersions listed in Table 1; radiative cooling and self-regulated jet feedback are then enabled at the restart time, which defines $t = 0$ Myr for the analysis. We define t_{rain} as the time, measured from this restart, at which the first cold clouds and filaments form, condense, and produce a sustained increase in sink accretion.

Throughout the analysis, the subgrid feedback prescription is kept fixed (see §2), so variations in jet power directly trace variations in \dot{M}_\bullet . This allows us to focus on the resolved problem: how condensation, stirring, uplift, and meso-scale transport regulate the coupling between multiphase gas and SMBH feeding. Variations in the unresolved disc-to-horizon physics or jet geometry are deferred to future `BLACKHOLEWEATHER` work.

Following the multiscale and multiphase framework adopted throughout this work, we divide the gas into four characteristic radial domains, defined primarily with respect to the meso-scale region where precipitation, jet-induced circulation, and inward transport are most directly coupled. As per `BLACKHOLEWEATHER` framework (Gaspari et al. 2020), we refer to the *micro* scale as the inner region where gas is already approaching the unresolved accretion zone ($r < 0.1$ kpc), the *meso* scale as the central transport and coupling layer between condensation and SMBH feeding ($0.1 \text{ kpc} < r < 1 \text{ kpc}$), the inner *macro* scale as the region where extended filaments, uplift, and halo circulation connect to the meso-scale ($1 \text{ kpc} < r < 10 \text{ kpc}$), and the outer *macro* scale as the large-scale hot-atmosphere reservoir ($r > 10 \text{ kpc}$). To track how the gas changes across these regions, we separate it into five temperature-defined phase proxies (tied to realistic observational bands): molecular-temperature gas ($T < 2 \times 10^2 \text{ K}$), cold atomic gas ($2 \times 10^2 \text{ K} \leq T < 1.6 \times 10^4 \text{ K}$), warm gas ($1.6 \times 10^4 \text{ K} \leq T < 1.16 \times 10^6 \text{ K}$), soft-X-ray-temperature gas ($1.16 \times 10^6 \text{ K} \leq T < 5.8 \times 10^6 \text{ K}$), and hard-X-ray-temperature gas ($T > 5.8 \times 10^6 \text{ K}$).

As in C26a, we use the terms *stormy*, *rainy*, and *cloudy* as compact descriptors of CCA weather states, again mainly in relation to how the meso-scale connects large-scale condensation to micro-scale feeding. A stormy state denotes spatially extended, strongly mixed, and bursty precipitation, in which condensation spreads from the meso-scale into the surrounding halo and drives

intermittent feeding. A rainy state denotes a more compact and coherent condensation cycle, where the meso-scale remains efficiently connected to the central sink and sustains SMBH fueling. A cloudy state denotes a multiphase atmosphere in which cold gas is still present, especially on meso- and macro-scales, but its coupling to direct SMBH feeding is weakened. These labels summarize the combined thermodynamic, kinematic, and accretion behaviour; they are not separate simulation inputs.

The results are presented as follows. We first study the time and space variability and accretion properties in §3.1. In §3.2 we investigate the inflow and outflow properties of the multiphase gas. The kinematic state of gas is described in §3.3.1 under different BH weather conditions. The C-ratio profiles are finally discussed in §3.3.2 to probe different CCA regimes.

3.1. Variability & accretion properties

SMBH accretion is intrinsically time-variable, and in a self-regulated feedback model this variability directly modulates the jet power. Observationally, such variability can appear in optical/UV continuum luminosity changes and in the broadening or narrowing of emission-line profiles (Tozzi et al. 2022; Cammelli et al. 2025; Paolillo & Papadakis 2025). In the CCA framework, this flickering behaviour is expected to arise from the stochastic arrival, interaction, and partial angular-momentum cancellation of multiphase clouds and filaments near the unresolved accretion region. In our simulations, gas accreted through the pc-scale sink is mapped onto the unresolved central engine through the prescription described in §2. We therefore use $\dot{M}_\bullet(t)$ as a tracer of the time-dependent CCA feeding cycle and of the corresponding jet-power response. As shown in the no-jet simulations of B26b, CCA feeding naturally produces scale-dependent variability: the relative amplitude of \dot{M}_\bullet fluctuations increases with the temporal window over which the accretion history is sampled, producing flicker-like accretion histories rather than white-noise fluctuations. Here we examine how this intrinsic CCA variability is modified once the accretion rate is coupled to a self-regulated kinetic jet in the `high-turb` and `low-turb` reference runs.

3.1.1. Feeding the SMBH

Figure 1 shows the accretion-rate evolution of the two reference runs, `high-turb` and `low-turb`, in units of the corresponding Bondi accretion rate, \dot{M}_{Bondi} , with time normalized to t_{rain} . This normalization provides a useful hot-mode baseline (Bondi 1952): values close to unity indicate accretion rates comparable to smooth, quasi-spherical capture from the hot atmosphere, whereas strongly super-Bondi values mark the transition to multiphase, precipitation-driven feeding. Bondi-like prescriptions have been widely adopted as subgrid SMBH accretion models in cosmological and galaxy-scale simulations (e.g. Springel et al. 2005; Booth & Schaye 2009), and have also been used as observationally motivated estimates of hot-mode feeding in nearby systems (e.g. Allen et al. 2006).

In the present CCA context, however, \dot{M}_{Bondi} should be interpreted only as a reference scale rather than as a predictive accretion model, because the flow is turbulent, multiphase, anisotropic, and directly affected by jet feedback. We define

$$\dot{M}_{\text{Bondi}} = 4\pi \lambda(\gamma) \frac{(GM_\bullet)^2 \rho_\infty}{c_{s,\infty}^3}, \quad (14)$$

where $\lambda(\gamma)$ is a dimensionless factor of order unity that depends on the adiabatic index γ (here assumed 5/3), while ρ_∞

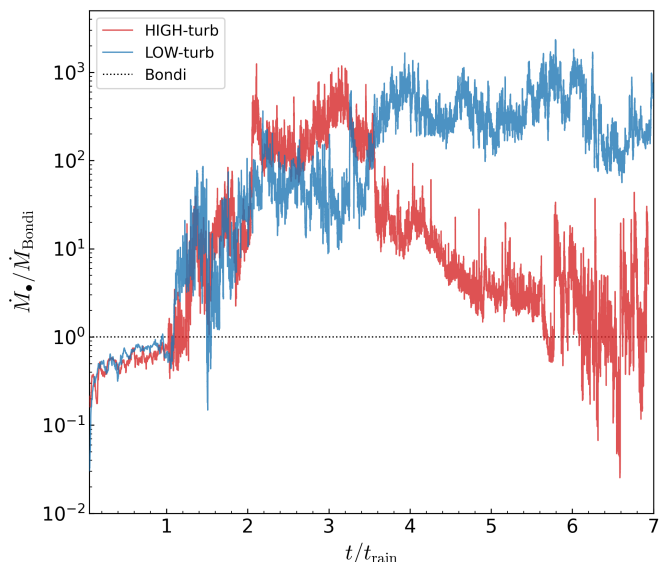


Fig. 1. Black hole accretion rate as a function of time in the *low-turb* and *high-turb* case, as indicated in the legend. The simulation time has been normalized to visually compare the different runs starting from the time they have been restarted with the jet and the cooling on. During active jet phases, \dot{M}_\bullet reaches large super-Bondi values (cf. B26a,b). The dotted horizontal line marks the Bondi reference accretion rate, $\dot{M}_\bullet / \dot{M}_{\text{Bondi}} = 1$.

and $c_{s,\infty}$ are the ambient density and sound speed measured in the hot atmosphere at large radii. The assumptions entering this expression, namely a smooth, homogeneous, radially symmetric, and adiabatic hot inflow, are not satisfied in realistic jet-regulated CCA atmospheres. For this reason, and consistently with B26a,b, we use \dot{M}_{Bondi} only to quantify departures from the idealized hot-mode limit and to identify when precipitation boosts the accretion rate into a genuinely multiphase feeding regime.

During the first normalized rain time, both runs evolve in a similar way, slowly rising from sub-Bondi values toward $\dot{M}_\bullet / \dot{M}_{\text{Bondi}} \sim 1$. The *high-turb* case remains slightly lower than *low-turb* in this initial stage, consistent with the stronger large-scale stirring already discussed in C26a, which enhances mixing and temporarily hinders the delivery of gas to the sink. At these early times, the jet has not yet had enough time to reshape the inner multiphase structure in a substantial way. Once condensation sets in, both runs move rapidly into a super-Bondi regime, marking the transition from smooth hot feeding to genuine CCA. The subsequent evolution closely mirrors the thermodynamic “weather” states identified in C26a. In the *high-turb* run, the accretion ratio rises abruptly after $t/t_{\text{rain}} \sim 1$ and reaches a short but intense super-Bondi episode, peaking at $\dot{M}_\bullet / \dot{M}_{\text{Bondi}} \sim 10^2 - 10^3$ around $t/t_{\text{rain}} \sim 2-3.5$. This is the accretion counterpart of the early *stormy* phase: the atmosphere hosts vigorous precipitation, broad phase mixing, and rapid delivery of dense condensates toward the nucleus. After $t/t_{\text{rain}} \sim 3.5$, however, the accretion rate declines by more than an order of magnitude and becomes increasingly intermittent, reaching on average values slightly larger or comparable to the Bondi case. This behaviour is consistent with the later *cloudy* phase identified in the previous paper, in which cold clouds and filaments are still present at micro- and meso-scales but are less efficiently channeled into the sink region. In other words, the high-turbulence atmosphere

Table 2. Broken power-law fits to the accretion-rate PSD. We fit the smoothed PSD with $P(f) = A \begin{cases} (f/f_b)^{\alpha_1}, & f < f_b \\ (f/f_b)^{\alpha_2}, & f \geq f_b \end{cases}$, where f_b is the break (knee) frequency. Reported uncertainties correspond to the formal 1σ fit errors.

Run	α_1	α_2	$\log_{10}(f_b / [\text{Myr}^{-1}])$
<i>low-turb</i>	-1.24 ± 0.06	-3.81 ± 0.19	1.55 ± 0.06
<i>high-turb</i>	-1.16 ± 0.05	-3.77 ± 0.15	1.48 ± 0.05
<i>high-stormy</i>	-1.24 ± 0.04	-3.66 ± 0.12	1.48 ± 0.04
<i>high-cloudy</i>	-0.79 ± 0.09	-3.50 ± 0.20	1.15 ± 0.07

continues to form multiphase structure, but its late-time precipitation is less effective at sustaining direct SMBH feeding.

The *low-turb* run follows a different path. After the initial transition to super-Bondi accretion, the feeding history remains elevated for a much longer interval, with $\dot{M}_\bullet / \dot{M}_{\text{Bondi}}$ persistently in the $\sim 10^2 - 10^3$ range over several t_{rain} . This prolonged active state is fully consistent with the results of C26a, which showed that the low-turbulence run retains a denser and more persistent molecular/cold reservoir within the inner kpc. In this calmer atmospheric state, condensation is less explosive but more effectively retained near the centre, so that the SMBH continues to be fueled over an extended period rather than through a single dominant burst. Thus, while the *high-turb* run produces an earlier and more impulsive accretion episode, the *low-turb* run supports a longer-lived, reservoir-fed mode of CCA. Read together with the multiphase diagnostics of C26a, Figure 1 shows that ambient turbulence controls not only the onset of condensation, but also how efficiently the resulting cold gas can feed the SMBH: either through short-lived *storm*-fed accretion episodes or through a longer-lasting *rainy*, reservoir-supported mode. Thus, the key distinction is not whether condensation occurs in both runs, but whether the resulting cold gas remains dynamically connected to the central sink.

3.1.2. Accretion-rate spectral analysis

To characterize the variability of the accretion flow beyond the time-domain behaviour, we compute the power spectral density (PSD) of \dot{M}_\bullet , shown in Figure 2. The PSD is calculated, in a given simulation time range, as the square modulus of the Fourier transform per frequency unit of the \dot{M}_\bullet . We restrict the displayed and fitted range to $f \lesssim 10^3 \text{ Myr}^{-1}$.² Solid thin transparent lines indicate the pure PSD of the \dot{M}_\bullet in the frequency space, while thick lines denote the average interpolated binned PSD in the corresponding spectral range. We attempt to fit the interpolated PSD by using a double power-law function, which parameters and relative errors are reported in Table 2. The upper panel compares the full \dot{M}_\bullet time series of the *high-turb* and *low-turb* runs over the range $1 - 7 t_{\text{rain}}$, once radiative cooling sets in. In both cases the PSD is well described by a broken power law, with a low-frequency slope close to $f^{-1.2}$ and a much steeper high-frequency tail, $f^{-3.8}$. This implies that the accretion variability is not white-noise dominated, but instead displays a flicker-like/pink-noise behaviour (cf. Gaspari et al. 2017) on long timescales, which steepens into a strongly red spectrum at high frequencies, as also found in the pure CCA feeding simulations (B26b). The two runs therefore share a similar under-

² Although the \dot{M}_\bullet is output at every integration time step, $\sim 3 \times 10^{-6}$ Myr, the PSD at higher frequencies is dominated by a quasi white-noise flattening arising from the finite numerical resolution of the simulation.

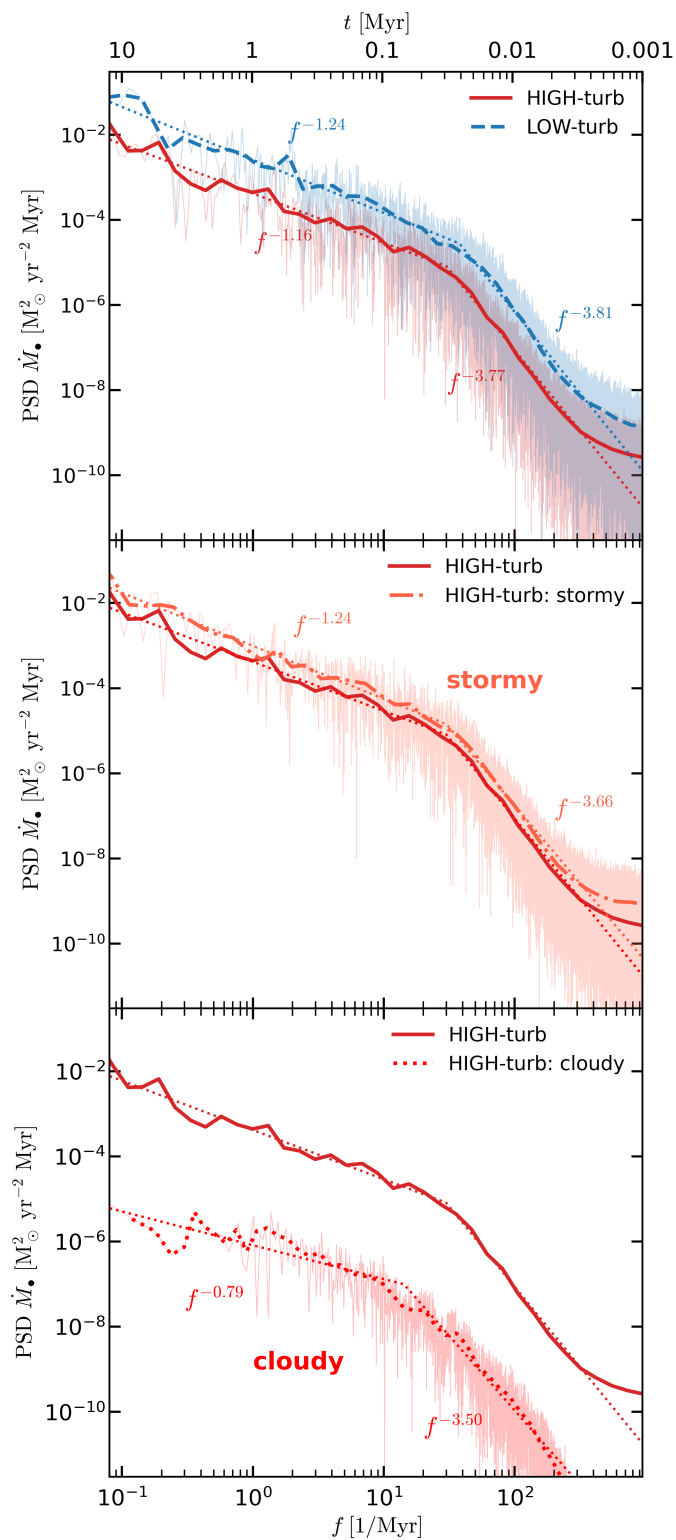


Fig. 2. Power spectral density (PSD) as a function of frequency in the low-turb and high-turb case, as indicated in the legend, from the time the jet and the cooling have been restarted. Light solid lines show the direct PSD from the \dot{M}_* reported in Figure 1, while dark solid lines depict the smoothed version using a Gaussian filter. Dashed lines instead represent a broken power law fit to the dark solid lines. Coloured text indicates the correspondent slope of the fitted curves in separate regimes. Upper panel: full ($t/t_{\text{rain}} = 1 - 7$) evolution of the high-turb and low-turb cases, with a low-frequency slope close to $f^{-1.2}$ and a much steeper high-frequency tail, $\sim f^{-3.8}$. Middle/bottom panels: the full high-turb PSD is compared to the stormy and cloudy regime, respectively.

lying variability cascade, but differ substantially in normalization: the low-turb run carries systematically more power at nearly all frequencies, consistent with its more persistent super-Bondi feeding history and longer-lived central cold reservoir. The high-turb run, instead, shows a lower overall PSD normalization once the full evolution is considered, reflecting the fact that its early bursty accretion phase is followed by a later suppression of direct SMBH feeding.

We also note that the two cases show remarkably similar break frequencies, $f_b \approx 30 \text{ Myr}^{-1}$ ($\log(f_b/[Myr^{-1}]) \approx 1.48$) in the high-turb and $f_b \approx 37.8 \text{ Myr}^{-1}$ ($\log(f_b/[Myr^{-1}]) \approx 1.55$) in the low-turb, corresponding to break times of $t_b \approx 0.033 \text{ Myr}$ and $t_b \approx 0.026 \text{ Myr}$, respectively. This indicates that the typical timescales for the multiphase gas to lose temporal coherence (likely due to cloud collisions) result similarly in the two reference cases. This behaviour differs, in terms of break frequency, from the study presented in B26b where f_b shows a remarkable separation due to the strength of the turbulence driving ($\approx 9 \text{ Myr}^{-1}$ in the low-turbulence scenarios versus $\approx 53 \text{ Myr}^{-1}$ of the high-turbulence case). This suggests that in the case of jet-regulated CCA, the dependence on the turbulence driving strength is somehow alleviated by the extra term introduced by the jet turbulence itself, which affects the collisional dynamics of multiphase clouds and filaments, and therefore impacts their accretion signal down to the central SMBH. As discussed in B26b, the implications of t_b are physically motivated when considering that t_b in our simulations is much longer than the integration timestep over which \dot{M}_* is updated ($\sim 10^{-6} \text{ Myr}$), and far shorter than the correlation time implemented in the turbulence driving (see Table 1, $t_{\text{corr}} = 30 \text{ Myr}$).

This interpretation becomes clearer when the two weather states identified in the high-turb run are isolated. In the middle panel, the stormy interval retains a low-frequency slope of $\propto f^{-1.24}$, i.e. close to pink noise (cf. Gaspari et al. 2017), together with a steeper red-noise high-frequency tail ($\propto f^{-3.66}$) and a consistent $f_b \approx 30 \text{ Myr}^{-1}$. This indicates a regime of correlated, scale-coupled accretion variability, driven by cloud-cloud collision, in which precipitation events remain linked across a broad range of timescales. In the lower panel, the cloudy interval is instead characterized by a much lower overall normalization and a flatter low-frequency slope, $\propto f^{-0.79}$, i.e. intermediate between white and pink noise, while the high-frequency tail retains the trend of a red noise ($\propto f^{-3.50}$). A possible interpretation is that the cloudy phase affects the long-timescale accretion variability: the large, coherent rain episodes typical of the stormy state are replaced by weaker and less correlated feeding events, even though rapid small-scale flickering is still present. This behaviour is consistent with the thermodynamic picture from C26a, in which the stormy phase corresponds to vigorous precipitation directly feeding the nucleus, whereas the cloudy phase still forms cold clouds and filaments on micro- and meso-scales but couples them less efficiently to the sink.

A complementary comparison is provided by P26b, where jet feedback includes spin-dependent precession and the turbulence-driving history is varied. In those simulations, the low-frequency slopes also remain close to $P(f) \propto f^{-1}$, while the main differences appear in the PSD normalization and in the high-frequency damping. This suggests that the low-frequency flicker-noise regime is common to the different CCA realizations, whereas the detailed PSD shape depends on the specific coupling between turbulence, jet feedback, and the multiphase inflow.

3.1.3. Eddington ratios

To fully characterize the accretion properties of our simulations, it is useful to study not only the absolute contribution of the \dot{M}_\bullet but also its role relative to the mass of the central BH. Figure 3 shows the probability density function (PDF) of the BH accretion rates normalized to the Eddington accretion rate \dot{M}_{Edd} , corresponding to the Eddington luminosity $L_{\text{Edd}} = 4\pi GM_\bullet m_p c / \sigma_T$, where $\sigma_T = 6.65 \times 10^{-25} \text{ cm}^2$ is the Thomson scattering cross section, $m_p = 1.67 \times 10^{-24} \text{ g}$ is the proton mass and $c = 3 \times 10^{10} \text{ cm s}^{-1}$ is the speed of light. Note that the PDF is obtained using equally spaced logarithmic bins, such that the integral in $d \log(\lambda)$ is equal to 1. With a fiducial SMBH mass $M_\bullet = 2.8 \times 10^8 M_\odot$ and a reference radiative efficiency $\epsilon_{\text{rad}} = 0.1$, this gives

$$\dot{M}_{\text{Edd}} = \frac{L_{\text{Edd}}}{\epsilon_{\text{rad}} c^2} \simeq 6.2 M_\odot \text{ yr}^{-1}. \quad (15)$$

The ratio $\lambda_{\text{Edd}} \equiv \dot{M}_\bullet / \dot{M}_{\text{Edd}}$ provides a dimensionless metric that is key to evaluating the role of the central engine in relative terms, and relatively easy to access in different spectral bands. Observationally, λ_{Edd} is usually inferred as $L_{\text{bol}} / L_{\text{Edd}}$, rather than as a direct mass-accretion-rate ratio. The bolometric luminosity is commonly estimated either from multi-wavelength SED modelling or from band-dependent bolometric corrections applied to optical/UV, X-ray, infrared, or emission-line tracers (e.g. Lamas-tra et al. 2009; Duras et al. 2020; Gupta et al. 2025), while M_\bullet is inferred from reverberation-mapped or single-epoch virial estimators in unobscured AGN, and from host-galaxy scaling relations, dynamical modeling, or indirect X-ray/line proxies in obscured and low-luminosity systems (e.g. Vestergaard & Peterson 2006). Our simulated λ_{Edd} should therefore be interpreted as an accretion-rate analogue rather than as a forward-modeled luminosity ratio, since converting \dot{M}_\bullet into L_{bol} would require assumptions about radiative efficiency, accretion mode, obscuration, and bolometric corrections.

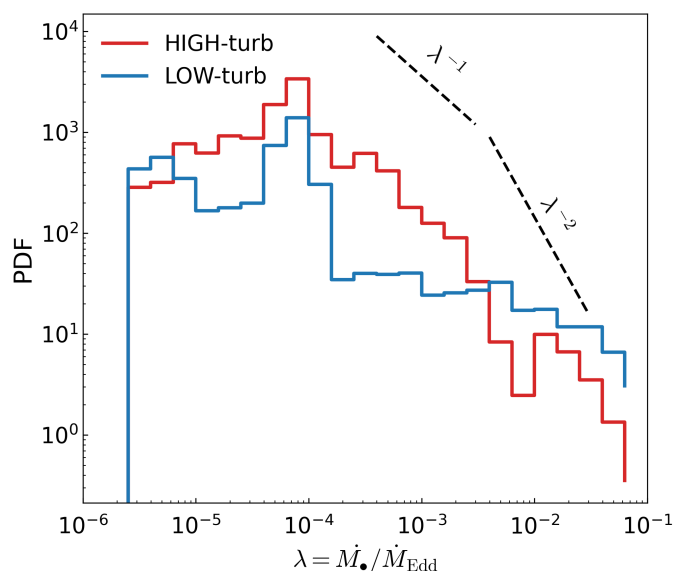


Fig. 3. Probability density function (PDF), normalized to unity in logarithmic space, of the Eddington ratios $\lambda_{\text{Edd}} \equiv \dot{M}_\bullet / \dot{M}_{\text{Edd}}$ as derived from the \dot{M}_\bullet for the 2 reference runs, high-turb (red solid line) and low-turb (blue solid line), computed at each time step. Black dashed guide lines indicate approximate slopes of -1 and -2 in the high- λ_{Edd} tail.

Focusing on our fiducial high-turb and low-turb runs, both distributions peak at $\lambda_{\text{Edd}} \simeq 8 \times 10^{-5}$, extend down to $\lambda_{\text{Edd}} \sim 10^{-6}$, and develop high- λ_{Edd} tails reaching rare episodes up to $\lambda_{\text{Edd}} \sim 10^{-1}$, as expected in mechanically-dominated AGN feedback (e.g. Russell et al. 2013). Overall, the simulated feeding histories span nearly five orders of magnitude, consistent with the broad, strongly time-variable accretion-rate distributions inferred for AGN populations (e.g. Kauffmann & Heckman 2009; Aird et al. 2012). The low absolute normalization mainly reflects the large SMBH mass adopted in our setup, $M_\bullet = 2.8 \times 10^8 M_\odot$ (Table 1), so that even substantial sink accretion corresponds to modest Eddington-scaled rates. In addition, the relatively short simulated time span, $\lesssim 100 \text{ Myr}$, limits the sampling of very rare near-Eddington events that would be better captured over Gyr-scale duty cycles. Both runs therefore remain predominantly in a low-Eddington, mechanically dominated regime, broadly consistent with radio-mode AGN activity (e.g. Best & Heckman 2012; Heckman & Best 2014), while still showing intermittent excursions toward higher accretion states relevant to wind/feedback-regulated systems.

Despite their similar peak values, the two reference runs differ in the shape of the high- λ_{Edd} tail. In the high-turb run, the distribution is approximately skewed toward large λ_{Edd} , with a logarithmic slope close to -1 over $\lambda_{\text{Edd}} \sim 10^{-4} - 10^{-2}$, reminiscent of the scale-free, flicker-like accretion/feedback variability discussed in micro-scale disk-wind models (Fiore et al. 2024), before steepening to ~ -2 for $\lambda_{\text{Edd}} \gtrsim 10^{-2}$. By contrast, the low-turb run displays a flatter tail, with slopes shallower than -1 across the sampled λ_{Edd} range. This indicates that, although both systems remain in a low-Eddington state on average, the low-turb atmosphere more frequently sustains relatively high Eddington-scaled feeding episodes, consistent with its larger $\dot{M}_\bullet / \dot{M}_{\text{Bondi}}$ and stronger inner inflow (see Figures 1 and 4, in the next section). The flatter PDF in low-turb, which is not seen in the no-jet B26b comparison, suggests that jet-regulated CCA modifies the mapping between turbulence and SMBH feeding. As also suggested by the morphological maps in C26a, weaker ambient stirring plausibly reduces cloud and filament fragmentation, limiting meso-scale decoupling and allowing condensates across a broader mass range to remain connected to the central accretion region.

3.2. Inflow & Outflow

In order to further characterize the CCA cycle, in this section we study the properties of the inflowing and outflowing material as a function of time and radius. Operatively, we classify as *inflow/outflow* each gas cell that has negative/positive radial velocity with respect to the SMBH, respectively. At radius r , we define:

$$\dot{M}_{\text{in}}(r) = \int \rho \max(-v_r, 0) dA,$$

and

$$\dot{M}_{\text{out}}(r) = \int \rho \max(v_r, 0) dA.$$

where dA is the surface element at radius r , ρ the gas density and v_r the radial component of the velocity.

This operational definition measures radial gas circulation in the simulation and does not necessarily coincide with observationally inferred inflow or outflow rates. In observations, the interpretation depends on the tracer phase, emissivity weighting,

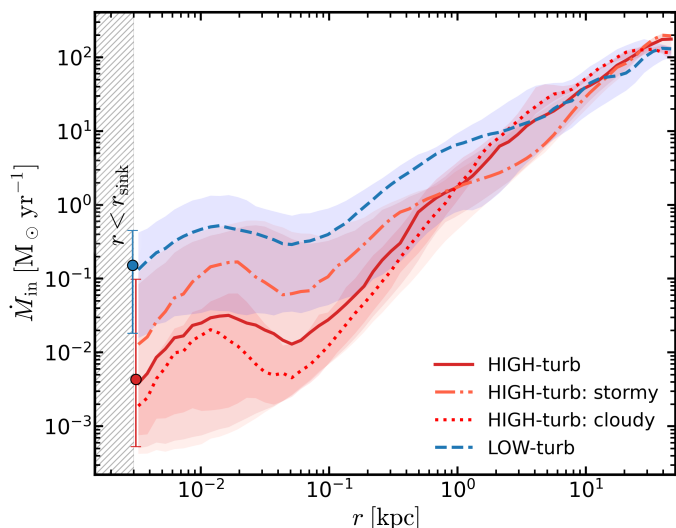


Fig. 4. Inflow rate as a function of the radius in the *low-turb* and *high-turb* case, as indicated in the legend. For the latter, we subdivide the two different weather phases, *stormy* (dash-dotted) and *cloudy* (dotted lines). Lines show the 50th percentile throughout the whole duration of the simulation, with the shaded areas indicating the 1 sigma levels at any given radius (16th and 84th percentiles). For comparison, colored circles at $r = r_{\text{sink}}$ display the time median of \dot{M}_\bullet for the *low-turb* (blue) and *high-turb* (red) cases.

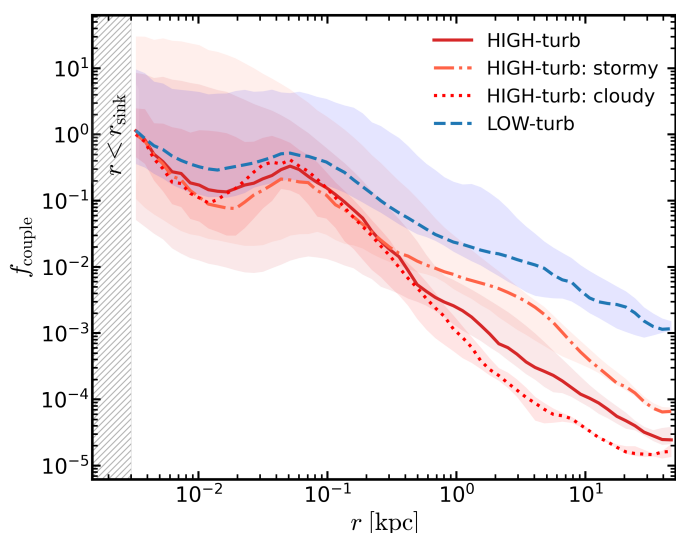


Fig. 5. Inflow coupling efficiency as a function of the radius in the *low-turb* (blue dashed line) and *high-turb* (red solid line) case. As in Figure 4, lines indicate the 50th percentile throughout the total duration of the simulation, with the shaded areas showing, at any given radius, the 1 sigma levels (16th and 84th percentiles).

projection effects, assumed geometry, and line of sight; we discuss these caveats in §4. Since a direct one-to-one mapping is not possible, throughout this section we discuss how these ideal inflows/outflows manifest in simulated data.

3.2.1. Radial accretion properties

Figure 4 shows the time-averaged radial inflow profile (median lines with 1 sigma uncertainty shaded areas) for the two reference runs, together with the *stormy* and *cloudy* decomposition of the *high-turb* case. At the sink radius, for compari-

son, the circles report the time median sink accretion rate as measured in Figure 1 with 1 sigma error bars (calculated over the elapsed time of the simulations). Outside the central kpc, all curves approach similar values and rise with radius following a power law trend (roughly $\propto r^{1.5}$), indicating that the large-scale atmosphere (within 100 kpc) provides a comparable mass supply in both turbulence regimes and it is only marginally affected by the presence of the jet as shown by their comparable slopes at the same radii in B26b. The main differences emerge inside the inner kpc. The *low-turb* run maintains a systematically larger inward mass flux down to the sink scale, consistent with its denser and more persistent central cold reservoir and with its prolonged super-Bondi accretion history. By contrast, the *high-turb* run shows a much weaker inflow in the central $\lesssim 0.1$ kpc, despite comparable supply at larger radii. This implies that the reduction in late evolution of the BH feeding in the *high-turb* atmosphere is not caused by a lack of gas at meso- or macro-scales, but by a reduced efficiency in funnelling the condensed component through the inner transport bottleneck (Fig. 5). At the sink micro-scale, the median SMBH accretion rates are $\sim 4 \times 10^{-3} M_\odot \text{ yr}^{-1}$ for *high-turb* and $\sim 0.15 M_\odot \text{ yr}^{-1}$ for *low-turb*, illustrating the large difference in final coupling to the SMBH.

The *stormy/cloudy* decomposition of the *high-turb* run, shown in the dash-dotted and densely dotted lines, makes this point particularly clear. During the *stormy* phase, the radial inflow profile is enhanced with respect to the *cloudy* phase throughout the inner kpc, showing that precipitation is more directly coupled to SMBH feeding. In the *cloudy* phase, instead, the inflow rate is strongly suppressed at the smallest radii, while remaining comparable at larger radii (see also C26a). This again supports the interpretation introduced above: in the *cloudy* state, cold clouds and filaments still form and survive on micro- and meso-scales, but a smaller fraction of them is able to penetrate all the way to the sink.

To quantify the coupling between large-scale inflow and SMBH feeding, we define a scale-dependent coupling efficiency:

$$f_{\text{couple}}(r) \equiv \frac{\dot{M}_\bullet}{\dot{M}_{\text{in}}(r)}, \quad (16)$$

where $\dot{M}_{\text{in}}(r)$ is the mass inflow rate measured at radius r . This quantity provides a direct measure of how efficiently gas crossing a given scale is able to reach the SMBH. In Figure 5, we compute f_{couple} as a function of the radii and, since the flow is highly time-variable and non-steady, this quantity is interpreted statistically through percentile-based measures. The coupling efficiency declines with radius in both runs, indicating that only a fraction of the inflowing gas at large scales is able to reach the SMBH. This reflects the inherently non-steady and multiphase nature of the flow, where mixing, stirring, and phase transitions progressively reduce the net inward mass transport. A clear systematic offset is observed between the two turbulence regimes. At all radii beyond the innermost region ($r > r_{\text{sink}}$), the *low-turb* run exhibits a significantly higher coupling efficiency than the *high-turb* case. This indicates that, although both atmospheres are able to condense multiphase gas, only the *low-turb* configuration maintains an efficient inward transport toward the SMBH.

The largest divergence occurs at meso-scales ($\sim 0.1 - 1$ kpc), where the *high-turb* run shows a rapid drop in f_{couple} by more than an order of magnitude relative to the *low-turb* case. This demonstrates that strong turbulence driving combined with the

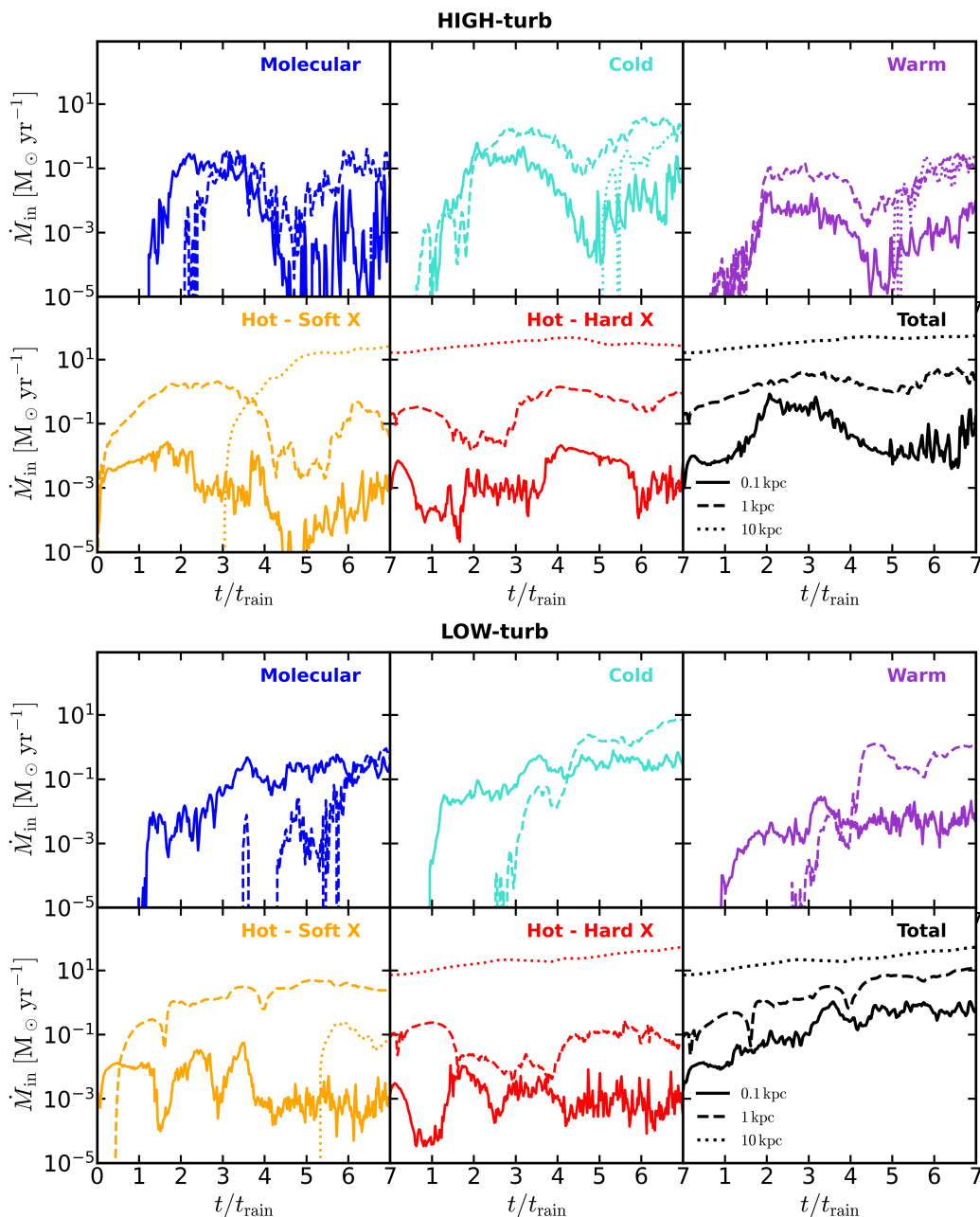


Fig. 6. Phase-separated mass inflow rates for the high-turb (top figure) and low-turb (bottom figure) runs as a function of time normalized to t/t_{rain} . Inflow is defined from gas with negative radial velocity, and each panel shows the corresponding \dot{M}_{in} for a given thermal phase: molecular, cold, warm, soft X-ray, hard X-ray, and the total (black lines). Solid, dashed, and dotted lines refer to $r = 0.1, 1,$ and 10 kpc, respectively. In the high-turb scenario, once condensation is established, the inflow becomes increasingly multiphase on micro- and meso-scales, while the outer halo remains dominated by the hot phases. The low-turb condensed inflow remains more persistent at late times, consistent with a longer-lived central cold reservoir and sustained SMBH feeding.

jet-induced turbulence disrupts the coherent inward cascade, effectively decoupling condensation from SMBH feeding. In terms of the weather framework, this behaviour suggests that, during the transition from a *stormy* to a *cloudy* state, cold gas remains abundant at meso-scales in the high-turb run, it is increasingly mixed, redistributed, or trapped in circulation, and only a small fraction is able to reach the central sink. By contrast, the low-turb run retains a more coherent, *rainy*-like configuration, in which the condensed gas remains dynamically coupled to the SMBH.

The decomposition of the high-turb run into *stormy* and *cloudy* phases clarifies the physical origin of this behaviour. Dur-

ing the *stormy* phase, the coupling efficiency remains relatively high across meso-scales, indicating that extended precipitation is still dynamically connected to the SMBH. In this regime, the larger jet-injected turbulence enhances condensation and promotes intermittent but effective inward transport, consistent with a bursty feeding. In contrast, the *cloudy* phase is characterized by a sharp decline in f_{couple} at the meso-scale ($r \gtrsim 300$ pc), with values suppressed by up to 2 orders of magnitude with respect to the *stormy* phase. This demonstrates that, although cold gas remains abundant, it is no longer efficiently funneled to the centre. At the micro-scale, $r \lesssim 0.1$ kpc, the *stormy* and *cloudy* curves are

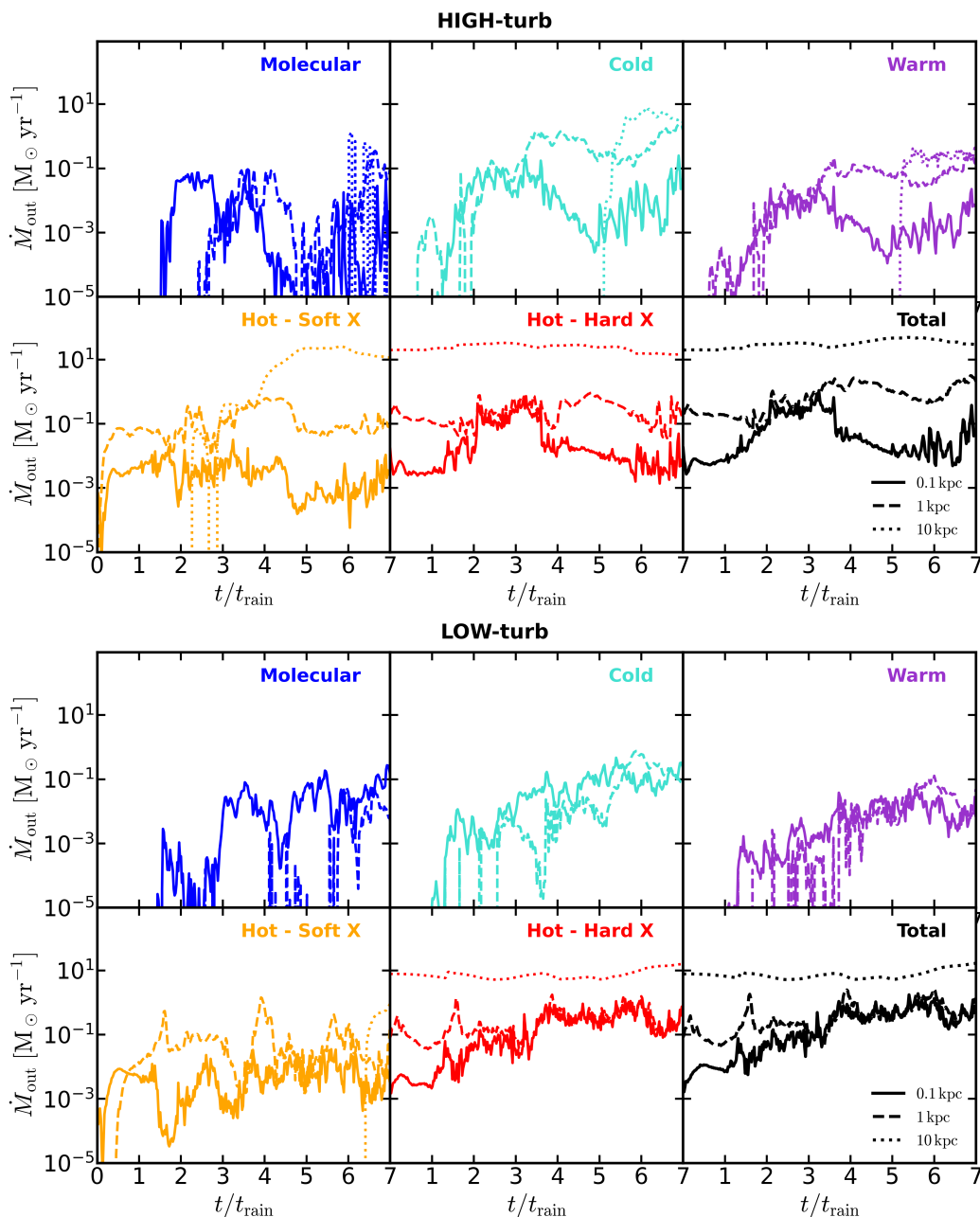


Fig. 7. Phase-separated mass outflow rates for the high-turb (top figure) and low-turb (bottom figure) runs as a function of time normalized to t/t_{rain} . Outflow is defined from gas with positive radial velocity, and each panel shows the corresponding \dot{M}_{out} for a given thermal phase: molecular, cold, warm, soft X-ray, hard X-ray, and the total (black lines). Solid, dashed, and dotted lines refer to $r = 0.1, 1,$ and 10 kpc, respectively. In the high-turb case, the molecular, cold and warm outflows are more extended and intermittent, consistent with stronger uplift and entrainment by the more disturbed medium. In the low-turb atmosphere, cool outflows are mostly confined to the inner kpc and are less intermittent, indicating a lower entrainment/uplifting due to the weaker turbulent stirring efficiency.

more similar because the lower \dot{M}_{\bullet} in the cloudy phase is partly compensated by the reduced inflow reaching the same radii.

We interpret this transition as a shift from a transport-dominated to a mixing-dominated regime. In the *stormy* state, filamentary structures and clouds retain sufficient coherence to sustain inward transport. In the *cloudy* state, the cold phase is redistributed into a population of drifting clouds embedded in a disturbed hot halo. As a result, the multiphase medium becomes dynamically decoupled from SMBH feeding.

These results highlight that the presence of multiphase gas alone is not sufficient to ensure efficient accretion. The key quantity is the transport efficiency across the meso-scale, which

is strongly regulated by the turbulent state of the atmosphere. In this framework, the *stormy-to-cloudy* transition marks the point at which the feedback-regulated system shifts from a precipitation-fed regime to a decoupled multiphase state, where cold gas survives but is no longer effectively accreted by the SMBH.

3.2.2. Time-dependent multiphase transport across scales

The phase-separated fluxes provide a complementary view. Figure 6 displays the inflow rates calculated at different radii representative of the 3 spatial regimes, i.e. micro- (solid, 0.1 kpc),

meso- (dashed, 1 kpc) and macro-scales (dotted, 10 kpc) as labelled in the legend. In observational terms, these ideal inflow/outflow patterns would likely translate into broader, multi-component ionized or molecular profiles in *stormy* phases, and narrower, more coherent cold-gas kinematics in *rainy* phases, although the exact mapping depends on projection and emissivity weighting. On the other hand, large-scale hot outflows are visible primarily as X-ray cavities and shocks, rather than clear line-of-sight (LOS) velocity features in the optical or in molecular lines.

In the *low-turb* run (bottom), the total inflow increases with radius and is dominated at large scales by the hot phases, especially the hard- and soft-X-ray components. This indicates that the outer halo is largely governed by hot, volume-filling circulation. Once condensation is established ($t/t_{\text{rain}} \gtrsim 1$), however, the molecular-temperature, cold, and warm inflow rates at 0.1 and 1 kpc remain sustained ($\gtrsim 0.1 M_{\odot} \text{ yr}^{-1}$) for several rain times. Thus, while the total mass circulation is largest on macro-scales, the condensed component remains efficiently connected from the meso-scale down to the micro-scale, naturally supporting the persistent SMBH feeding typical of the rainy state. This behaviour matches the longer-lived central cold reservoir described in C26a, also typically observed in millimeter emission (e.g. Elford et al. 2024), and naturally explains the persistent super-Bondi \dot{M}_{\bullet} discussed above, typical of the *rainy* weather where the gas feeding the innermost region is progressively transformed from a predominantly hot steady inflow at macro-scales into a multiphase, cold-dominated inflow at micro- and meso-scales.

By contrast, in the *high-turb* run shown in the top graph, the same phases behave more episodically. During the *stormy* weather, the molecular and cold inflow rates at 0.1 kpc rise rapidly and become comparable to those in the *low-turb* ($\sim 0.1 M_{\odot} \text{ yr}^{-1}$), indicating a vigorous precipitation directly feeding the SMBH. However, after $t/t_{\text{rain}} \sim 3.5$, the inner inflow in the condensed phases collapses by 2-3 orders of magnitude, while significant cold and molecular inflow can still be present at meso- and macro-scales. It is important to note that the *cloudy* phase is not completely lacking inflowing condensed gas, but the coupling between meso- and micro-scale precipitation, eventually feeding the SMBH, becomes much weaker. This decoupling corroborates what was suggested by the accretion history in Figure 1 and the thermodynamic phase distributions shown in C26a.

Outflow rates are shown in Figure 7. At $r = 10$ kpc, the outward flux is dominated by the hot phases in both runs, with the hard-X-ray-temperature component providing the largest contribution. This large-scale hot outflow should be interpreted mainly as part of the volume-filling turbulent weather circulation of the halo, further stirred and reorganized by the jet, rather than as purely jet-ejected material. By contrast, the cool phases are more intermittent and primarily concentrated at smaller radii, consistent with uplifted/entrained condensates and local fountain motions rather than a coherent large-scale escaping cold wind.

In the *high-turb* case, the transition from *stormy* to *cloudy* is particularly evident. During the *stormy* phase, up to $t/t_{\text{rain}} \sim 4$, the outflow rates show flicker-like variability across phases, with the molecular-temperature, cold, and warm components becoming increasingly prominent as cooling proceeds. In the innermost shell ($r = 0.1$ kpc), the cool outflow peaks around $t/t_{\text{rain}} \sim 3$, reaching $\sim 0.1 M_{\odot} \text{ yr}^{-1}$. After $t/t_{\text{rain}} \gtrsim 4$, the inner outflow weakens, broadly following the decline in \dot{M}_{\bullet} and the associated reduction in the central jet response. At the same time, cool outward transport becomes more radially extended: molecular-temperature, cold, and warm gas remain visible at the 1 kpc shell, and during the late *cloudy* phase ($t/t_{\text{rain}} \gtrsim 5$) cold and warm ma-

terial can intermittently reach the 10 kpc shell. The cold component can approach rates of order unity to a few $M_{\odot} \text{ yr}^{-1}$, while the warm and molecular-temperature components remain lower and more intermittent. This late stage also shows larger relative variability, of order ~ 1 dex and even higher in the molecular-temperature phase, consistent with a more disturbed and turbulent multiphase medium (see also §3.1.2).

By contrast, the *low-turb* run shows a more locally confined recycling cycle. The molecular-temperature, cold, and warm outflows are strongest within the inner $\lesssim 1$ kpc, with typical rates spanning $\sim 10^{-5} - 0.1 M_{\odot} \text{ yr}^{-1}$. Weak and intermittent cool outward motions can appear at 10 kpc, but they are much less persistent and less massive than in the *high-turb* run. The soft-X-ray-temperature component also moves outward mainly on micro- and meso-scales, reaching the outer radial shell only at later stages ($t/t_{\text{rain}} \gtrsim 6$), qualitatively echoing recent X-ray studies of type-2 quasars in which extended soft-X-ray emission traces ionized gas shaped by AGN-driven outflows and multiphase feedback (e.g. Trindade Falcão et al. 2026). The hard-X-ray-temperature component, instead, is present at all epochs and radii, as expected for the hot phase that carries most of the volume-filling turbulent circulation. At 0.1 and 1 kpc, its outflow rate is more directly coupled to the central feedback response, since the injected jet material has $T_{\text{jet}} = 5 \times 10^8$ K and the injected mass flux scales with the SMBH accretion rate (Eq. 6). At 10 kpc, however, the hot outflow is smoother and remains around $\sim 10 M_{\odot} \text{ yr}^{-1}$, suggesting that jet propagation, turbulent weather circulation, and halo-scale mixing wash out much of the short-timescale variability injected near the centre.

Overall, the mass-flux diagnostics reinforce the interpretation emerging from the thermodynamic and accretion analyses. The large-scale hot supply and volume-filling circulation are broadly comparable in the two runs, but the weather state determines how efficiently condensed gas is transported from meso- to micro-scales. In the *low-turb* run, and during the *stormy* phase of *high-turb*, the condensed phases remain dynamically connected to the central sink and can sustain efficient SMBH feeding. In the late-time *cloudy* phase of *high-turb*, instead, multiphase gas is still present on meso- and inner macro-scales, but a larger fraction is recycled, uplifted, or dispersed through the inner halo rather than being funneled into the SMBH.

3.3. Kinematics – CCA diagnostics

We now focus on two complementary diagnostics rooted in the CCA framework of Gaspari et al. (2018): the k-plot, which traces projected multiphase gas kinematics, and the C-ratio, which links cooling and turbulent mixing. Together, they provide observationally motivated probes of how the simulated weather states regulate meso-scale condensation and feeding.

3.3.1. k-plots

To characterize the kinematics of the multiphase atmosphere, we construct the *kinematical plot* (k-plot) following Gaspari et al. (2018), in which we plot the line-of-sight velocity offset, $\log_{10}|v_{\text{los}} - v_{\text{sys}}|$, versus the line-of-sight velocity dispersion, $\log_{10}(\sigma)$, for beam-based projected samples (to properly mock observational datasets, see also Maccagni et al. 2021). Specifically, we set $v_{\text{sys}} = 0$ km/s (we assume on average a null global bulk velocity of the gas with respect to the origin) and compute mass-weighted values, projected along the z axis (hence the jet axis), for any given gas phase in separate radial ranges accord-

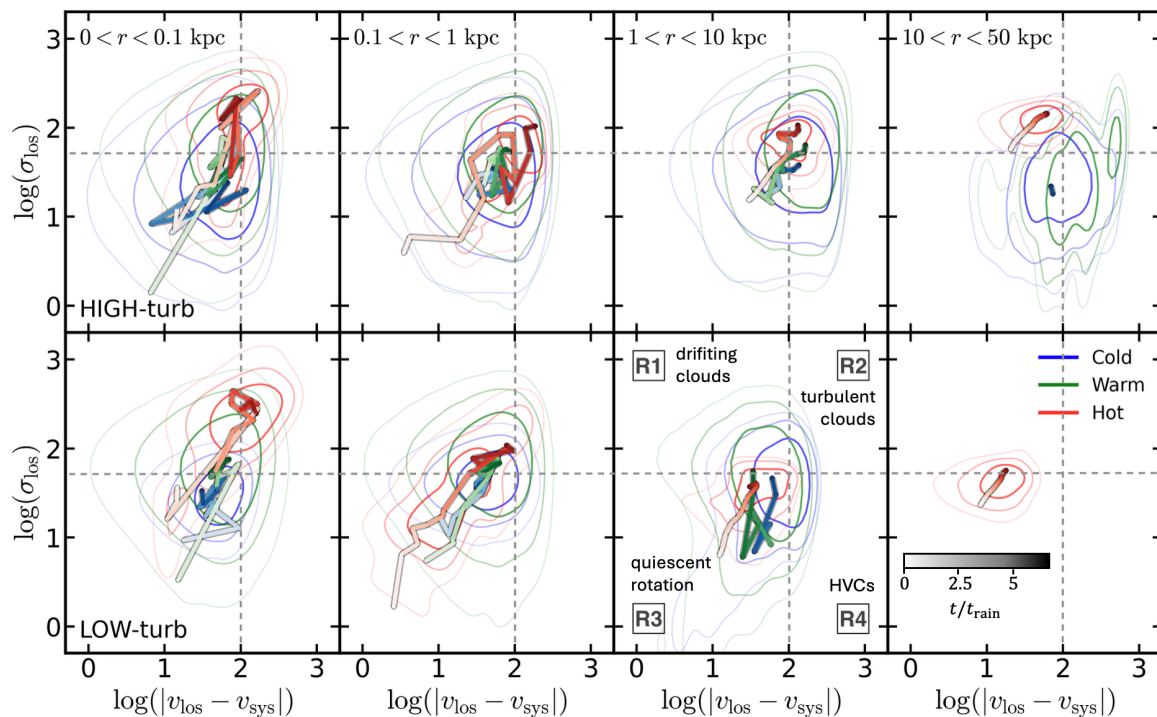


Fig. 8. k-plot: evolution of projected gas kinematics in the k-plot plane, $\log_{10}(|v_{\text{los}} - v_{\text{sys}}|/\text{km s}^{-1})$ versus $\log_{10}(\sigma_{\text{los}}/\text{km s}^{-1})$, shown for four radial shells from micro- to outer macro-scale. The top row shows high-turb, while the bottom row shows low-turb. Blue, green, and red contours show the time-integrated distributions of the cold, warm, and hot phase components, respectively. For visual clarity, the molecular-temperature and cold phases are grouped as “cold”, while the soft- and hard-X-ray-temperature phases are grouped as “hot”. Contours enclose the 85th, 92nd, and 97th percentiles of each phase distribution. Coloured tracks show the temporal evolution of the median position of each phase in the diagram, with darker colours corresponding to later times over the interval $t/t_{\text{rain}} = 0 - 7$. Grey dashed lines mark the reference thresholds $|v_{\text{los}} - v_{\text{sys}}| = 100 \text{ km s}^{-1}$ and $\sigma_{\text{los}} = 50 \text{ km s}^{-1}$, which separate the heuristic kinematic regions discussed in the text.

ing to the definition of micro-, meso-, inner and outer macro-scale. It is important to note that the projected beam size is retrieved by considering a resolution buffer of 256 elements per side, where the projected area is adjusted to each radial shell (i.e., $l_{\text{side}} = 2 \times r_{\text{shell,max}}$). Although the k-plot does not uniquely separate inflow from outflow, it provides an observationally motivated framework for classifying quiescent, rotating, drifting, turbulent, and high-velocity gas components in velocity-dispersion space.

As in B26b, we use the gray dashed lines as heuristic guides and divide the k-plot into four projected kinematic regions: drifting clouds (R1), with modest bulk velocity offsets but enhanced line broadening; strongly turbulent gas (R2), with both large velocity offsets and high dispersion; quasi-quiescent or rotation-dominated gas (R3), with low offsets and low dispersion; and high-velocity clouds (R4), with large projected bulk motion but comparatively lower internal dispersion. The kinematic regimes are defined by $|v_{\text{los}} - v_{\text{sys}}| = 100 \text{ km s}^{-1}$ and $\sigma_{\text{los}} = 50 \text{ km s}^{-1}$. This representation is particularly useful because different regions of the diagram map naturally onto distinct cloud kinematics: the R3 quadrant corresponds to relatively quiescent clouds, the R2 area to turbulent clouds with both large velocity offsets and large internal dispersion, the R1 region to drifting clouds with modest bulk motion but large dispersion, and the R4 section to high-velocity clouds (HVCs) with stronger bulk motion but lower internal turbulence, as also observed in absorption-line studies (e.g. Tremblay et al. 2016). The mutual interplay among the different phases, in terms of overlap and extent in the k-plot, is representative of the physical state of the gas. In particular, during CCA episodes, relatively quiet material is expected to

progressively drift towards a chaotic state, with typical velocities of $50 - 100 \text{ km s}^{-1}$ (Gaspari et al. 2018).

Figure 8 shows how the projected gas kinematics evolve in the k-plot for the two reference runs, separated into four radial shells and three broad thermal components. For visual clarity, we group the molecular-temperature and cold phases into a single cold component, and the soft- and hard-X-ray phases into a single hot component. The percentile contours show the time-integrated distribution of each phase, while the tracks follow the temporal evolution of the median kinematic state from early to late times. This representation allows us to track whether each phase remains kinematically ordered, becomes increasingly turbulent and mixed, or shifts toward regions of larger projected bulk motion. In this way, the k-plot links the simulated weather states to observable multiphase kinematic signatures.

A first common feature of both runs is the radial ordering of the kinematics. In general, all phases at most radii broadly follow a diagonal transition from the quiescent quadrant (R3) to the strong turbulence quadrant (R2). In the inner shells, all phases tend to occupy regions of relatively large velocity dispersion (from 10 km s^{-1} up to $\sim 500 \text{ km s}^{-1}$ in the hot phases), while at the same time the intrinsic velocities increase from a few $\times 10 \text{ km s}^{-1}$ up to $\geq 100 \text{ km s}^{-1}$ reflecting the stronger coupling between turbulence, jet forcing, and multiphase condensation near the centre. Moving outward, the characteristic loci become more compact, and the phase trajectories shorten, indicating that the kinematics becomes progressively less variable and more dominated by the global weather circulation. The outermost macro-scale shell ($10 < r < 50 \text{ kpc}$) is the most stable in both runs and is largely set by the hot component, consistent with

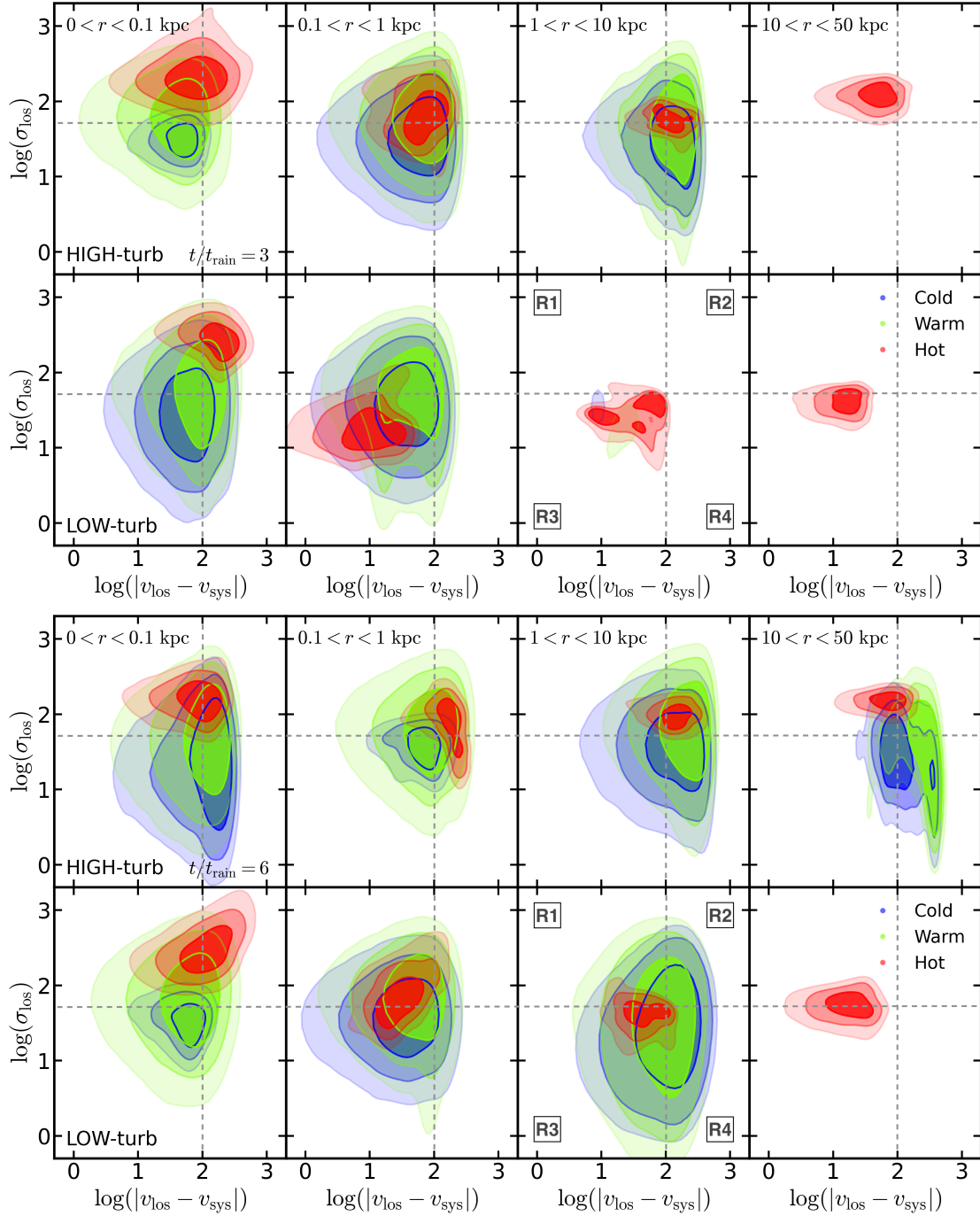


Fig. 9. As in Fig. 8, k-plots of the projected gas in radial bins, from the micro to the outer macro scale, for the *high-turb* and *low-turb* reference runs at two representative evolutionary stages. The upper block shows $t/t_{\text{rain}} = 3$, while the lower block shows $t/t_{\text{rain}} = 6$; within each block, the *high-turb* run is shown in the upper row and the *low-turb* run in the lower row. As in Fig. 8, the molecular-temperature and cold phases are grouped as “cold”, while the soft- and hard-X-ray phases are grouped as “hot”. Solid contours and filled regions show the distribution of each phase in the corresponding radial bin, with hot, warm, and cold gas shown in red, green, and blue, respectively. The filled levels enclose the 85th, 92nd, and 97th percentiles, from lighter to darker shading.

the fact that this region remains thermodynamically dominated by the diffuse hot X-ray atmosphere.

The main difference lies in the degree of kinematic mixing between phases in the inner kpc (see the first two columns). In the *high-turb* run, the trajectories are broader, more irregular, and more strongly overlapping. The cold and warm phases do not remain confined to a narrow locus; instead, they wander substantially in both $|v_{\text{los}} - v_{\text{sys}}|$ and σ_{los} , often approaching the kinematic region occupied by the hot component. As a result, the con-

densed gas randomly samples a wide range of dynamical states, from relatively quiescent clouds (R3) to turbulent clouds (R2), and from high-velocity components (R4) to drifting clouds (R1). This is the kinematic counterpart of the *stormy* multiphase atmosphere identified in C26a: stronger turbulence and a more disturbed jet cocoon continuously stir the condensed gas, enhance its kinematic overlap with the hot phase, and broaden the distribution of cloud velocities and dispersions. The hot phase also evolves rapidly from a more quiescent locus toward a turbulence-

dominated regime, following the diagonal drift described above. Thus, in the *high-turb* run, the atmosphere is not only thermodynamically more mixed, but also kinematically more entangled.

The *low-turb* run instead displays a more coherent evolution. The phase trajectories are more compact and, particularly in the inner two shells, the cold, warm, and hot components remain more clearly separated through time. Their evolution is also smoother, with less pronounced wandering across the velocity-dispersion plane. This behaviour is consistent with the calmer weather state inferred from the morphological and phase-space diagnostics: condensation still occurs and the phases remain dynamically active, but the overall flow is less chaotic and the multiphase gas remains organized enough to sustain a persistent inward reservoir. The kinematic separation between phases is therefore better preserved in *low-turb* than in *high-turb*, indicating a more coherent coupling between precipitation and SMBH feeding.

The difference is especially informative in the inner kpc, where the simulations connect meso-scale precipitation to the central sink. In the *high-turb* case, the time trajectories of the cold and warm phases span a larger range in σ_{los} , with the colored contours covering $\gtrsim 1$ dex at the 97th percentile, and remain closer to the hot component for a substantial fraction of the evolution. This indicates stronger stirring and repeated recycling between inflow, uplift, and local fountain motions. In the *low-turb* case, the same phases occupy tighter loci, with the cold component typically remaining within $\lesssim 100 \text{ km s}^{-1}$ and spanning $\lesssim 0.5$ dex at the 97th percentile. Their evolution also shows smaller excursions through the k-plane, consistent with a more coherent coupling between condensation and inward transport. Thus, the k-plot supports the interpretation already suggested by the \dot{M}_* , PSD, and inflow diagnostics: the *high-turb* run is characterized by a more chaotic, extended, and mixed multiphase circulation, consistent with the *stormy-to-cloudy* weather transition, whereas the *low-turb* run supports a more ordered and persistent condensation cycle, limited to the inner few kpc and consistent with a *rainy* CCA regime.

Figure 8 summarizes the time-integrated distributions and median evolutionary tracks, while Figure 9 shows instantaneous k-plots at $t/t_{\text{rain}} = 3$ and 6. The upper and lower blocks correspond to these two epochs, respectively; within each block, the *high-turb* and *low-turb* runs are shown in the first and second rows. As before, columns show the four radial shells, and colors indicate the cold, warm, and hot phase components. The contours trace the instantaneous phase distributions, with increasing opacity marking the 85th, 92nd, and 97th percentiles. This view highlights how the *stormy*, *cloudy*, and *rainy* weather states populate projected velocity-dispersion space.

At $t/t_{\text{rain}} = 3$, the *high-turb* case occupies a broader and more asymmetric region of the kinematic plane, especially in the inner shell. The cold and warm phases are shifted toward the R1 region, while the hot phase shows systematically larger σ_{los} ($\gtrsim 100 \text{ km s}^{-1}$) and $|v_{\text{los}} - v_{\text{sys}}|$. This suggests that, at the micro-scale, the condensed gas is no longer confined to a narrow quiescent locus, but is being stirred and partially separated kinematically from the hot phase, consistent with the strong central accretion episode shown in Fig. 1. At meso- and inner macro-scales, the three phase components remain more dynamically mixed, indicating that the disturbed state extends beyond the immediate sink region. This is the kinematic imprint of the *stormy* phase: the multiphase gas is strongly stirred, broadly mixed, and frequently displaced toward the turbulent-cloud and drifting-cloud regimes. By contrast, the *low-turb* run at the same epoch retains a larger fraction of its cold and warm gas in the lower-

left region of the plane (R3) at micro- and meso-scales. This is consistent with a population of more quiescent clouds that can remain coherently connected to inward precipitation and therefore contribute to the sustained \dot{M}_* . At the micro-scale, the hot component still shows large σ_{los} and $|v_{\text{los}} - v_{\text{sys}}|$ values, both $\gtrsim 100 \text{ km s}^{-1}$, but it shifts rapidly toward the R3 region at larger radii. The *rainy* phase is therefore characterized by a more compact and coherent kinematic evolution, in which multiphase structures build up gradually and remain better organized across scales. In this sense, the high-turbulence atmosphere is already in a *stormy*-dominated regime by $t/t_{\text{rain}} = 3$, whereas the low-turbulence atmosphere is still developing toward a full *rainy* weather state, consistent with the projections shown in C26a.

By $t/t_{\text{rain}} = 6$, the difference becomes even more diagnostic. In the *high-turb* run, the micro-scale cold and warm phases have moved away from the earlier *stormy* configuration and now occupy a broader region of the k-plot, extending mainly toward the R2 and R4 quadrants. Their σ_{los} distribution spans up to ~ 2 dex around $\sim 50 \text{ km s}^{-1}$, while $|v_{\text{los}} - v_{\text{sys}}|$ reaches $\sim 200\text{--}300 \text{ km s}^{-1}$. This indicates that a substantial fraction of the condensed gas is now more dispersed, turbulent, and characterized by large projected velocity offsets. Rather than forming a coherent inward reservoir, the cold and warm material is increasingly stirred and recycled, consistent with the weaker central accretion and reduced inner inflow shown in Figs. 1 and 4. At the same time, the hot locus changes only mildly with respect to $t/t_{\text{rain}} = 3$, suggesting that the hot phase has already reached a strongly stirred, quasi-steady kinematic state by this stage. At meso-scales, the cold and warm phases remain broadly similar in centroid position to their earlier configuration, although the contour area is reduced, consistent with less efficient late-time production and/or survival of centrally connected multiphase gas. At macro-scales, cold and warm gas have reached the outer shell with relatively large $|v_{\text{los}} - v_{\text{sys}}| \sim 100\text{--}300 \text{ km s}^{-1}$ and modest-to-intermediate dispersions of a few tens to $\sim 100 \text{ km s}^{-1}$, in agreement with C26a, where late-time cold filaments extend to $\sim 10\text{--}20$ kpc. Overall, this supports the picture developed from the accretion and inflow diagnostics: at $t/t_{\text{rain}} = 6$, clouds and filaments are still present, but their coupling to direct SMBH feeding is weaker, and a larger fraction of the condensed gas participates in a turbulent, cloudy inner-halo circulation rather than in coherent inward transport.

The *low-turb* run instead follows a more stable and coherent evolution. Between $t/t_{\text{rain}} = 3$ and 6, the micro-scale phases show only modest changes, with the cold component becoming more compact and less mixed with the warmer and hotter gas while remaining relatively quiescent. This behaviour is consistent with the sustained accretion rate \dot{M}_* , since the condensed gas remains organized into a central reservoir rather than being dispersed into a broader turbulent circulation. At the meso-scale, the hot phase gradually shifts along the diagonal toward larger velocity dispersions and velocity offsets, reflecting the continued action of feedback and ambient turbulence. The cold and warm phases, however, move only mildly in the k-plot, indicating that the multiphase gas largely preserves its coherent kinematic structure at this scale. This points to a gradual transition from quiescent clouds to moderately stirred condensates, as expected for a *rainy* weather state. Because the spatial growth of the rain cycle is slower and more ordered than in the *high-turb* run, cold and warm components appear at the inner macro-scale only at later times, where they show broader scatter in both σ_{los} and $|v_{\text{los}} - v_{\text{sys}}|$. Even there, the phase loci remain more compact and show smaller excursions through the plane than in

the **high-turb** case, indicating a more coherent kinematic evolution. Overall, the **low-turb** run is consistent with a calmer weather cycle in which condensed gas is progressively stirred while remaining comparatively well coupled to inward transport.

An interesting point of comparison can be drawn against the recent simulations shown in B26a,b within the same framework, where no-jet runs investigate the impact of turbulence driving and radiative cooling alone on CCA. Commonly to both studies, we see that the *stormy* phase is usually well represented by a broad and mixed multiphase gas population across the **high-turb** k-plane while the *rainy* weather is associated with more compact and kinematically distinct loci consistent with a more coherent dynamics. However, the k-plots shown in B26b (Figure 9) depict a situation where the phase distribution and the stochastic motion of gas condensates are partially slimmed down with respect to this work. This suggests that, despite the large-scale turbulent stirring provides sufficient kinematic support to sustain a multiphase CCA state, the jet-driven recycling and uplift added here are able to further amplify and favor multiphase mixing and gas circulation, resulting in broader and more overlapping contours highlighted in the k-plot diagnostic.

The two kinematic snapshots reinforce the broader interpretation of the simulations. The **high-turb** run undergoes a genuine weather transition, from a *stormy* phase dominated by strongly turbulent condensates to a *cloudy* phase in which cold and warm gas remain present but become more dispersed, drifting, and weakly connected to direct SMBH feeding. The **low-turb** run does not show such an abrupt change of state; instead, it evolves more smoothly from quiescent clouds toward a moderately turbulent but still comparatively ordered multiphase flow. The k-plot therefore provides an observationally motivated diagnostic for distinguishing different modes of jet-regulated CCA: *stormy* systems should occupy broader and more mixed regions of the velocity-dispersion plane, *cloudy* systems should show extended but weakly coupled drifting condensates, and *rainy* systems should retain more compact and coherent cold/warm loci.

3.3.2. C -ratios

We finally consider the ratio between the local cooling time and the eddy gyration time, typically referred to as C -ratio. This quantity compares the timescale on which gas can radiatively lose thermal energy with the timescale on which turbulence creates/mixes gas perturbations. Following Gaspari et al. (2018), the dimensionless condensation ratio is

$$C \equiv \frac{t_{\text{cool}}}{t_{\text{eddy}}} \quad (17)$$

where the cooling time can be approximated as

$$t_{\text{cool}} \simeq \frac{3k_B T}{n_e \Lambda}, \quad (18)$$

with T being the gas temperature, n_e the electron density, and $\Lambda(T, Z)$ the emissivity function. The eddy gyration/turnover time quantifies how quickly subsonic turbulence stirs the gas on a given scale, and it can be estimated via

$$t_{\text{eddy}} = 2\pi \frac{r^{2/3} L^{1/3}}{\sigma_{v,L}}, \quad (19)$$

where L is the turbulence injection scale and $\sigma_{v,L}$ is the velocity dispersion measured at that scale. This expression assumes a

Kolmogorov cascade, $\sigma_v(\ell) \propto \ell^{1/3}$ (Kolmogorov 1941), as found for subsonic turbulence in simulations (e.g. Gaspari & Churazov 2013; Fournier et al. 2025). At radius r , the characteristic eddy size is expected to be of order r (larger eddies cannot be contained locally), hence $\ell \sim r$.³

This quantity provides a direct measure of the propensity of the atmosphere to undergo turbulence-regulated condensation. Large values, $C \gg 1$, indicate that the cooling time is much longer than the eddy turnover time. In this regime, turbulent stirring and mixing act faster than radiative losses, efficiently dispersing overdense perturbations before they can cool nonlinearly. By contrast, $C \ll 1$ implies that cooling proceeds much faster than turbulent mixing, so the flow tends toward a more coherent cooling-flow regime rather than a turbulence-regulated CCA cascade. The CCA regime is therefore expected to occur when C approaches unity, within the typical ~ 0.3 dex scatter around unity (i.e. $t_{\text{cool}} \sim t_{\text{eddy}}$; see also B26b, Gaspari et al. 2018). In this regime, turbulence can seed nonlinear density perturbations while radiative cooling amplifies them into multiphase clouds and filaments. Such condensation can persist even in the presence of AGN feedback, provided the jet does not fully erase the local balance between cooling and turbulent mixing, as suggested by observations of galaxies, groups, and clusters across different spectral bands (e.g. Juráňová et al. 2020; Maccagni et al. 2021; Temi et al. 2022; Olivares et al. 2022; Singha et al. 2023; Mehdipour et al. 2023; Lepore et al. 2025).

We split the volume into two regions using a bipolar cone of half-opening angle 15° centered on the fixed jet axis. This cone is a diagnostic mask, not the numerical injection geometry: the injected jet remains collimated along $\pm z$ without an imposed opening angle. It is therefore not meant to track jet-launched or entrained material directly, but only to separate the jet-channel/interface region from the off-axis ambient atmosphere. We stress that, in the present analysis, the C -ratio is computed for all temperature phases using the eddy turnover time inferred from the hot turbulent cascade, $t_{\text{eddy,hot}}$. This choice is physically motivated because the C -ratio criterion is fundamentally a hot-phase diagnostic: it identifies where turbulence in the volume-filling atmosphere allows radiative cooling to amplify perturbations into multiphase condensation. It should therefore not be interpreted as a direct measurement of the internal eddy time of already-condensed filaments or clouds. For the cooler phases, the plotted C -ratio values should instead be regarded as phase-tracking diagnostics relative to the ambient hot turbulent field.⁴

Figure 10 compares the radial profiles of the C -ratio across emission-band proxies for two reference runs, **high-turb** (left panels) and **low-turb** (right panels), while the outside/inside jet-cone regions are depicted in the upper/lower figures, respectively. Different line styles and colors indicate the various gas phases, while the color intensity traces the progression of the simulation in time as a function of t/t_{rain} . Vertical lines and related shaded areas show the average spatial extent in radius of each phase with their 1 sigma uncertainty over the range of snap-

³ Given the large dynamic SMR range, estimating the turbulent amplitude from $\sigma_v(r)$ (capped at the driving scale L) measured in spherical radial shells and $L \sim r$ is a reasonable numerical approximation.

⁴ A fully phase-dependent estimate would require measuring the characteristic size and internal velocity dispersion of individual cold clouds and filaments. Since these structures become progressively smaller and more fragmented at lower temperatures, using the large-scale hot-phase eddy time may underestimate the intrinsic C -ratio of the cooler gas. A clump-by-clump analysis of this kind is beyond the scope of the present work and will be investigated in a future study.

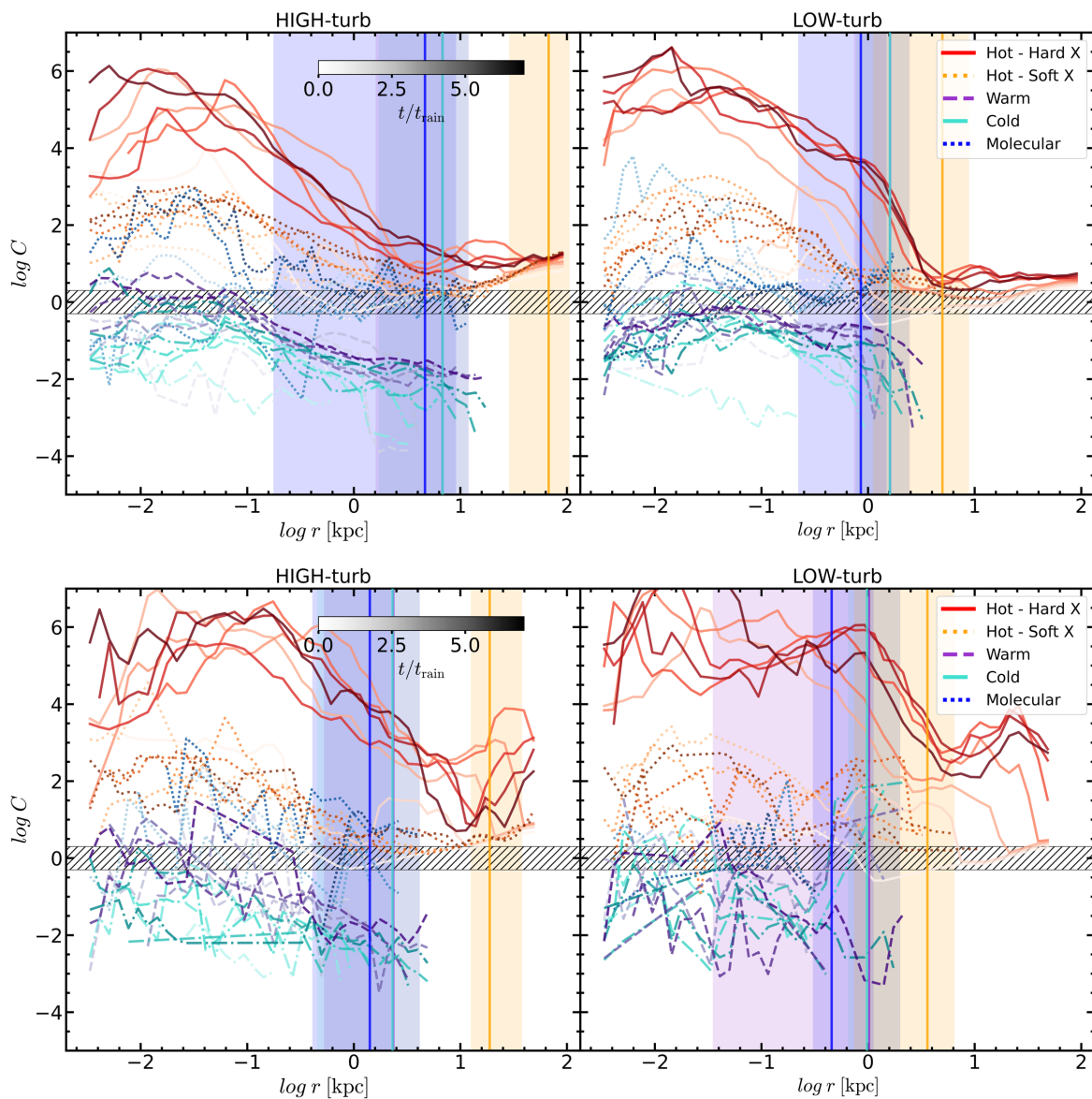


Fig. 10. Radial evolution of the C -ratio for different temperature-defined emission-band proxies, comparing the high-turb (left panels) and low-turb (right panels) runs. Curves with the same color and line style refer to the same gas phase, as indicated in the legend. Within each phase, lighter to darker shades encode increasing evolutionary time in units of t/t_{rain} , as shown by the color bar. Vertical lines and shaded bands mark the mean maximum radial extent of each phase and its scatter over the analyzed snapshots. The grey shaded region, $0.5 < C < 2$, indicates the approximate condensation-prone regime where $t_{\text{cool}} \sim t_{\text{eddy}}$. Upper panels show gas outside the bipolar jet cone, defined by a 15° half-opening angle around the fixed jet axis, while lower panels show gas inside the cone. The hard-X-ray phase reaches the largest central C -ratio values and shows the steepest radial decline. The soft-X-ray and warm phases most clearly trace the approach to the condensation-prone band, while the cold and molecular-temperature phases should be interpreted as already-condensed gas tracked relative to the hot turbulent cascade. For the cooler phases, the plotted C -ratio should be interpreted relative to the hot-phase turbulent cascade, rather than as a direct measure of the internal eddy time of individual clouds or filaments.

shots considered in the plot. Figure 10 can be read in three complementary ways. First, the vertical position of each curve distinguishes mixing-dominated gas, $C \gg 1$, condensation-prone gas, $C \sim 1$, and rapidly cooling gas, $C \ll 1$.⁵ Second, the radial interval near $C \sim 1$ identifies where CCA rain can be sustained. Third, the inside/outside cone comparison traces feedback anisotropy: the hot jet channel remains at high C -ratio values, while the sur-

⁵ For the already-condensed cold and molecular-temperature phases, however, values below $C \sim 1$ should not be read as a separate cooling-flow criterion, but as the location of post-condensation gas relative to the hot-phase eddy time.

rounding atmosphere more readily approaches the CCA threshold.

We start by analyzing the C -ratio outside the jet cone, shown in the upper panels. Similarly to B26a, the hard-X-ray phase in the high-turb run shows relatively low values, $C \sim 1 - 10$, at radii larger than a few kpc, while the C -ratio rises steeply toward the centre, reaching $C \sim 10^4 - 10^6$ at $r < 1$ kpc. This radial trend is mainly driven by the rapid decrease of t_{eddy} toward small radii, whereas the cooling time of the hard-X-ray gas varies more mildly. The hard-X-ray phase therefore approaches the condensation-prone band, $C \approx 0.5 - 2$, primarily at $r \sim 1 - 10$ kpc during the early stages of the run, $t/t_{\text{rain}} \sim 1 - 2$. This ra-

dial interval marks the region where the hot atmosphere is most susceptible to extended nonlinear condensation, consistent with the broad multiphase fragmentation characteristic of the *stormy* weather state identified in C26a. At later times, the remaining hard-X-ray gas is displaced toward larger C -ratio values, $C \gtrsim 10$, indicating that turbulent stirring and mixing dominate over radiative cooling in this phase. The soft-X-ray phase, by contrast, lies systematically closer to the condensation-prone regime, with C -ratio values lower by about 2 – 3 dex, especially at micro- and meso-scales. This suggests that the soft-X-ray-temperature gas acts as the intermediate cooling stage between the volume-filling hot atmosphere and the warmer and colder condensates. Its proximity to $C \sim 1$ over a broad radial range marks the soft-X-ray phase as the main thermodynamic gateway between the volume-filling hot atmosphere and the warmer and colder condensates. Consistently, the maximum radial extensions shown by the vertical bands reach a few tens of kpc for the soft-X-ray phase and $\lesssim 10$ kpc for the warm, cold, and molecular-temperature components.

In the low-turb run, the hard-X-ray phase follows a qualitatively similar radial behaviour, but the transition from the outer, flatter profile to the inner rising profile is steeper, and the C -ratio remains closer to unity at $r \gtrsim 5$ kpc. This indicates that condensation-prone conditions are reached more coherently, although over a more compact radial region, as traced by the smaller radial extent of the soft-X-ray, warm, cold, and molecular-temperature bands. Thus, reaching $C \sim 1$ in the hot phase is a necessary indicator of precipitation-prone gas, but the radial extent of the colder phases also depends on the amplitude of perturbations and the ability of condensates to survive and remain connected across scales. This behaviour is consistent with B26b, where no jet was included, suggesting that outside the directly affected jet cone the large-scale CCA condition is still primarily regulated by the ambient turbulent cascade and cooling structure. Following the condensation cascade toward cooler phases, we find that the average radial extent of the warm, cold, and molecular-temperature components is more confined than in the high-turb run. In particular, the cold and molecular-temperature phases remain mostly within the inner kpc, supporting the interpretation of a compact and coherent *rainy* weather state, as also found in B26b. The molecular-temperature C -ratio profiles provide a particularly useful imprint of this weather transition. In the late-time high-turb run, the molecular-temperature profiles show a marked increase toward the centre, especially below 0.1 kpc. This is driven by the lower central cold-gas density in the *cloudy* phase relative to the earlier *stormy* phase. For approximately fixed temperature and cooling function, $t_{\text{cool}} \propto 1/n$, so a density decrease lengthens the cooling time and increases the C -ratio. By contrast, the low-turb run does not show this central rise. Instead, it evolves toward higher molecular-temperature gas densities, keeping the inner C -ratio closer to $C \sim 1$ and favoring a persistent, reservoir-fed *rainy* inflow.

By contrast, inside the jet-cone, shown in the lower panels, the hard-X-ray phase remains at systematically high C -ratio values, $C \sim 10^2 - 10^6$, in both the high-turb and low-turb runs, especially at small radii $r \lesssim 1$ kpc. This indicates that the jet channel is comparatively hostile to sustained in-situ condensation. The kinetic outflow excavates a hot, low-density, dynamically stirred region in which turbulent mixing and jet heating dominate over local radiative cooling. The gas outside the cone, less directly affected by the injected jet power, is therefore more likely to approach the multiphase rain threshold. The warm and cold phases inside the cone nevertheless show broader tem-

poral variability, including intermittent excursions toward the condensation-prone regime, $C \sim 1$. These episodes likely trace transient cooling, uplift, shocks, and mixing layers driven by jet-induced instabilities. They may seed local and short-lived condensation events, but they do not indicate that the jet channel is the dominant site of long-lived precipitation. Finally, inside the cone the high-turb and low-turb runs show no strong systematic difference within the intrinsic scatter of the C -ratio profiles. This suggests that, close to the jet axis, jet-driven stirring and heating dominate over the imposed ambient turbulent driving, reducing the contrast between the two turbulence amplitudes.

One robust result is the strong geometrical dependence of the condensation criterion. Outside the jet cone, the warm, cold, and molecular-temperature phases remain closer to the condensation-prone regime, $C \sim 1$, over an extended radial interval, especially from the meso-scale to the inner macro-scale. This indicates that sustained precipitation is favoured primarily in the ambient stirred atmosphere and jet-interface regions, rather than along the jet channel itself, as confirmed also in the simulations with the inclusion of the SMBH spin (P26b). This picture is consistent with the multiphase morphology and inflow diagnostics, where cold gas can form and circulate outside the directly heated jet cone before part of it is transported inward. This differs from the pure feeding simulations of B26b, where the cold and molecular phases remain much more cooling-dominated, with typical molecular values $C \sim 10^{-2} - 10^{-3}$. The comparison suggests that jet uplift and feedback-driven mixing/reheating keep part of the condensed gas closer to the hot turbulent atmosphere. For the hot phases, the outside-cone behaviour broadly agrees with the no-jet results of B26a,b, suggesting that the large-scale cooling-versus-turbulence balance remains primarily controlled by the ambient turbulent cascade. The cooler phases, however, show C -ratio values larger by about ~ 1 dex, especially at $t/t_{\text{rain}} \gtrsim 5$. This likely reflects the additional stirring and velocity dispersion induced by the jet as its energy couples to the surrounding atmosphere. Thus, even outside the nominal jet cone, feedback can broaden the turbulent cascade and make the condensation-prone regime more intermittent, particularly in the late-time cloudy state.

The inside/outside comparison is particularly informative for the high-turb case. During the *stormy* stage, the region outside the cone is driven rapidly toward the condensation-prone regime, $C \sim 1$, over a substantial radial interval, matching the epoch of strongest precipitation and SMBH accretion. At later times, as the system enters the *cloudy* phase, the relevant phases still remain close to the condensation band outside the cone, but in a less coherent and more temporally variable way. This indicates that clouds can continue to form, while their coupling to direct SMBH feeding becomes weaker. Conversely, inside the cone the hot phases remain displaced toward high C -ratio values, confirming that the jet channel itself is not the primary site of sustained condensation.

Overall, the two turbulence regimes differ mainly in how broadly, coherently, and persistently they approach the condensation-prone regime. In the high-turb run, the C -ratio profiles show larger temporal excursions and a broader radial spread, especially outside the cone and for the warm and cold phases. This reflects the stormier atmospheric conditions already identified through the kinematic and accretion diagnostics: stronger stirring broadens the region over which gas can be driven toward $C \sim 1$, but also makes that condition more intermittent and irregular. The low-turb run instead displays a smoother and more coherent evolution. Outside the cone, the

cooler phases approach $C \sim 1$ more steadily and over a more compact radial range, while inside the cone the separation between the hot and condensed gas remains cleaner. This behaviour is consistent with a calmer *rainy* weather state, in which condensation is less explosive but more persistently maintained.

Taken together, the C -ratio diagnostics reinforce the main physical picture emerging from C26a and B26a,b. Multiphase condensation is favoured where turbulent stirring and radiative cooling become comparable, $t_{\text{cool}} \sim t_{\text{eddy}}$, and this condition is reached most naturally in the ambient atmosphere surrounding the jet rather than within the jet channel itself. The main jet-regulated effect is therefore geometrical: feedback displaces sustained precipitation away from the hot jet channel and into the surrounding turbulent atmosphere and interface layers. Outside the cone, the *high-turb* run reaches the condensation-prone regime over a broader and more intermittent radial interval, consistent with an extended *stormy* precipitation cycle followed by a weaker *cloudy* coupling phase. The *low-turb* run approaches the same regime more coherently and over a more compact region, supporting a persistent, reservoir-fed *rainy* state. Inside the cone, cold and warm gas can still show intermittent excursions toward $C \sim 1$, likely associated with jet-driven mixing layers, shocks, and local instabilities. These episodes should be interpreted as transient interface events, not as evidence that the jet channel is the dominant site of long-lived CCA rain.

4. Discussion

This work assesses how the kinematic and temporal diagnostics of CCA are modified when radiative cooling and driven turbulence are coupled to an actively responding kinetic jet. Together with the companion thermodynamic analysis presented in C26a, our results show that the jet does not simply oppose cooling through local heating, nor merely advect pre-existing condensates. Instead, it dynamically reorganizes the multiphase medium, changing where gas condenses, how it moves across scales, and how efficiently it remains connected to SMBH feeding.

The two turbulence regimes, *high-turb* and *low-turb*, differ primarily in their inner transport efficiency. Outside the central few hundred parsecs, both runs maintain substantial mass circulation, showing that the halo continues to supply gas to the inner atmosphere. The decisive difference appears closer to the sink. In the *low-turb* case, condensed gas remains well coupled from meso- to micro-scales and sustains a long-lived, super-Bondi accretion state. In the *high-turb* case, late-time inner inflow is strongly suppressed despite the continued presence of multiphase gas at larger radii. Strong ambient turbulence therefore does not shut off condensation altogether; it weakens the final coupling between condensation and SMBH feeding.

This behaviour maps naturally onto the weather states introduced in C26a. The early *high-turb* evolution corresponds to a *stormy* phase, with spatially extended precipitation, strong phase mixing, bursty super-Bondi accretion, and rapidly multiphase inflow/outflow channels. At later times, the same run evolves toward a *cloudy* state: clouds and filaments persist across meso- and inner macro-scales, but are less coherently connected to inward transport. The accretion-rate PSD flattens at low frequencies, and the SMBH feeding rate drops by more than an order of magnitude relative to the preceding *stormy* episode. The *low-turb* run instead maintains a smoother *rainy* mode, in which the central cold reservoir survives longer and sustains the feedback cycle.

The variability, kinematic, and C -ratio diagnostics reinforce this picture. Both runs show flicker-like low-frequency variability and steep high-frequency red-noise tails, consistent with a scale-coupled multiphase feeding cycle rather than stochastic white noise. However, the *low-turb* case retains more variability power over long intervals because its central reservoir remains connected to the SMBH, whereas the *high-turb* case transitions from a bursty *stormy* episode to a weaker *cloudy* regime with reduced long-timescale coherence. The k-plot provides the corresponding kinematic view: the *high-turb* run evolves from strongly turbulent condensates to more dispersed and drifting clouds, while the *low-turb* run evolves more monotonically from quiescent to moderately stirred multiphase gas. The C -ratio profiles show that condensation-prone conditions are preferentially reached in the ambient atmosphere surrounding the jet channel, rather than inside the jet cone itself. In the *high-turb* run this regime spans a broader radial range, whereas in the *low-turb* run it remains more centrally concentrated. These diagnostics identify the meso-scale as the mediating layer of jet-regulated CCA, where turbulence, cooling, uplift, and phase conversion determine whether precipitation becomes extended and weakly coupled, or compact and efficiently feeding.

An instructive comparison is provided by P26a,b, where the jet direction is coupled to the evolving SMBH spin rather than kept fixed. Even modest reorientations, of order a few degrees, allow the jet-driven stirring to be distributed over a larger effective volume, making the transition toward a *sunny* state more persistent and more evident than in the fixed-axis runs analyzed here. In P26b, this broader coupling more efficiently depletes cold gas from the central and meso-scale regions, producing a clearer *stormy-to-sunny* evolution. By contrast, in the fixed-axis *high-turb* run presented here, cold gas is never fully erased: cold inflow and outflow remain visible down to micro- and meso-scales even during the late *cloudy* phase, although their coupling to SMBH feeding is strongly reduced. The radius–time maps in Appendix B and Fig. B.1 make this distinction explicit. A *cloudy* state retains multiphase gas with inefficient central accretion, whereas the *sunny* state reported in P26b corresponds to a more sustained clearing of the cold component from the inner transport region.

Several limitations should be taken into account. The present work analyzes two benchmark turbulence amplitudes and adopts a fixed-axis, purely kinetic, mass-loaded jet. The simulations neglect magnetic fields, anisotropic conduction and viscosity, cosmic rays, star formation, stellar feedback, detailed molecular chemistry, and relativistic jet propagation. For instance, allowing a reorientation of the jet axis would certainly introduce an extra turbulence and mixing component, especially where magnetic pressure and reconnection are expected to play a key role in shaping the gas environmental conditions, hence affecting multiphase condensation and CCA precipitation.

A potential caveat of the present setup is the adoption of a purely kinetic jet, neglecting explicit thermal, magnetic, or cosmic-ray feedback channels at injection. In realistic AGN environments, the relative partition between these components is uncertain and likely varies with the accretion state and micro-physical conditions. However, our choice should also be viewed as a controlled and physically motivated simplification. By injecting feedback purely in kinetic form, we avoid imposing an a priori subgrid prescription for thermal energy deposition, allowing thermalization to arise self-consistently through shocks, turbulent dissipation, and mixing in the resolved flow. This approach enables us to isolate the role of kinetic energy in regu-

lating multiphase condensation and gas transport, and to assess how efficiently mechanical feedback alone can couple to the atmosphere across scales.

In addition, unresolved disk-to-horizon physics is captured through fixed subgrid efficiencies, and the k-plot diagnostics are not yet full emissivity-weighted mock observations including projection and instrumental effects. These choices isolate how ambient turbulence modifies jet-regulated CCA, but likely underestimate the diversity of coupling pathways in realistic systems.

4.1. Comparison with observed AGN feeding and feedback

The diagnostics introduced here can be compared with multiphase observations of AGN host galaxies, groups, and clusters, although they are not yet forward-modeled observables. Our coldest bin should be interpreted as a molecular-temperature proxy, not as a direct prediction of CO luminosity or molecular mass, and the warm and X-ray phases are not emissivity-weighted mock $H\alpha$, $[N II]$, or X-ray maps. With this caveat, the most useful comparisons are filament extent, projected velocity offsets, velocity dispersions, hot-gas turbulent amplitudes, accretion-state indicators, and condensation thresholds. Table A.1 summarizes a proposed mapping between simulation diagnostics and observational proxies. It should be read as a semi-quantitative guide rather than a one-to-one forward model. A direct comparison with individual systems will require synthetic CO, optical-line, and X-ray diagnostics, including emissivity weighting, projection, instrumental resolution, and tracer selection; this will be investigated in future work.

A recurring result from ALMA, MUSE, and X-ray studies is that cold molecular and warm ionized filaments are spatially associated with X-ray cavities, rims, and uplifted low-entropy gas (e.g. Tremblay et al. 2016, 2018; Temi et al. 2018; Vantyghem et al. 2019; Olivares et al. 2019, 2022, 2023, 2025). Observed structures span compact central reservoirs of $\lesssim 3$ kpc in some group and early-type systems, while cool-core cluster filaments commonly extend over $\sim 3 - 25$ kpc. Their projected velocities are typically $|v_{\text{los}} - v_{\text{sys}}| \sim 100 - 400 \text{ km s}^{-1}$, often below the free-fall expectation (Tremblay et al. 2018; Russell et al. 2019; Vantyghem et al. 2019; Olivares et al. 2019). This is broadly consistent with the jet-regulated CCA picture inferred here: the jet does not simply remove or heat the condensing gas; it excavates a hot channel and reorganizes the surrounding atmosphere, where cooling, entrainment, fallback, and feeding coexist.

In this context, the high-turb and low-turb runs map onto different parts of the observed multiphase phenomenology. The high-turb run resembles extended filamentary systems: during the *stormy* and *cloudy* phases, cold and warm gas reach the inner macro-scale and occupy kinematic loci with $|v_{\text{los}} - v_{\text{sys}}| \sim 100 - 300 \text{ km s}^{-1}$ and moderate dispersions, similar to uplifted or slowly raining filaments. The low-turb run instead remains more centrally concentrated, with molecular-temperature and cold gas predominantly retained in the inner kpc and only a slower extension to larger radii. Group-centred systems are especially relevant for the present setup. In NGC 5846, NGC 4636, and NGC 5044, ALMA CO(2-1) observations reveal off-centre molecular clouds with broad linewidths and irregular velocity structure, interpreted as unbound molecular associations drifting through the turbulent hot atmosphere rather than as a settled rotating disk (Temi et al. 2018, 2026). This provides a closer kinematic analogue to our cold-gas transport than a comparison based on molecular mass.

Warm ionized gas provides a complementary kinematic comparison. Optical/IFU and multiphase studies of radio galaxies, brightest group/cluster galaxies, and nearby AGN hosts show complex velocity fields, moderate line-of-sight offsets, and substantial broadening, indicating that the gas is often neither a settled disk nor a purely radial wind (Maccagni et al. 2021; Olivares et al. 2022; Singha et al. 2023; Eskenasy et al. 2024, 2026). This is the regime targeted by the k-plot. The thresholds used here, $|v_{\text{los}} - v_{\text{sys}}| = 100 \text{ km s}^{-1}$ and $\sigma_{\text{los}} = 50 \text{ km s}^{-1}$, lie within the observed multiphase range (Gaspari et al. 2018; Maccagni et al. 2021). In high-turb, cold and warm gas repeatedly cross these thresholds, especially in the inner and meso-scale regions. In low-turb, the phase loci remain more compact. Thus, cold or warm gas near cavities or radio structures should not be classified from morphology alone as inflow or outflow, as it can often trace local CCA circulation.

Two observed cases are particularly useful anchors. In Fornax A, MeerKAT, ALMA, and MUSE data reveal concurrent feeding and feedback, including turbulent clouds, a ~ 3 kpc filament, and a larger-scale multiphase outflow in the wake of the radio jets (Maccagni et al. 2021). This resembles the mixed inflow, uplift, and circulation seen in the high-turb run, although Fornax A includes additional environmental and merger-related complexity. On smaller scales, molecular absorption in cool-core BCGs reveals cold gas with velocities between -45 and 283 km s^{-1} relative to systemic, likely biased toward motions in the direction of the SMBH (Rose et al. 2019). This is particularly relevant to the inner k-plot interpretation: the low-turb run and the early high-turb *stormy* phase resemble efficiently coupled cold-cloud feeding, whereas the late high-turb *cloudy* phase retains cold gas with weaker final coupling to the sink.

The hot X-ray phase sets the turbulent background for condensation. Cavities, shocks, rims, and low-entropy structures show that AGN feedback acts mainly through mechanical redistribution rather than spatially uniform heating (Fabian 2012; Gitti et al. 2012; McNamara & Nulsen 2012). Direct calorimeter measurements now anchor the velocity scale across different hot-atmosphere states: Hitomi measured $\sigma_{\text{los}} = 164 \pm 10 \text{ km s}^{-1}$ in Perseus, while XRISM found $\sigma_v = 169 \pm 10 \text{ km s}^{-1}$ in A2029, representative of calm subsonic cores (Hitomi Collaboration et al. 2016; XRISM Collaboration et al. 2025). At the disturbed end, XRISM observations of the merging cluster Abell 2034 show much broader lines, with $\sigma_v \approx 470 \text{ km s}^{-1}$, a kinetic pressure fraction of $\sim 15\%$, and a possible turbulent Mach number $\mathcal{M} \sim 0.5$ (Heinrich et al. 2026). Indirect constraints from X-ray density and thermodynamic perturbations likewise indicate predominantly subsonic motions, but with substantial system-to-system variation (Gaspari & Churazov 2013; Hofmann et al. 2016). Thus, our runs should be compared with observations at the level of turbulence regime rather than as synthetic line profiles: low-turb ($\mathcal{M}_{3D} \sim 0.15$) is close to calm Hitomi/XRISM-like cores, whereas high-turb ($\mathcal{M}_{3D} \sim 0.4$) represents a more strongly stirred but still subsonic atmosphere, closer to disturbed XRISM systems such as Abell 2034.

The simulated accretion states are also broadly consistent with mechanically dominated AGN feedback. Both runs peak near $\lambda_{\text{Edd}} \approx 8 \times 10^{-5}$, span roughly $\lambda_{\text{Edd}} \sim 10^{-6} - 10^{-1}$, and spend a significant fraction of their evolution in a sub-Eddington regime. Observed Eddington ratios are inferred from bolometric luminosities, bolometric corrections, emission-line proxies, and black-hole mass estimates (Vestergaard & Peterson 2006; Lamastra et al. 2009; Duras et al. 2020), so our accretion-rate-based λ_{Edd} is only an analogue. Nevertheless, the simulated range is consistent with low-luminosity radio-mode systems and

broad survey-inferred accretion-rate distributions (Kauffmann & Heckman 2009; Aird et al. 2012; Best & Heckman 2012; Heckman & Best 2014), with rare high- λ_{Edd} excursions corresponding to short-lived CCA bursts rather than sustained quasar-like phases (Gaspari et al. 2017; Sądowski & Gaspari 2017). Nuclear wind studies provide a complementary small-scale view of the same feedback cascade (Mehdipour et al. 2023).

The C -ratio profiles provide the closest thermodynamic connection. In the CCA framework, nonlinear condensation is expected when radiative cooling and turbulent mixing operate on comparable timescales, namely $C \equiv t_{\text{cool}}/t_{\text{eddy}} \approx 1$. In Gaspari et al. (2018), the observed condensation locus is centered on $C \approx 1$, with multiphase systems falling within a band close to ± 0.3 dex interval around unity, similar to our reference shaded interval. Observed filaments generally lie within low-entropy, short-cooling-time regions, often where $t_{\text{cool}}/t_{\text{ff}} \lesssim 20$, and in several systems near $t_{\text{cool}}/t_{\text{eddy}} \sim 1$ (Olivares et al. 2019, 2022; Lepore et al. 2025). In our simulations, this $C \sim 1$ band is reached most systematically outside the jet cone. The high-turb run reaches it over a broader radial interval, consistent with extended *stormy/cloudy* filaments, whereas low-turb reaches it in a more compact central region, consistent with *rainy* precipitation.

These comparisons highlight the main observational implication: multiphase gas alone is not sufficient evidence for efficient SMBH feeding. In the high-turb *cloudy* state, cold and warm gas remain visible on meso- and inner macro-scales while the inner inflow and SMBH accretion rate decline. Systems with rich ALMA/MUSE filamentary structure but weak or intermittent instantaneous AGN activity may therefore occupy a weakly coupled *cloudy* state, in which condensation continues but only a small fraction of the cold phase reaches the central engine.

5. Conclusions

In this work, we have analyzed the kinematics, radial inflow/outflow structure, accretion variability, and CCA diagnostics of a jet-regulated multiphase atmosphere in a galaxy-group environment. This paper complements C26a, which focuses on the morphology and thermodynamics of the same simulations. Within the BLACKHOLEWEATHER framework, the setup builds on the CCA and self-regulated feedback picture developed in Gaspari et al. (2011, 2013, 2017, 2018) and on the no-jet turbulent CCA baseline of B26a,b. We examined two benchmark hydrodynamical simulations including radiative cooling, driven subsonic turbulence, pc-scale sink accretion, and a self-regulated kinetically-dominated jet. The two runs, high-turb and low-turb, differ only in the imposed turbulent driving strength, allowing us to isolate how ambient turbulence modifies jet-regulated feeding and feedback. Combined with C26a, our main conclusions are as follows.

1. Ambient turbulence regulates both the onset and the geometry of jet-regulated CCA. In high-turb, condensation starts later but develops into a broader, more filamentary, and more strongly mixed multiphase structure. In low-turb, condensation starts earlier and remains more coherent, centrally concentrated, and persistent.
2. The main difference between the two weather regimes is inner transport efficiency, not the large-scale gas supply. Both runs maintain substantial inflow and circulation at meso- and macro-scales, but only low-turb preserves strong coupling from meso- to micro-scales. In the late high-turb evolution, multiphase gas remains abundant outside the centre, while the final inward transport to the sink is weakened.
3. Once precipitation begins, both runs become genuinely CCA-fed. The accretion rate rises from Bondi-like values to strongly super-Bondi levels ($\dot{M}_*/\dot{M}_{\text{Bondi}} \sim 10^{2-3}$), showing that smooth hot-mode accretion is replaced by multiphase, precipitation-driven feeding. The low-turb run sustains this state for longer, whereas high-turb undergoes a short, intense *stormy* burst followed by a weaker and more intermittent *cloudy* mode. Both simulations remain mostly in a low-Eddington, mechanically dominated regime, with sporadic episodes at larger accretion states.
4. Accretion variability retains the flickering character expected for CCA. The accretion-rate PSDs show flicker-like low-frequency slopes and steep red-noise high-frequency tails, consistent with scale-coupled multiphase feeding. Compared with the no-jet B26b baseline, the two jet-regulated runs show more similar break frequencies, suggesting that jet-driven stirring partly reduces the direct dependence of the variability timescale on the imposed ambient turbulence. In high-turb, the transition to the *cloudy* phase lowers the PSD normalization and flattens the low-frequency slope, indicating that coherent, large-scale rain episodes are replaced by weaker and less correlated feeding events.
5. The phase-separated mass fluxes show how the weather state controls cold-gas transport. In high-turb, molecular-temperature, cold, and warm gas become more radially extended and participate in uplift, fountain-like recycling, and cloudy circulation out to kpc and inner-macro scales. In low-turb, cool recycling remains mostly confined to the inner kpc, favoring a compact reservoir that feeds the SMBH more persistently.
6. The k-plot and the C -ratio profiles provide complementary diagnostics of whether cold gas is dynamically linked to SMBH feeding. The C -ratio profile identifies where the hot atmosphere is thermodynamically prone to condensation, with $C = t_{\text{cool}}/t_{\text{eddy}} \sim 1$, and shows that this condition is reached most systematically outside the jet cone and around the jet-ambient interface. The k-plot shows whether the resulting cold and warm gas is coherent, turbulent, drifting, or dispersed. Neither diagnostic alone is sufficient: cold gas can form where condensation is allowed, yet still be dynamically decoupled from the sink. Together, the two diagnostics show that high-turb develops a broader condensation-prone region but more dispersed and weakly coupled cold gas, whereas low-turb maintains a more compact $C \sim 1$ region and more coherent cold/warm kinematics, favoring sustained feeding.

Overall, jet feedback acts primarily as a dynamical reorganization mechanism rather than as simple local heating. The kinetic jet excavates a hot, low-density channel, suppresses sustained in-situ condensation inside the cone, and promotes uplift, entrainment, mixing, and circulation in the surrounding atmosphere and jet-ambient interface. The *meso-scale* therefore does not merely host multiphase condensation: it sets the transport bottleneck between halo cooling and SMBH feeding. The presence of cold gas alone is insufficient to infer efficient accretion; what matters is whether the local weather state allows condensed gas to remain thermodynamically favoured, kinematically coherent, and dynamically connected to the inner inflow. In this sense, the BLACKHOLEWEATHER framework (Gaspari et al. 2020) provides the physical language and diagnostic structure needed to connect macro-scale halo cooling, meso-scale multiphase transport, and micro-scale SMBH feeding. This extends the B26a,b no-jet CCA results to the jet-regulated regime and

motivates the combined use of k-plot and C-ratio diagnostics to interpret multiphase AGN feeding and feedback.

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Appendix A: Synthetic observational mapping

Table A.1. Synthetic mapping between simulation diagnostics and observational proxies. The entries summarize how jet-regulated CCA weather states may appear in multiphase observations. They are intended as an interpretive guide, not as forward-modelled observables.

Simulation diagnostic	Observable proxy	Observed anchor	Weather-state cue
<i>Multiphase morphology</i>			
Cold/warm extent and distribution	CO, H α , [N II] morphology; compact reservoirs; off-centre clouds	Cold/warm gas ranges from compact nuclear reservoirs to kpc-scale filaments, often associated with cavities, rims, or uplifted low-entropy gas. ^a	high-turb: extended <i>stormy</i> or <i>cloudy</i> filaments. low-turb: compact <i>rainy</i> reservoir.
<i>Kinematics and circulation</i>			
Projected cold/warm motions	CO, H α , [N II] centroids; absorption components	Offsets of order 10 ² km s ⁻¹ are common in cold/warm filaments and nuclear absorbers, often below pure free-fall expectations. ^b	Large offsets with modest dispersion suggest uplift, fallback, or weakly coupled <i>cloudy</i> circulation.
Velocity dispersion / k-plot	CO and optical line widths; kinematic phase-space location	The adopted k-plot thresholds, $ v_{\text{los}} - v_{\text{sys}} = 100 \text{ km s}^{-1}$ and $\sigma_{\text{los}} = 50 \text{ km s}^{-1}$, sample the observed multiphase range. ^b	<i>stormy</i> : broad/turbulent loci. <i>rainy</i> : compact coherent loci. <i>cloudy</i> : drifting or dispersed gas.
Inflow/outflow ambiguity	Line asymmetry, centroid shifts, gas near bubbles and radio structures	Cold/warm gas near cavities or jets can trace uplift, fallback, local fountains, or feeding, rather than a single radial flow. ^b	CCA circulation: inflow, uplift, mixing, and fallback coexist.
<i>Hot atmosphere and condensation</i>			
Hot-phase stirring	X-ray calorimeter line broadening; X-ray fluctuations; cavities and rims	Direct spectroscopy and fluctuation analyses constrain predominantly subsonic hot-atmosphere motions, even in disturbed systems. ^c	low-turb: calmer atmosphere, compact precipitation. high-turb: stronger stirring, broader circulation.
$C = t_{\text{cool}}/t_{\text{eddy}}$	Cooling time plus turbulence constraints	CCA condensation is expected within an interval around unity $C \sim 0.5 - 2$. ^d	Extended $C \sim 1$: <i>stormy/cloudy</i> condensation. Compact central $C \sim 1$: <i>rainy</i> condensation.
<i>AGN state</i>			
Accretion level and variability	Nuclear luminosity, cavity power, radio duty cycle, long-term monitoring	Observed λ_{Edd} estimates are luminosity-based; radio-mode systems trace mechanically dominated, low-Eddington feedback with intermittent activity. ^e	<i>stormy</i> : bursty feeding. <i>rainy</i> : sustained coupling. <i>cloudy</i> : weak instantaneous feeding.

Notes. ^(a) Multiphase morphology and filament extent: Tremblay et al. (2016, 2018); Temi et al. (2018); Olivares et al. (2019); Vantyghem et al. (2019); Olivares et al. (2022, 2025); Russell et al. (2019). ^(b) Cold/warm kinematics, absorption, and k-plot diagnostics: Gaspari et al. (2018); Tremblay et al. (2018); Olivares et al. (2019); Rose et al. (2019); Maccagni et al. (2021); Singha et al. (2023). ^(c) Hot-atmosphere feedback and turbulence constraints: Fabian (2012); Gitti et al. (2012); McNamara & Nulsen (2012); Gaspari & Churazov (2013); Hofmann et al. (2016); Juráňová et al. (2020); Hitomi Collaboration et al. (2016); XRISM Collaboration et al. (2025); Heinrich et al. (2026). ^(d) CCA condensation thresholds and cooling/turbulence diagnostics (and hot/warm filament structure): Gaspari et al. (2018); Olivares et al. (2019, 2022, 2025); Lepore et al. (2025). ^(e) AGN accretion-state and mechanical-feedback context: Vestergaard & Peterson (2006); Lamastra et al. (2009); Kauffmann & Heckman (2009); Aird et al. (2012); Best & Heckman (2012); Heckman & Best (2014); Duras et al. (2020); Gaspari et al. (2017, 2020); Sądowski & Gaspari (2017); Mehdipour et al. (2023).

Appendix B: Time-radius maps

To further support the weather-phase distinction, Fig. B.1 shows radius–time 2D histograms of the inflow (top panels) and outflow (bottom panels) mass rates, \dot{M}_{in} and \dot{M}_{out} , for the combined molecular-temperature and cold gas phase ($T < 1.6 \times 10^4 \text{ K}$). The maps are computed as total mass fluxes in 50 radial shells over $0 < r < 50 \text{ kpc}$ and sampled every 1 Myr, with the high-turb run shown on the left and the low-turb run on the right. They highlight the difference between the late *cloudy* state in the fixed-axis high-turb run and the more prevalent *sunny* clearing reported in P26b. In the present *cloudy* state, cold gas is still present from micro- to inner-macro scales even at $t/t_{\text{rain}} \gtrsim 6$, but its coupling to central accretion is inefficient. In the P26b *sunny* state, by contrast, jet-axis reorientation more persistently suppresses the cold component across the inner transport region.

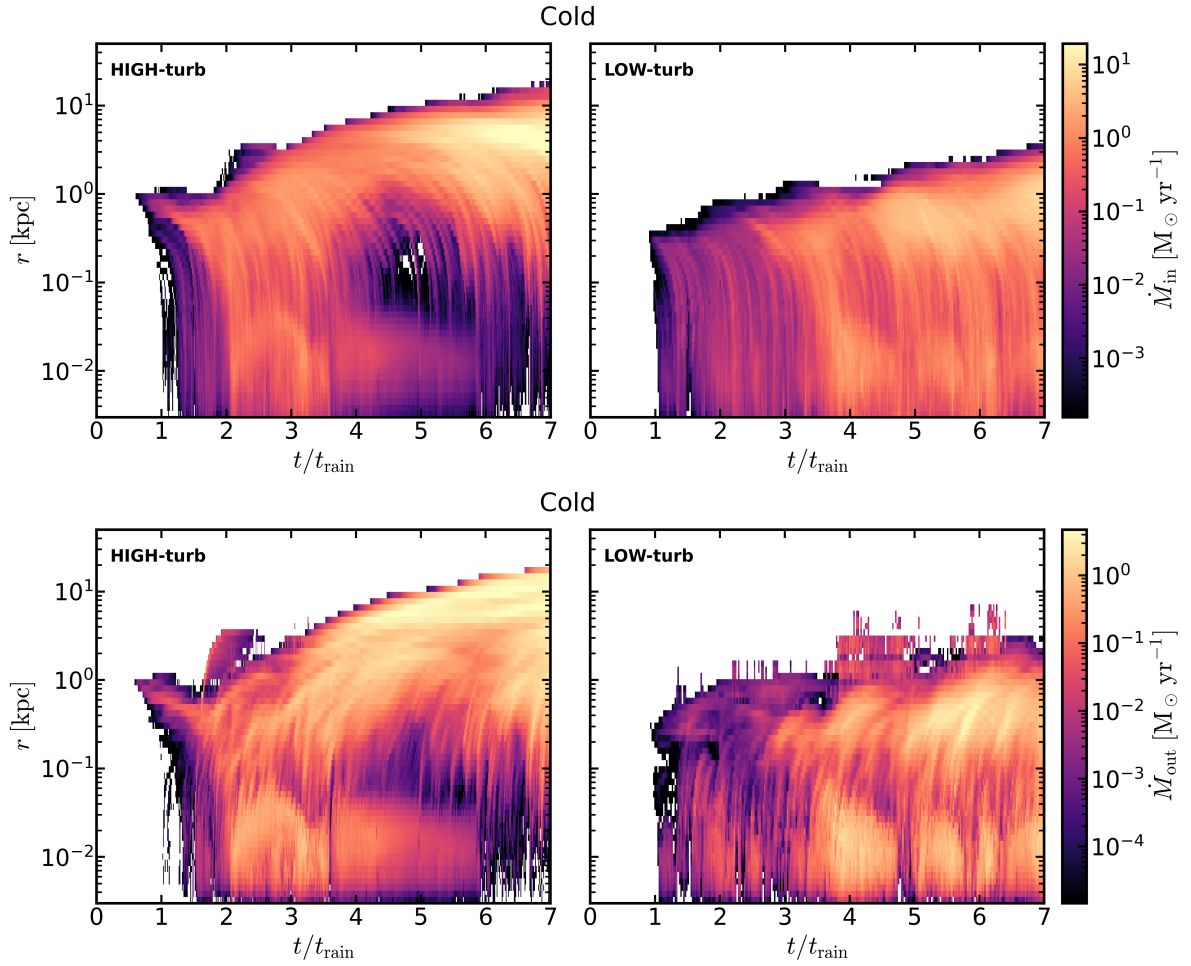


Fig. B.1. Radius–time 2D histograms of the inflow (upper panels) and outflow (lower panels) mass rates for the combined molecular-temperature and cold gas phase ($T < 1.6 \times 10^4$ K). The **high-turb** run is shown in the left panels and the **low-turb** run in the right panels. Mass rates are computed as the total inflowing or outflowing mass flux in each radial shell, using 50 radial bins sampled every 1 Myr.